



Ab initio local field effects for surface second-harmonic generation

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Response to a perturbation

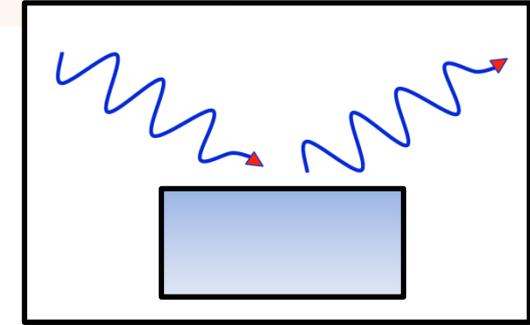
Perturbation

Electric field



Response

Polarisation



Linear Response

Nonlinear Response

$$P_i = \epsilon_0 \sum_j \chi_{ij}^{(1)} E_j + \epsilon_0 \sum_{jk} \chi_{ijk}^{(2)} E_j E_k + \epsilon_0 \sum_{jkl} \chi_{ijkl}^{(3)} E_j E_k E_l + \dots$$

Response to a perturbation

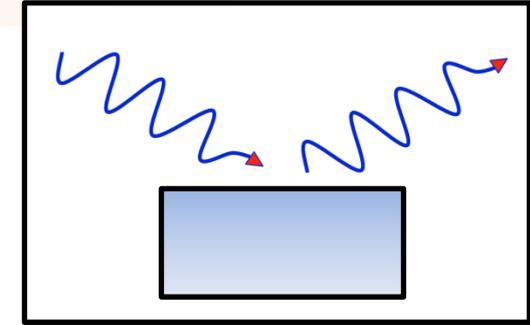
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$$\chi_{ij}^{(1)}$$

$$\chi_{ijk}^{(2)}$$

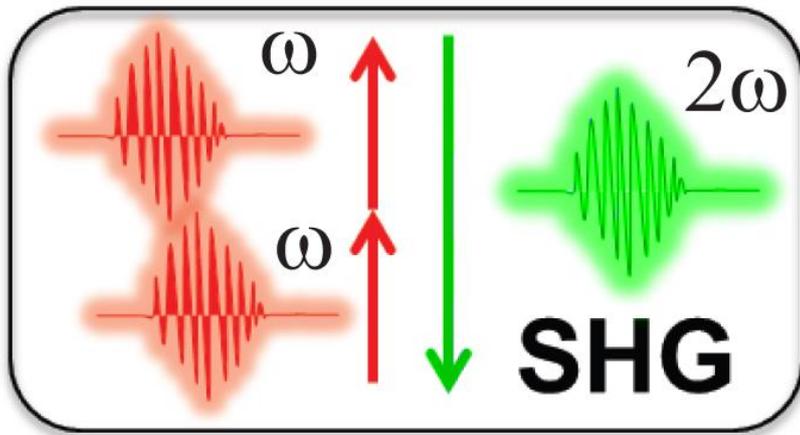
$$\chi_{ijkl}^{(3)}$$

- Absorption
- Refraction
- Loss function
- Second-harmonic generation
- Pockels effect
- ...
- Kerr effect
- Four-wave mixing
- ...

Second harmonic generation

$$P_i = \epsilon_0 \sum_j \chi_{ij}^{(1)} E_j + \epsilon_0 \sum_{jk} \chi_{ijk}^{(2)} E_j E_k + \epsilon_0 \sum_{jkl} \chi_{ijkl}^{(3)} E_j E_k E_l + \dots$$

First nonlinear term



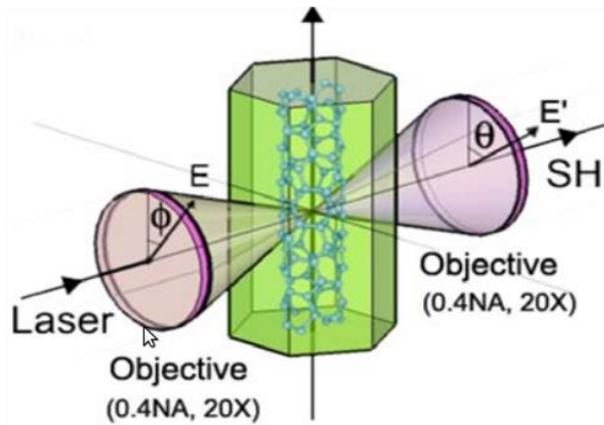
Dipole approximation

Centrosymmetric material : $\chi^{(2)} = 0$

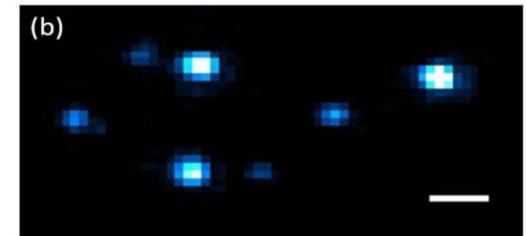
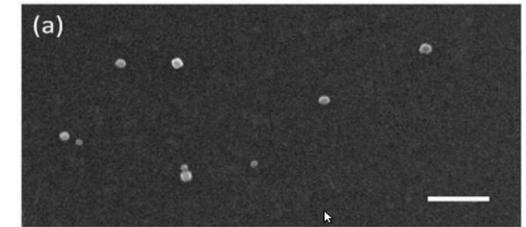
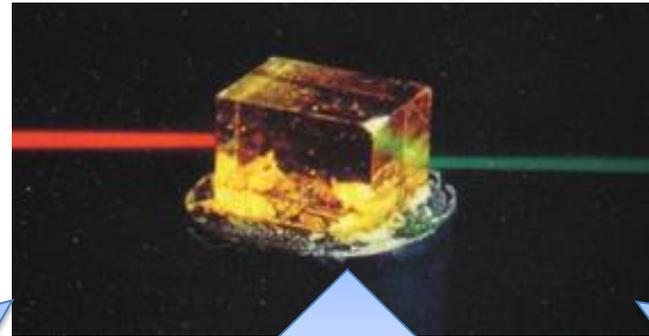
First nonlinear term : $\chi^{(3)}$

Non centrosymmetric material : $\chi^{(2)} \neq 0$

Applications of second harmonic generation (SHG)

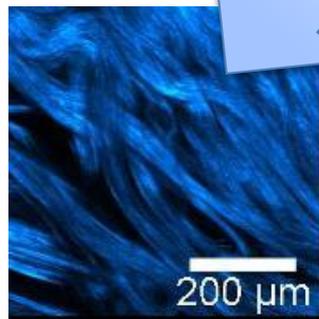
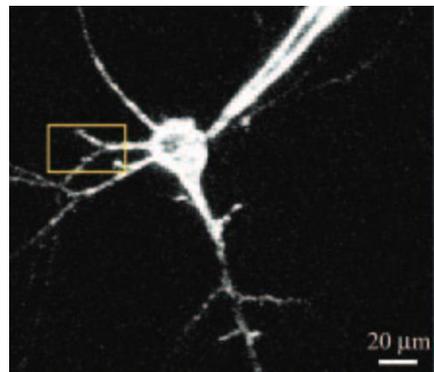
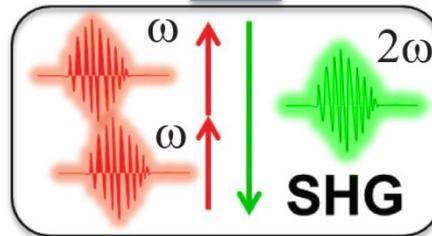


Nanotubes characterization
(*PRB 77 125428*)

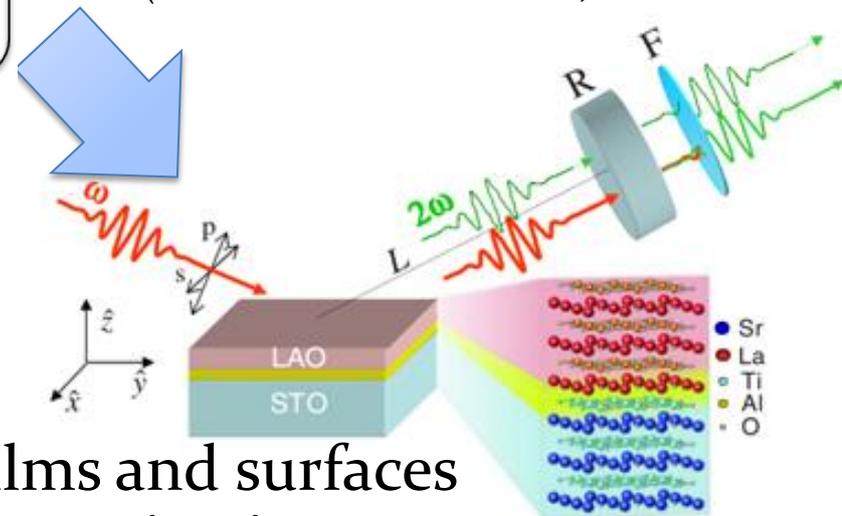


Nanoparticles imaging and microscopy

(*C-L Hsieh PhD thesis, Caltech 2011*)



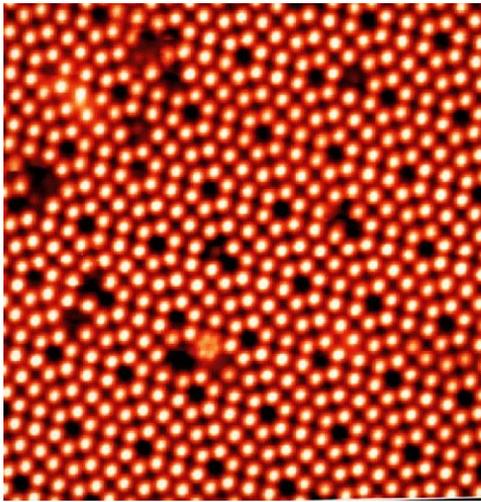
Biological tissues/neurons imaging
(*Biophys. J. 81 493; PNAS 103, 786*)



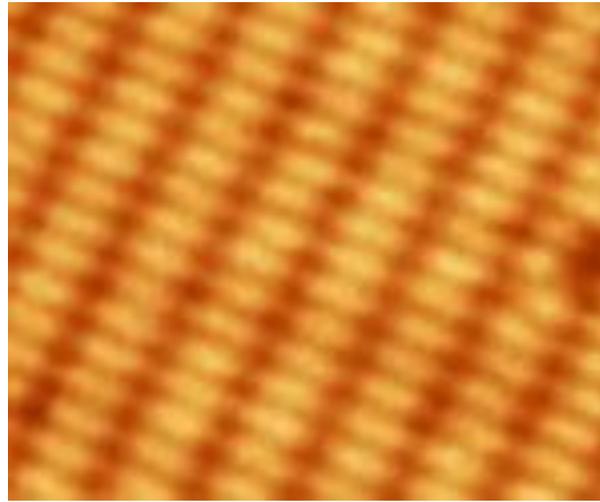
Thin films and surfaces characterization (*PRB 89 075110*)

Surfaces

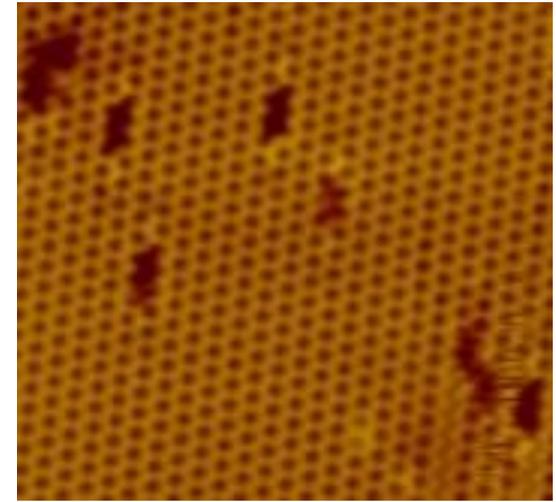
Different surfaces for the same material (e.g. Silicon)



Si(111) 7x7



Si(001) 2x1



Si(001) 4x2

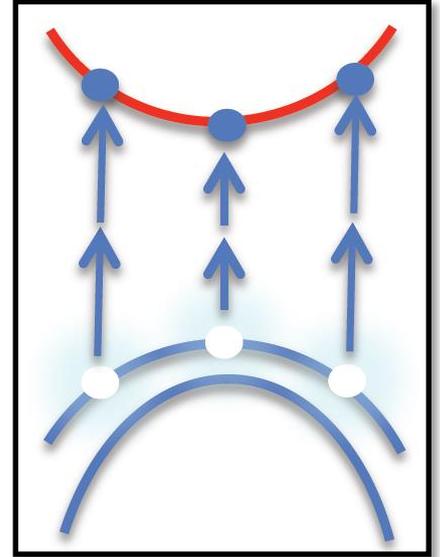
Surface = symmetry breaking

How do we get the spectrum for SHG

Second-order response

Independent Particles Approximation

All the electrons make independent transitions



$$\chi_{abc}^{(2)}(-2\omega, \omega, \omega) = \frac{-ie3}{\hbar^2 m^3 \omega^3 V} \sum_{nml} \int d\vec{k} \frac{1}{E_m - E_n - 2\omega - 2i\eta}$$
$$\times \left[f_{nl}(\vec{k}) \frac{p_{nm}^a(\vec{k}) \{ p_{ml}^b(\vec{k}) p_{ln}^c(\vec{k}) \}}{E_l - E_n - \omega - i\eta} + f_{ml}(\vec{k}) \frac{p_{nm}^a(\vec{k}) \{ p_{ml}^b(\vec{k}) p_{ln}^c(\vec{k}) \}}{E_m - E_l - \omega - i\eta} \right]$$

 2light

Local field effects

Total perturbation

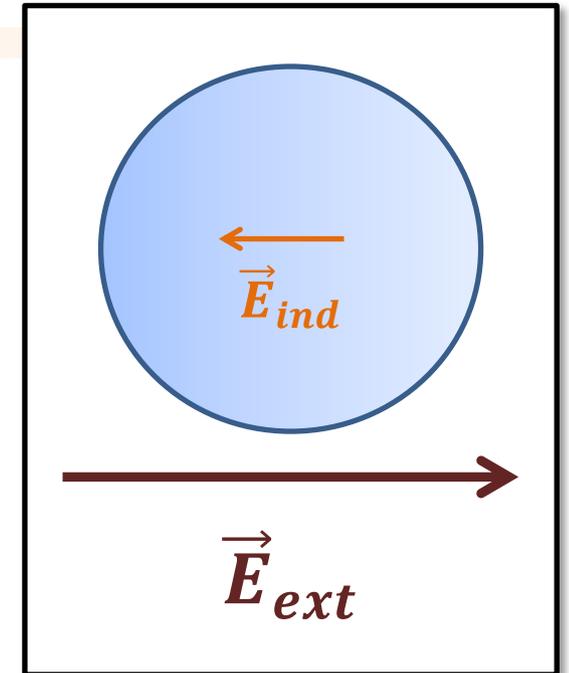


External perturbation
Macroscopic external field



Local fields

Microscopic fluctuations of the electric field on the nanoscale



Local field effects:
Difference between the microscopic and macroscopic responses

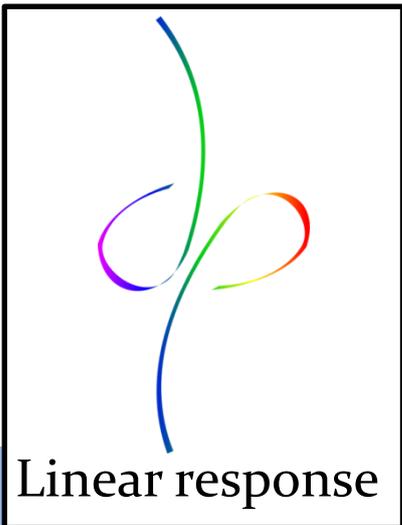
Macroscopic response and local field effects in TDDFT

Linear and Second-order Response Function in the framework of TDDFT

Time-Dependent Density Functional Theory

1st order $\left[1 - \chi_0^{(1)} (v + f_{xc}) \right] \chi^{(1)} = \chi_0^{(1)}$

$$f_{xc} = \frac{\partial V_{xc}}{\partial \rho}$$



Macroscopic response and local field effects in TDDFT

Linear and Second-order Response Function in the framework of TDDFT

Time-Dependent Density Functional Theory

1st order
$$\left[1 - \chi_0^{(1)}(v + f_{xc}) \right] \chi^{(1)} = \chi_0^{(1)}$$

$$f_{xc} = \frac{\partial V_{xc}}{\partial \rho}$$

2nd order

$$\begin{aligned} [1 - \chi_0^{(1)}(2\omega) f_{uxc}(2\omega)] \chi^{(2)}(2\omega, \omega) &= \chi_0^{(2)}(2\omega, \omega) \left[1 + f_{uxc}(\omega) \chi^{(1)}(\omega) \right]^2 \\ &+ \chi_0^{(1)}(\omega) g_{xc}(\omega) \chi^{(1)}(\omega) \chi^{(1)}(\omega) \end{aligned}$$



Linear response

$$g_{xc} = \frac{\partial^2 V_{xc}}{\partial \rho \partial \rho}$$



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Second-harmonic generation

Second-order response

State-of-the-art

Bulk materials and interfaces

- Independent particles - end of 1980's
- Local field and excitonic effects - TDDFT (2010) and Bethe-Salpeter equation (2005)

Surfaces

- Some calculations in independent particles - since 1994
- Effect of nonlocal operators (“scissors” but wrong formula) – since 2001
- Effect of nonlocal operators (“scissors” and pseudo-potential) – [1] 2015

[1] S. Anderson, N. T-D, B. Mendoza and V. Véniard, PRB 91, 07530 (2015)

Model of surface – Super-cells

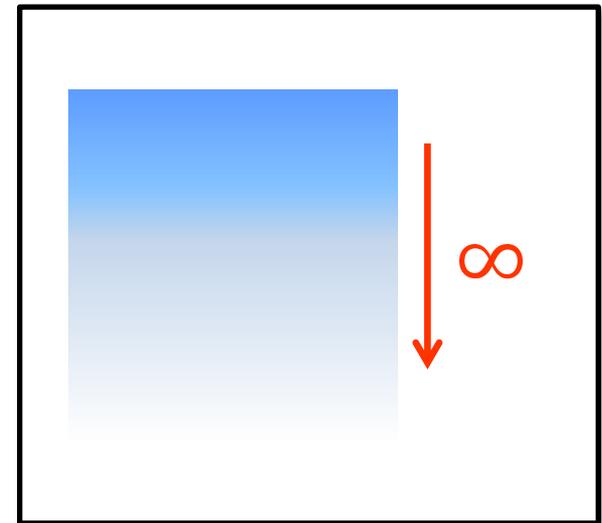
Surfaces

- Periodic in 2 directions
- Aperiodic in one direction

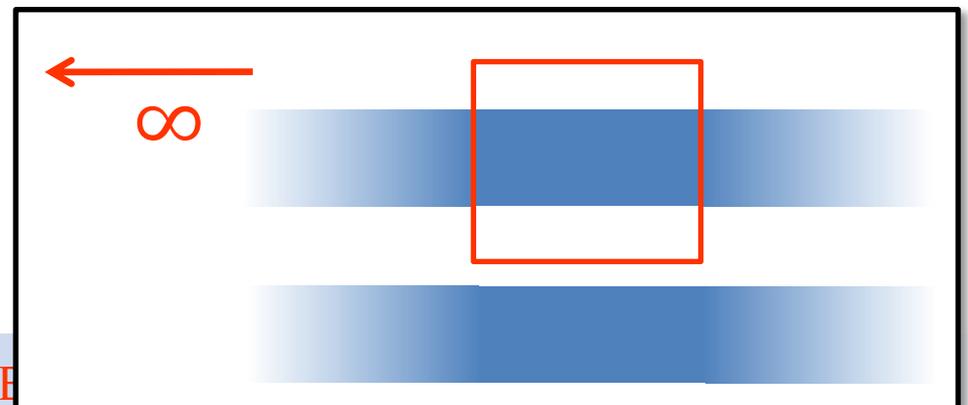
Periodic \Rightarrow Plane-waves

Aperiodic \Rightarrow Supercells (atoms + vacuum)

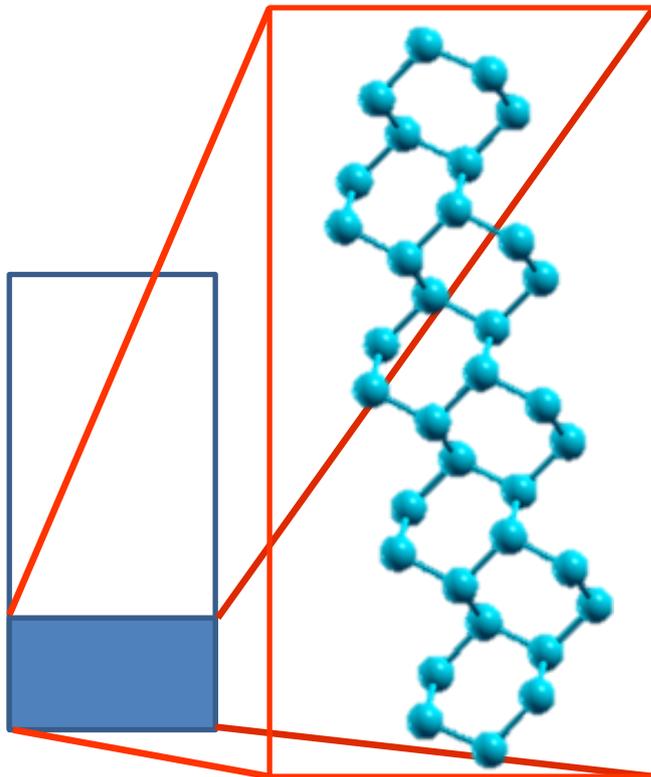
Semi-infinite material



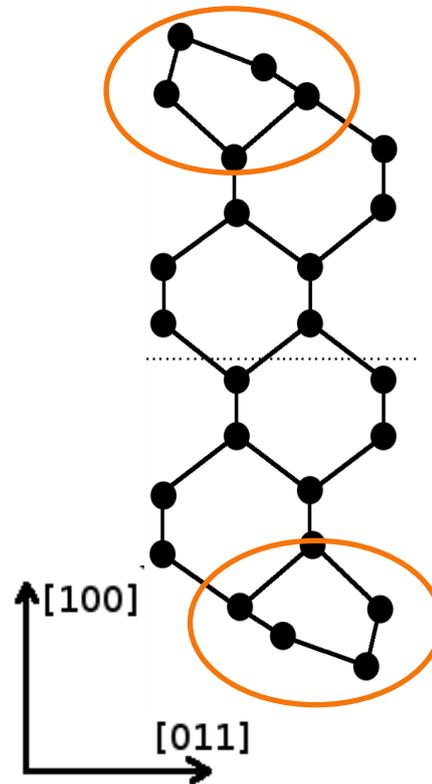
System with 2 surfaces (slab)



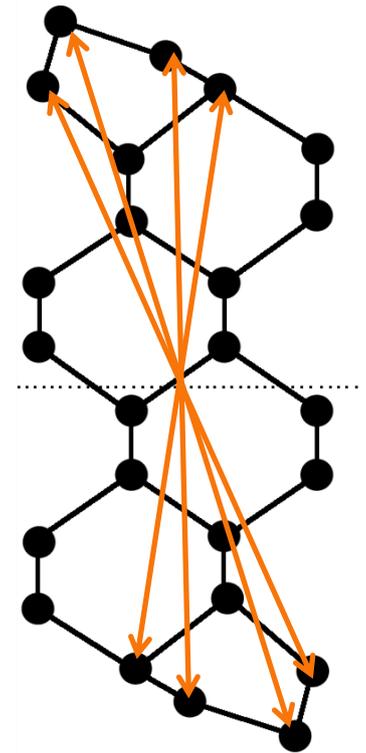
Model of surface – Super-cell



Construction of super-cell
(atoms + vacuum)



System with 2
surfaces (slab)



inversion symmetry

$$\chi^{(2)} = 0$$

(artificial)

Surface second-order response in Time-Dependent DFT

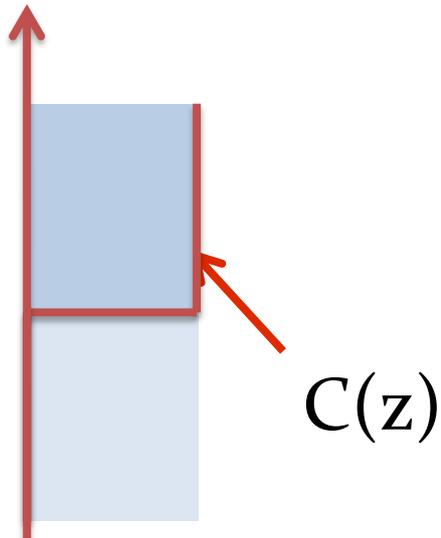
$$\chi_{abc}^{(2)}(-2\omega, \omega, \omega) = \frac{-ie^3}{\hbar^2 m^3 \omega^3 V} \sum_{nml} \int d\vec{k} \frac{1}{E_m - E_n - 2\omega - 2i\eta}$$

$$\times \left[f_{nl}(\vec{k}) \frac{\tilde{p}_{nm}^a(\vec{k}) \left\{ p_{ml}^b(\vec{k}) p_{ln}^c(\vec{k}) \right\}}{E_l - E_n - \omega - i\eta} + f_{ml}(\vec{k}) \frac{\tilde{p}_{nm}^a(\vec{k}) \left\{ p_{ml}^b(\vec{k}) p_{ln}^c(\vec{k}) \right\}}{E_m - E_l - \omega - i\eta} \right]$$

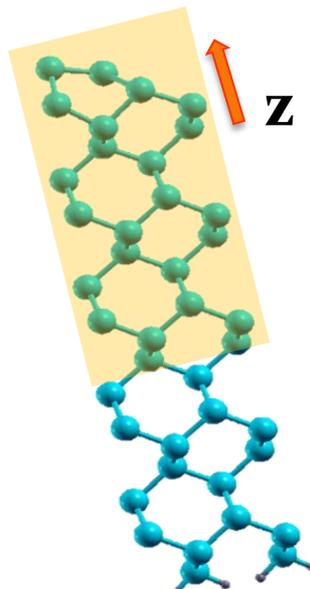
$$\tilde{p} = \frac{1}{2} (pC(z) + C(z)p)$$

$$p = i[H, r]$$

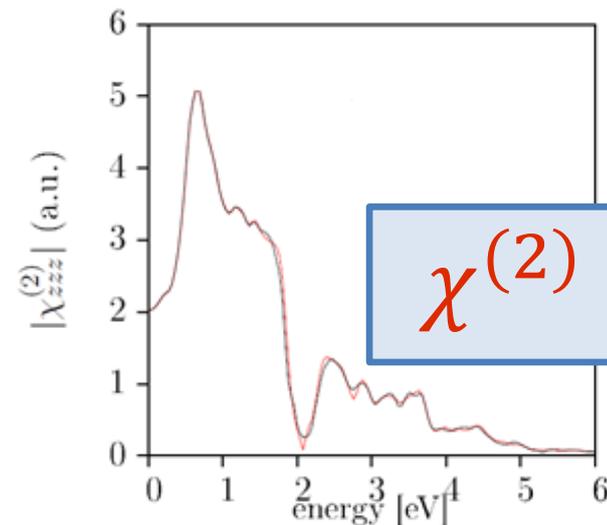
$C(z)$ function



ab initio calculation



Extracted Signal



S. Anderson, N. T-D, B. Mendoza and V. Véniard, PRB 91, 07530 (2015)

DPG Spring Meeting – Berlin - March 2015

Macroscopic response and local field effects

Macroscopic surface response

$$\chi_M^{(2)S}(\mathbf{q}, \mathbf{q}_1, \mathbf{q}_2, \omega_1, \omega_2) = \frac{-i}{2|\mathbf{q}||\mathbf{q}_1||\mathbf{q}_2|} \chi_{\rho\rho\rho}^S(\mathbf{q}, \mathbf{q}_1, \mathbf{q}_2, \omega_1, \omega_2) \\ \epsilon_M^S(\mathbf{q}, \omega) \epsilon_M(\mathbf{q}_1, \omega_1) \epsilon_M(\mathbf{q}_2, \omega_2)$$

2nd order Dyson equation

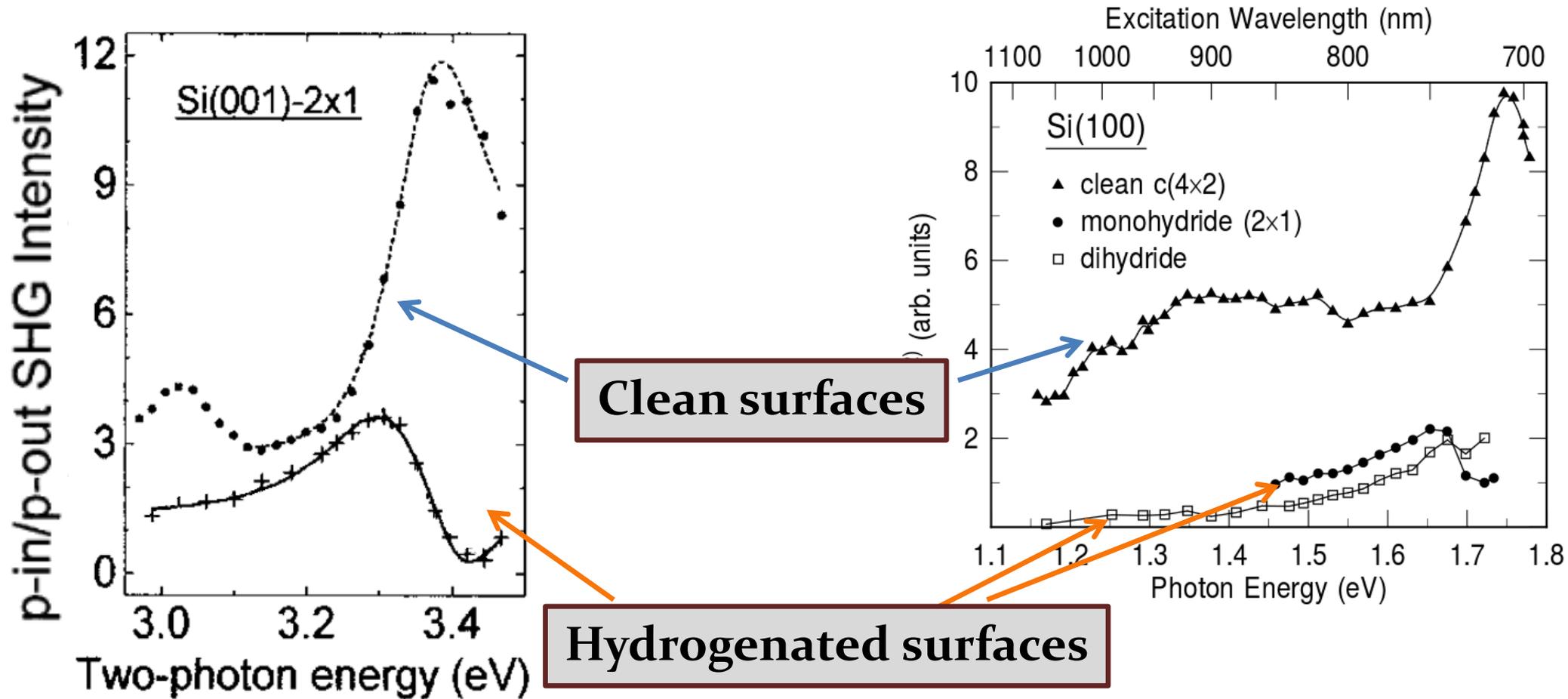
$$\left[Id - \chi_{\rho\rho}^{(0)S} v \right] \chi_{\rho\rho\rho}^{(2)S} = \chi_{\rho\rho\rho}^{(0)S} \left[Id + v \chi_{\rho\rho}^{(1)} \right] \left[Id + v \chi_{\rho\rho}^{(1)} \right]$$

$$\text{with } \left[Id - \chi_{\rho\rho}^{(0)} v \right] \chi_{\rho\rho}^{(1)} = \chi_{\rho\rho}^{(0)}$$

Including local field effects (LFE)

[1] N. T-D and V. Véniard, *in preparation*

Some experimental results for Si(001)

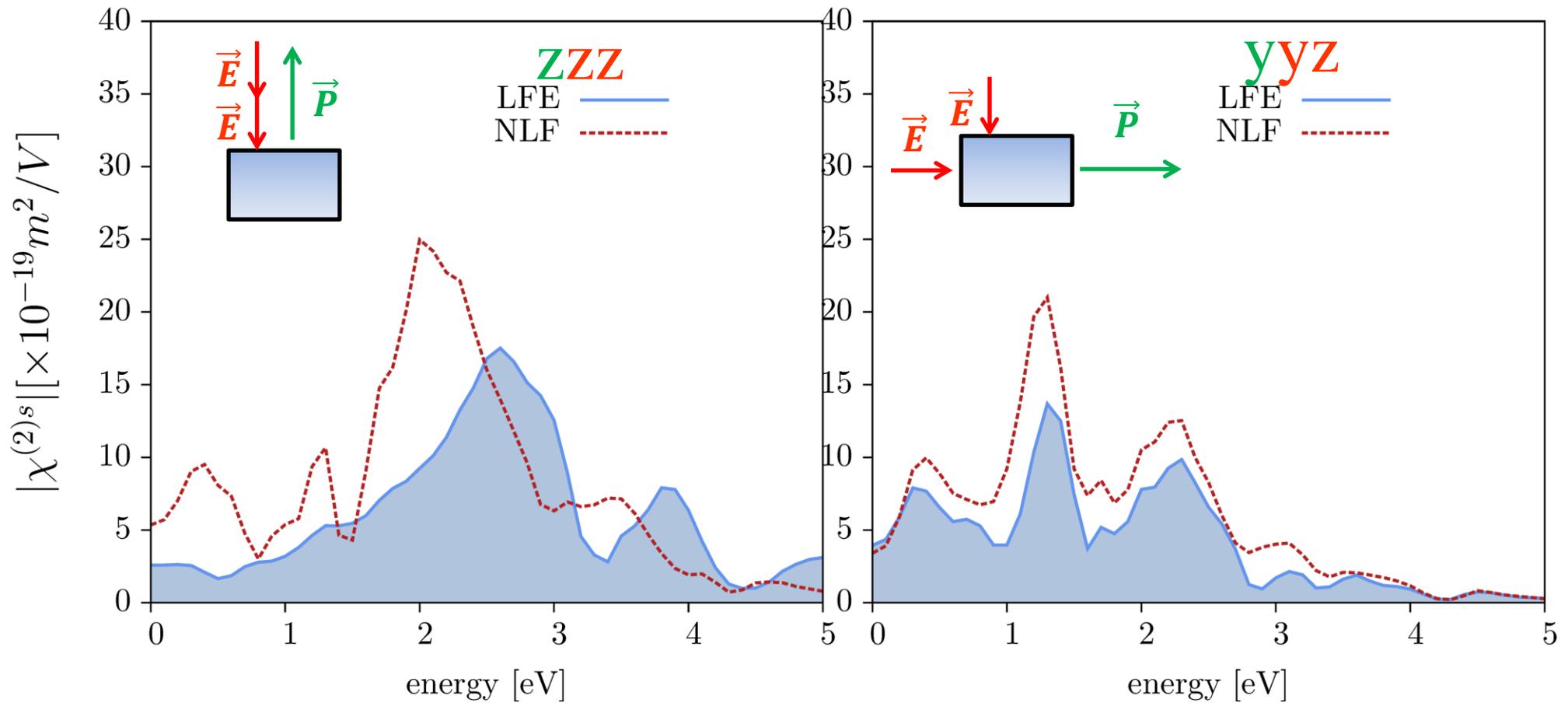


Adapted from
J. I. Dadap, *et al.*, Phys. Rev. B 56, 13367
(1997)

Taken from
U. Höfer, Appl. Phys. A 63, 533 (1996)

Local field effects (LFE) on $\chi^{(2)s}$

Clean Surface

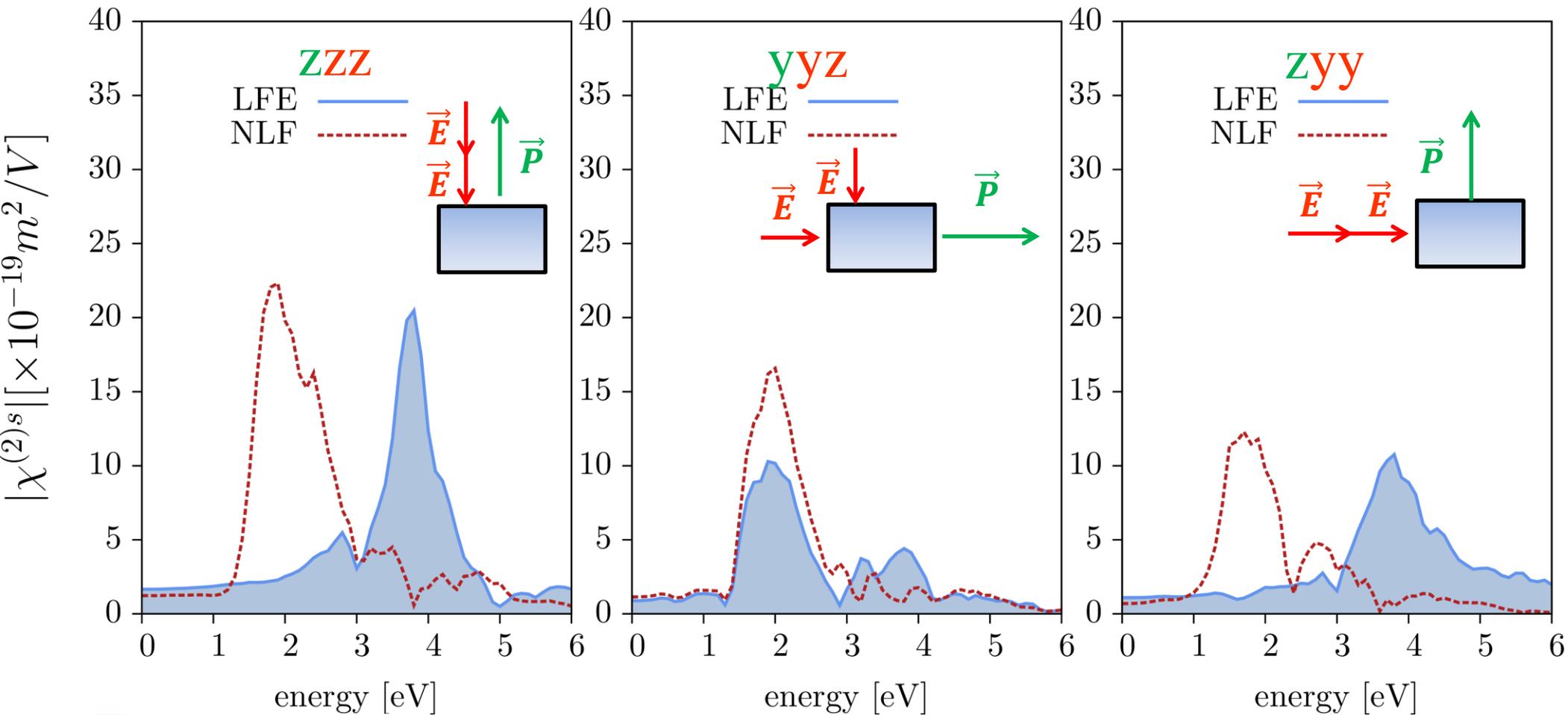


 2light

Si(001)2x1
24 Si layers – 256 k-points

Local field effects (LFE) on $\chi^{(2)s}$

Hydrogenated surface

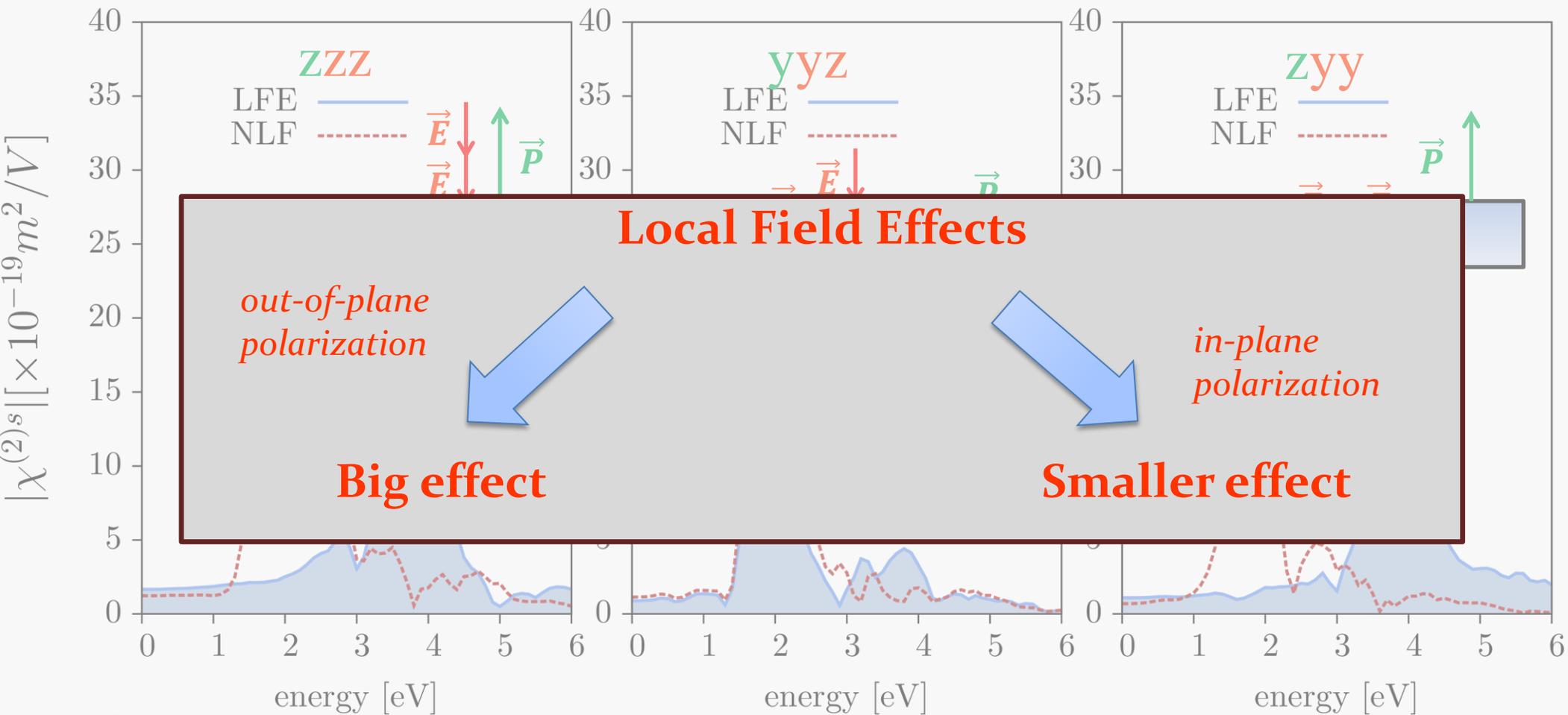


 2light

Si(001)_{1x1:2H}
 24 Si layers – 288 k-points

Local field effects (LFE) on $\chi^{(2)s}$

Hydrogenated surface

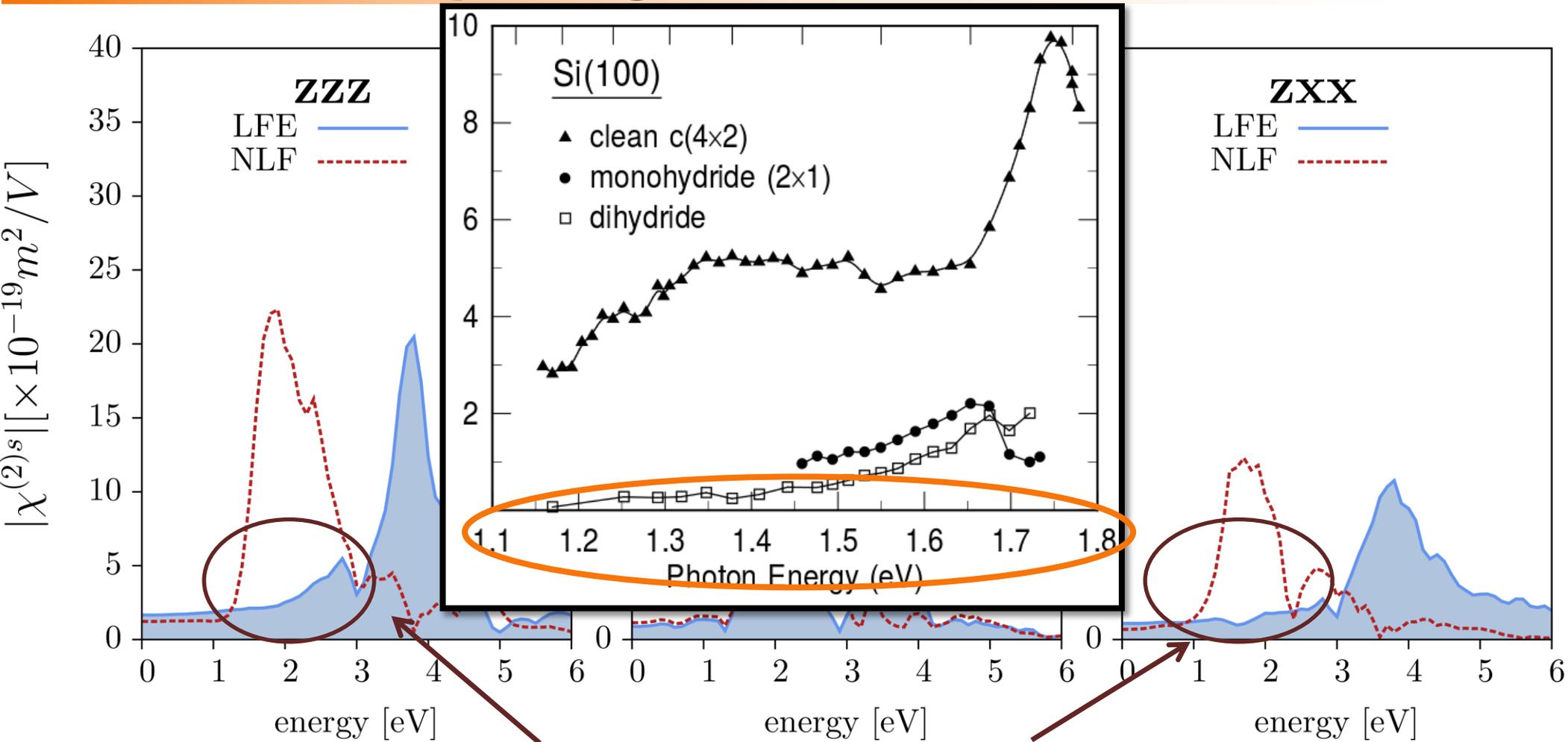


 2light

Si(001)1x1:2H
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Local field effects (LFE) on $\chi^{(2)s}$

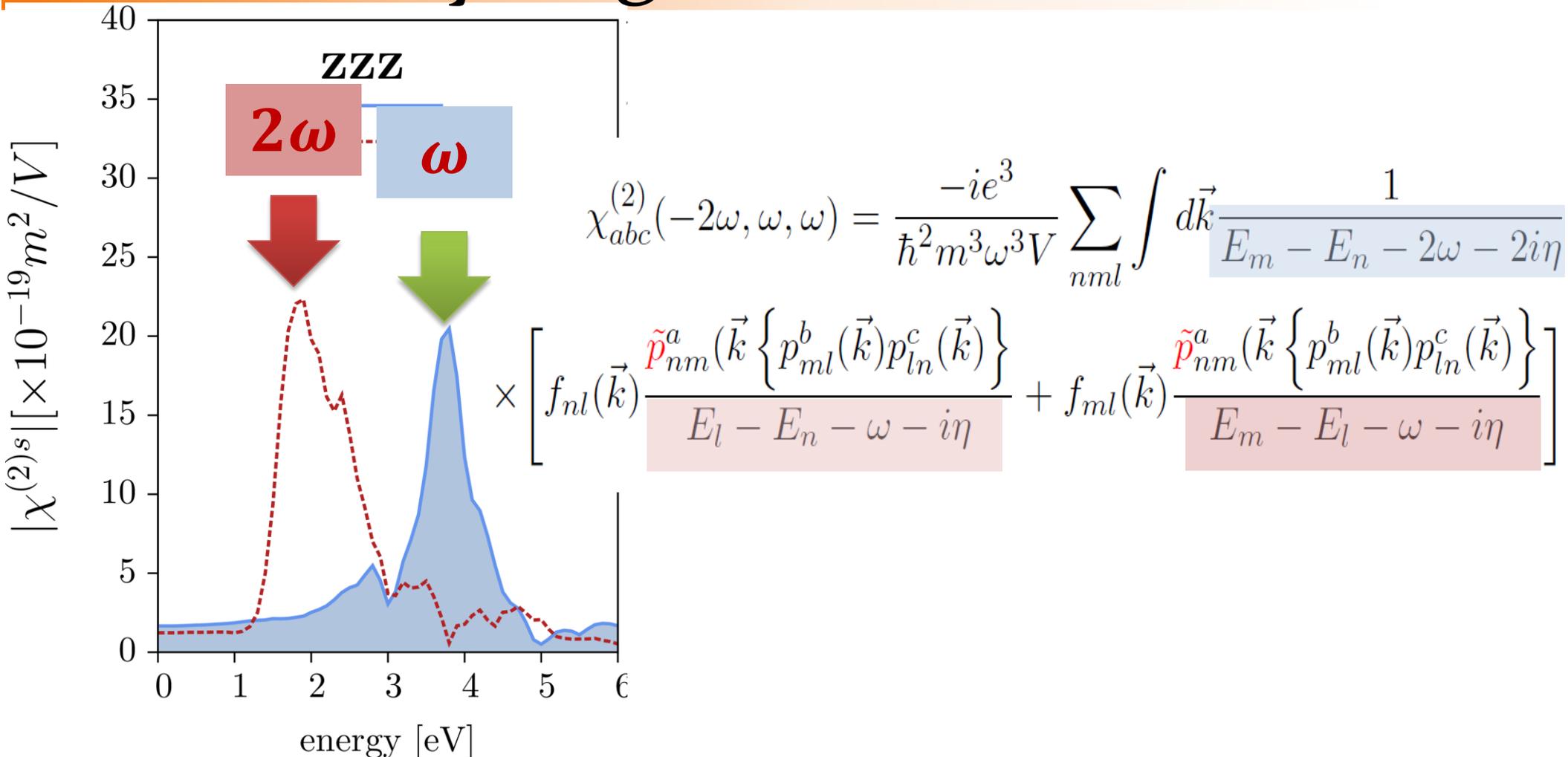
Hydrogenated surface



Strong quenching in region probed experimentally

Local field effects (LFE) on $\chi^{(2)s}$

Hydrogenated surface



**Redistribution of spectral weight
between ω and 2ω transitions**

Conclusion

Conclusion

- Formalism for including the local field effects on surface second harmonic generation (SSHG)
- Local field effects are important for describing SSHG
- Possible explanation of the difference between the clean and the hydrogenated surface signal

Perspective

- Further analysis of hydrogen and ad-atom effects
- Inclusion of other many-body effects

Thank you for your attention



Model of surface – Super-cells

What we want

