

Excitonic signatures in different spectroscopies from optical to scattering experiments, via ab initio many-body approaches

Francesco Sottile

LSI, Ecole Polytechnique, Palaiseau and ETSF - France

UK- France Seed meeting Edinburgh 30-31 May 2024
Quantum Effects in Energy Harvesting





ANR'S ROLE IN RESEARCH

CALLS FOR PROPOSALS

FUNDED PROJECTS AND IMPACT

FRANCE 2030

/ Funded projects and Impact / Search for a funded project / Funded projects



Stratégie nationale

PEPR Technologies Avancées des Systèmes Energétiques (TASE)

Multimodal approach combining IN-situ, ex-situ and Operando CharacTerizAtion with SimUlations for Highly RELiable Next Generation Photovoltaics

MINOTAURE



GENERAL INFORMATION

Acronym: MINOTAURE

Reference Number: 22-PETA-0015

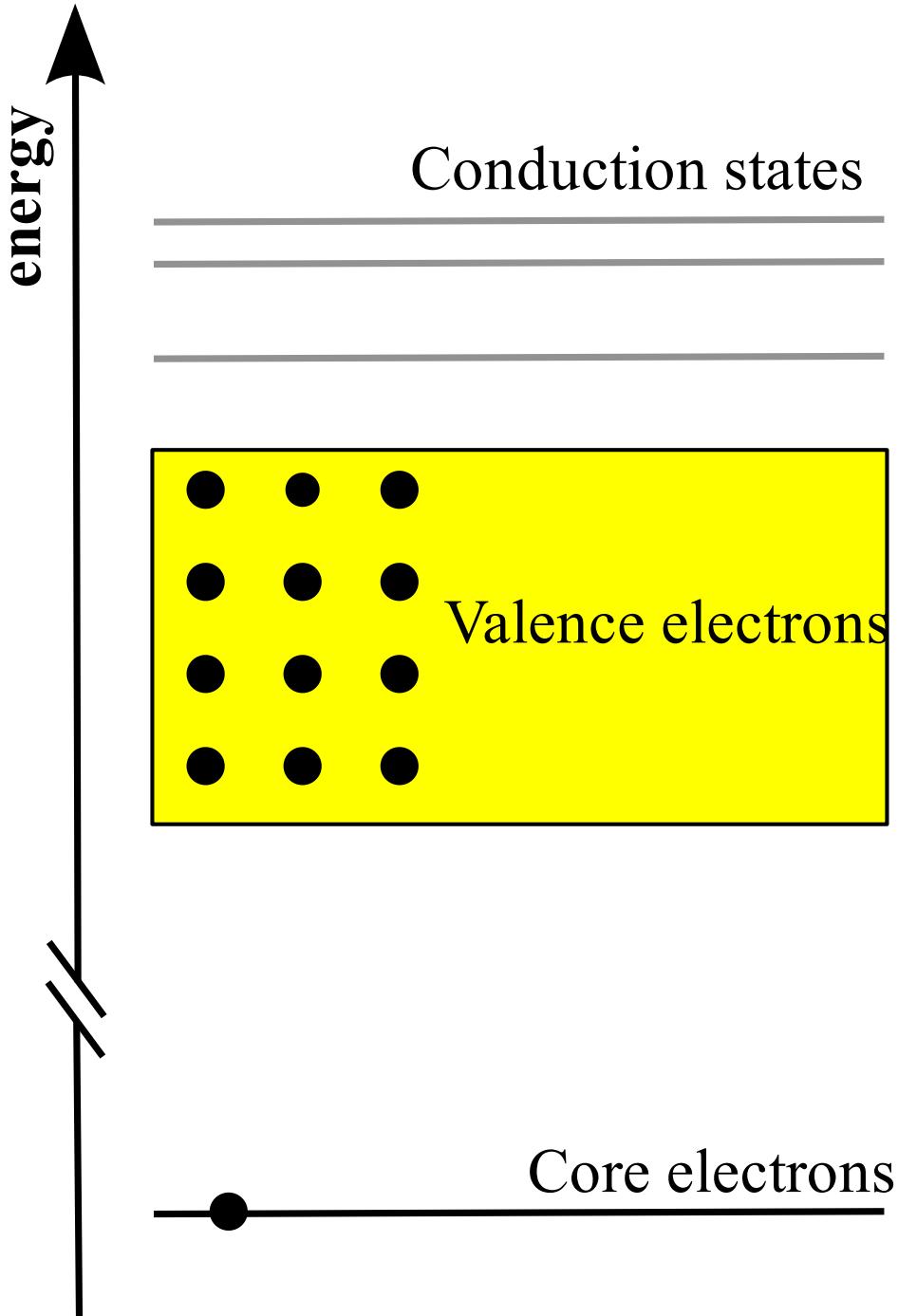
Project Region: Île-de-France

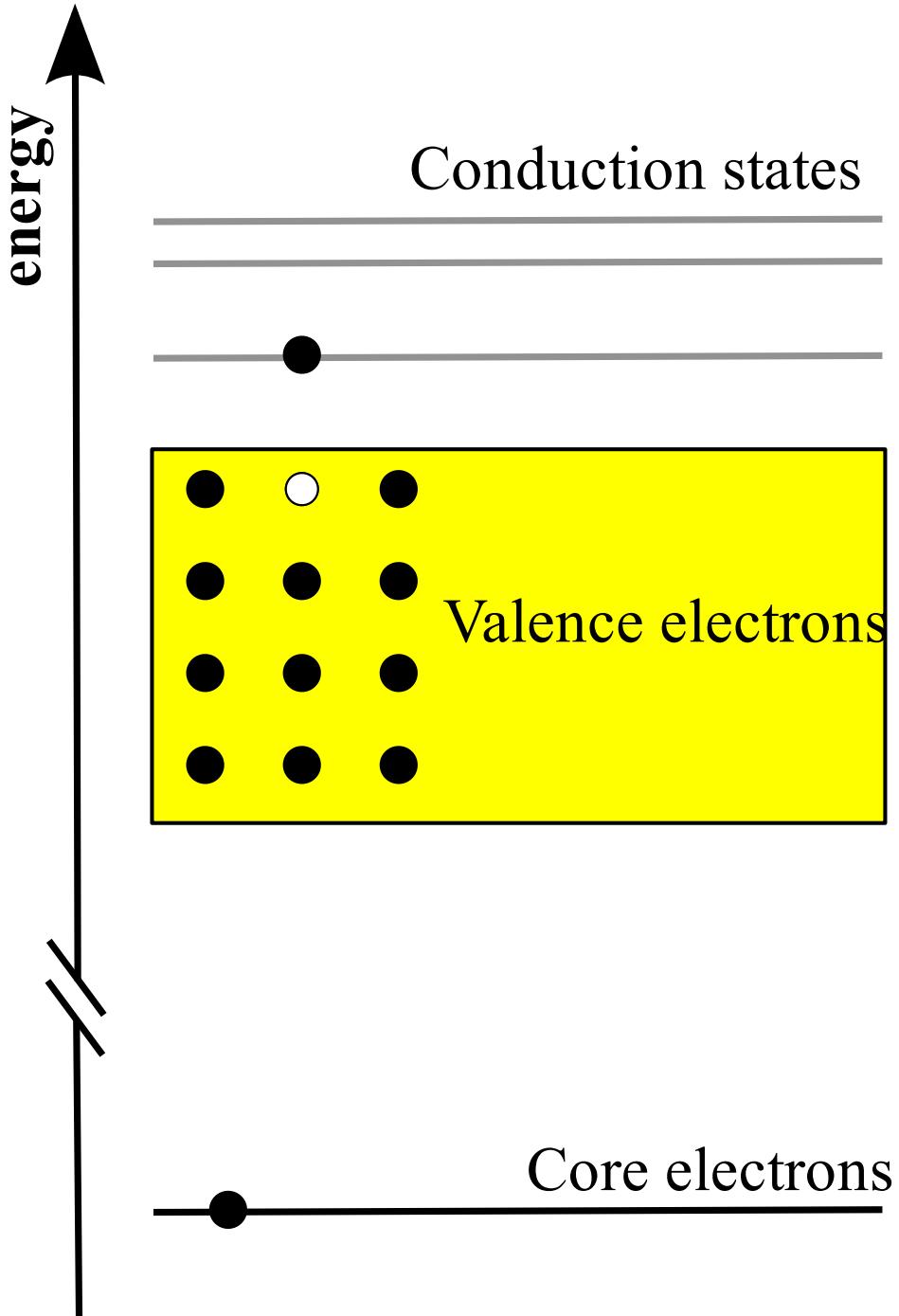
Discipline: 2 - SMI

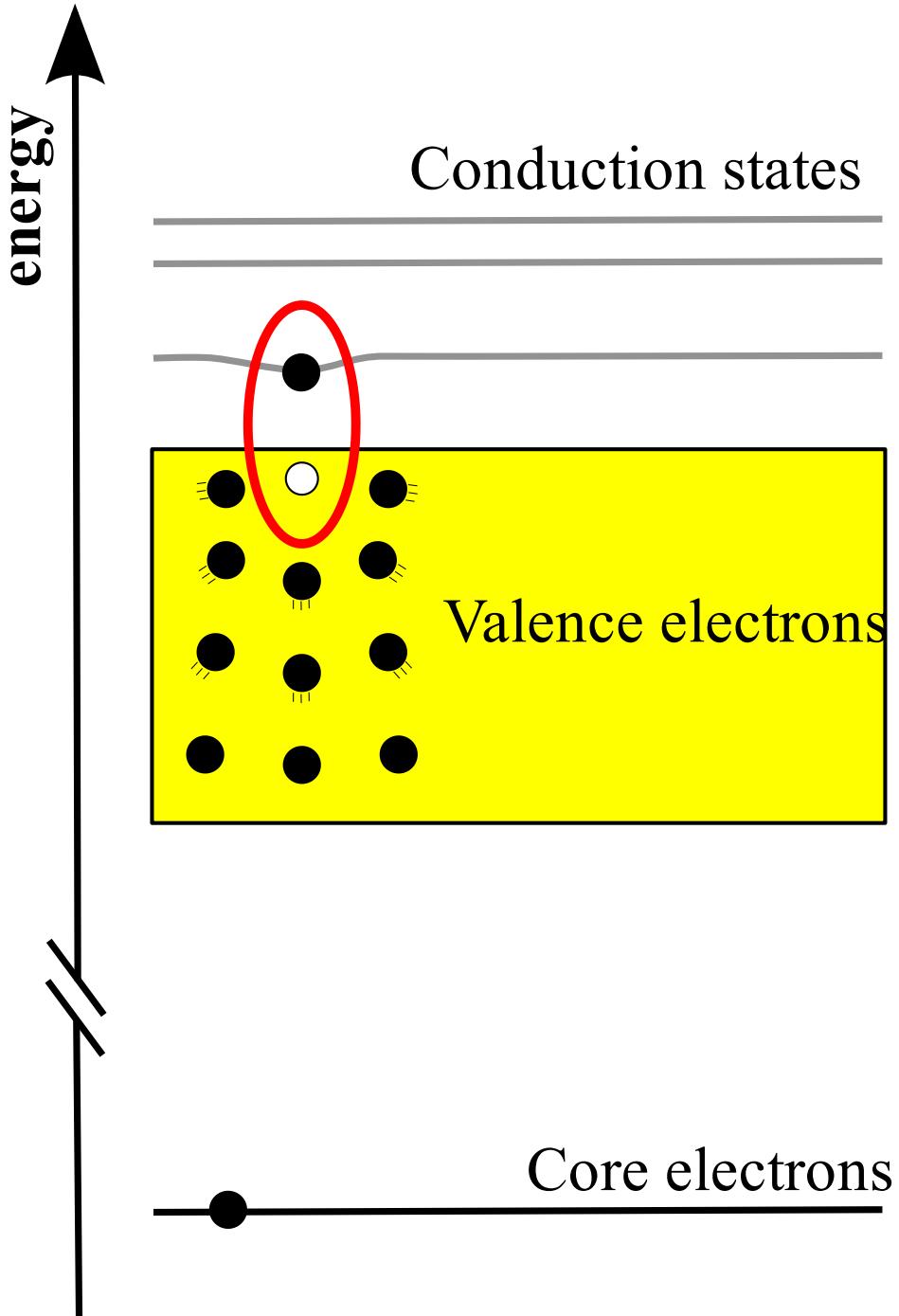
PIA investment: 5,049,988 €

Start date: November 2023

End date: October 2028





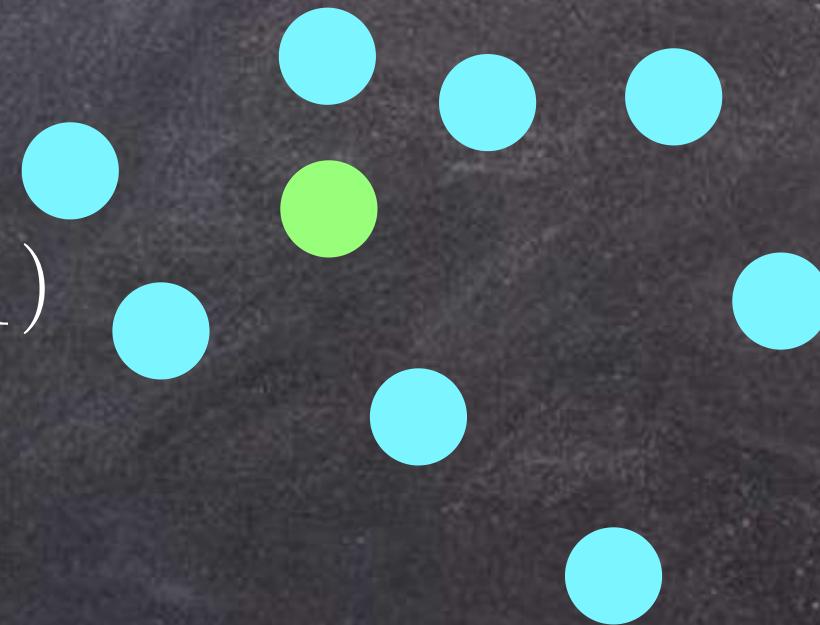


- 
- The background of the slide is a photograph of the Edinburgh skyline at dusk. The sky is a mix of blue and orange hues. In the foreground, there are several historic buildings, including a large castle-like structure on a hill to the left and a tall clock tower with a white face in the center-right. To the right, a tall Gothic-style monument with a statue on top is visible. The city lights are starting to come on, creating a warm glow against the cool tones of the sunset.
- Excitons via Green's functions many-body theory
 - Challenges
 - Results and accuracy

The Green's functions mb formalism

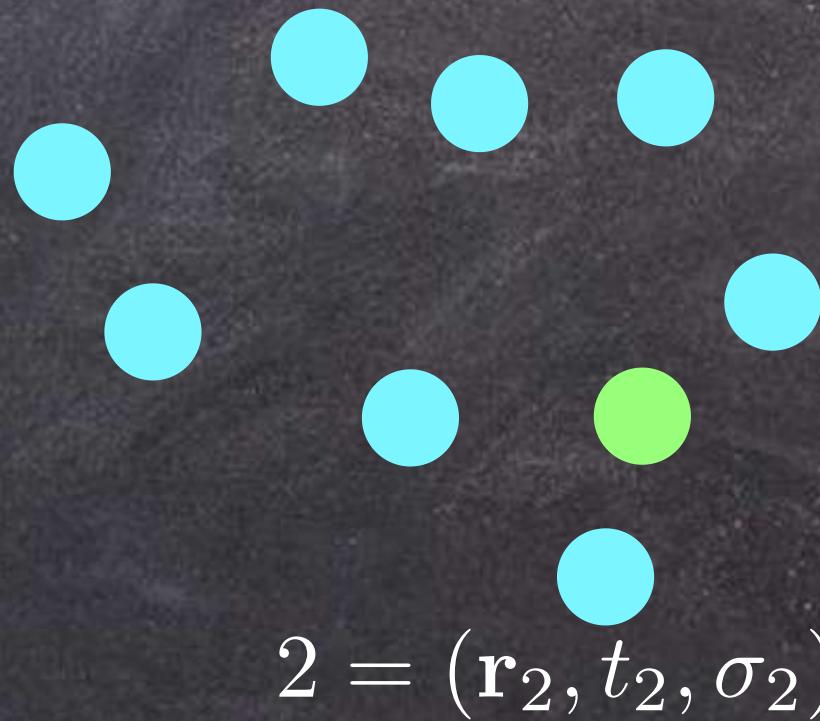
$$G(1, 2)$$

$$1 = (\mathbf{r}_1, t_1, \sigma_1)$$



The Green's functions mb formalism

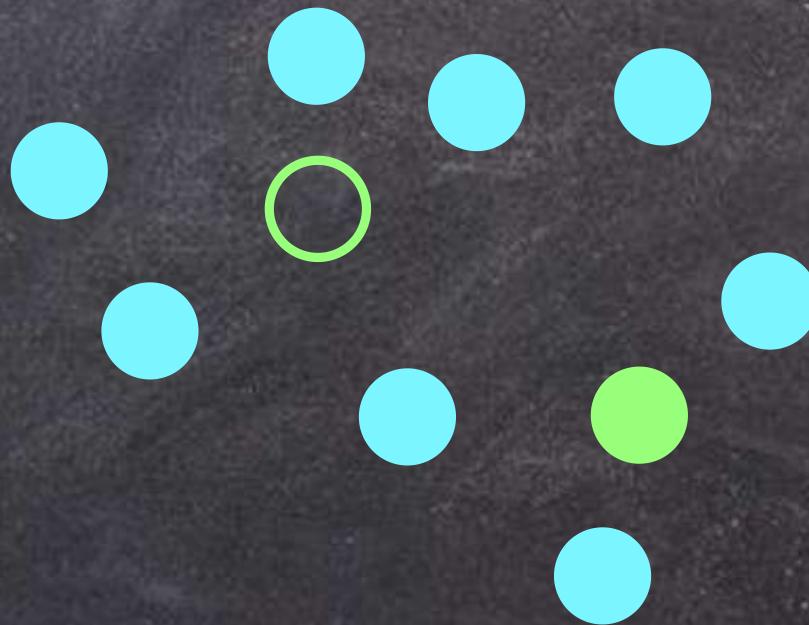
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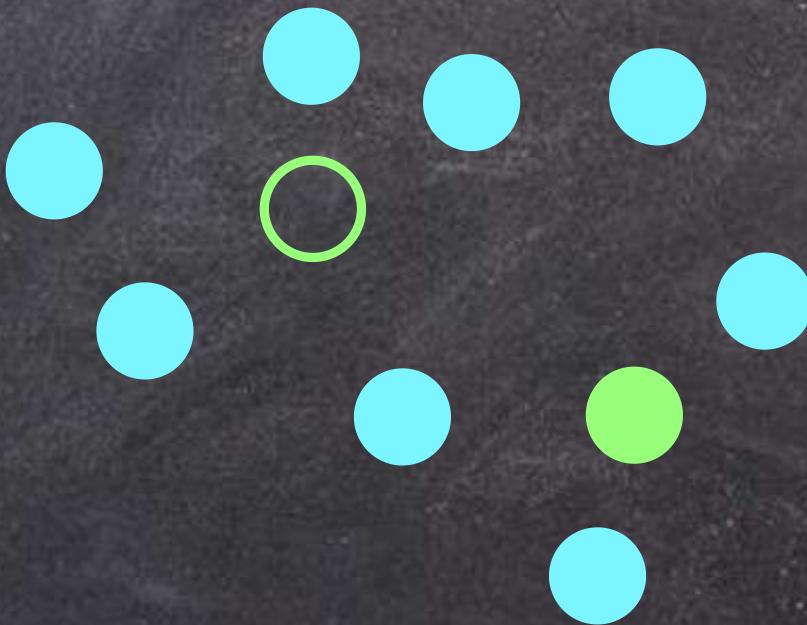
$$G(1, 2)$$

$$G^{(2)}(1, 2, 3, 4)$$



The Green's functions mb formalism

$$G(1, 2)$$



$$G^{(2)}(1, 2, 3, 4) - G(1, 2)G(3, 4) = -iL(1, 2, 3, 4) = \frac{\delta G(1, 3)}{\delta V_{ext}(2, 4)}$$

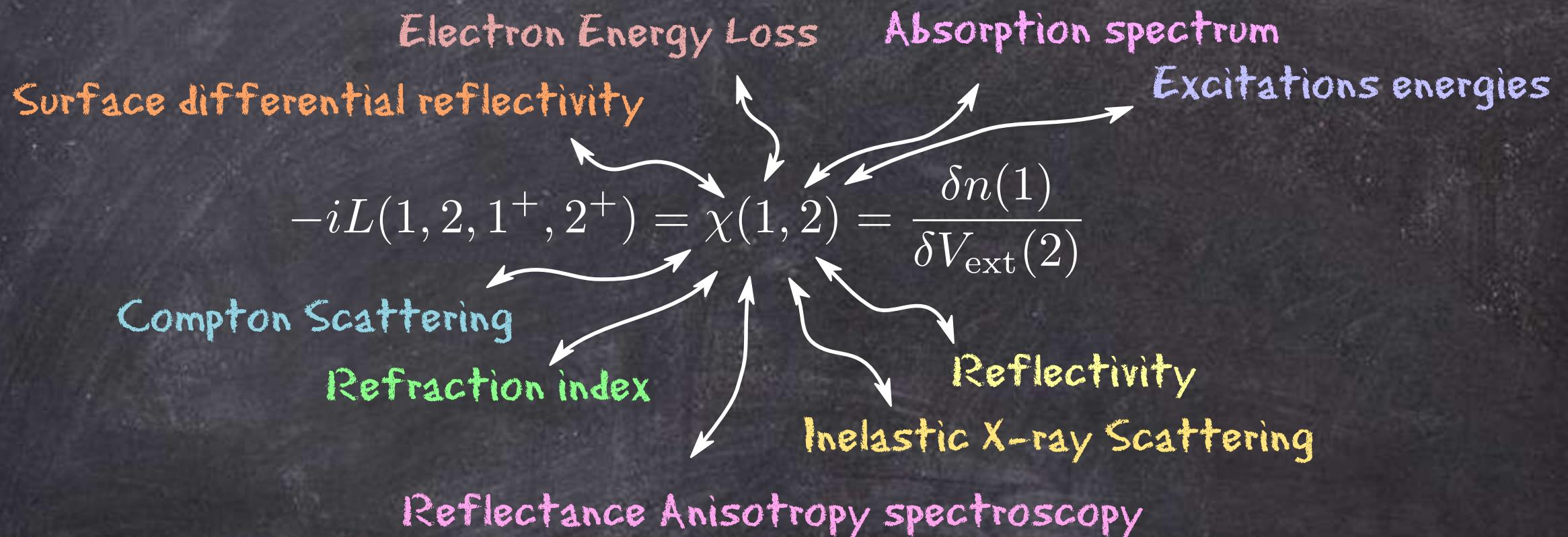
$$-iL(1,2,1^+,2^+)=\chi(1,2)=\frac{\delta n(1)}{\delta V_{\rm ext}(2)}$$

Electron Energy Loss

Excitations energies

$$-iL(1, 2, 1^+, 2^+) = \chi(1, 2) = \frac{\delta n(1)}{\delta V_{\text{ext}}(2)}$$

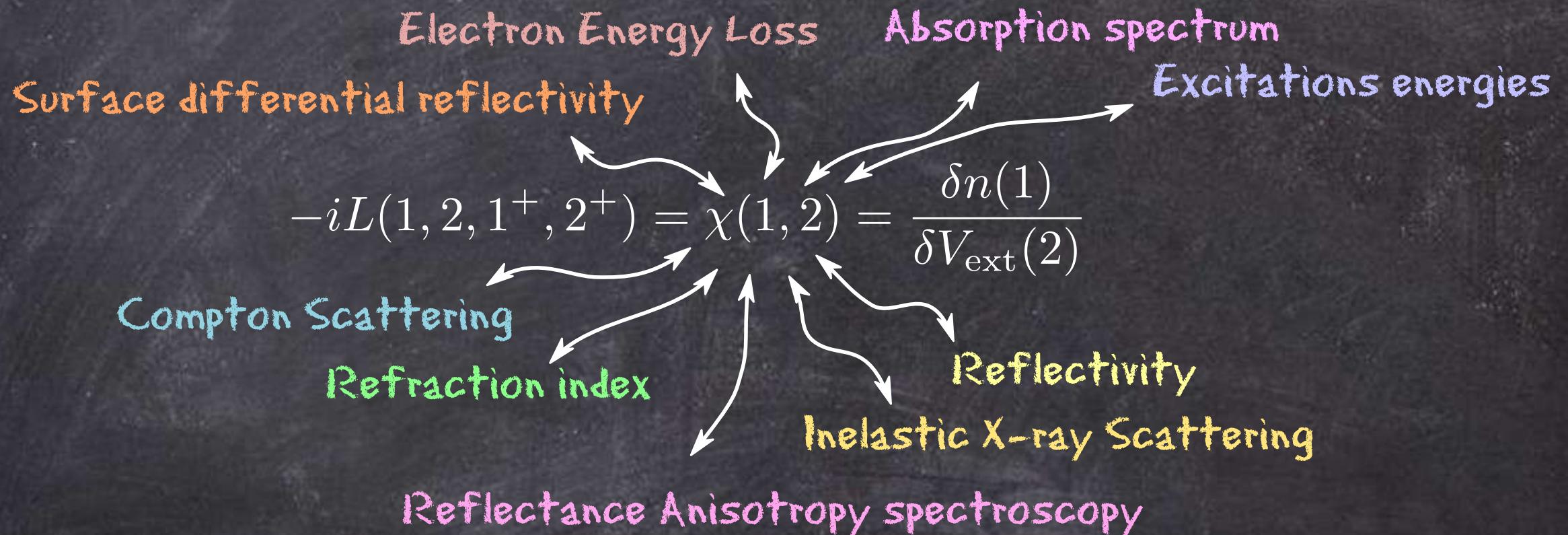
Inelastic X-ray Scattering



The Bethe-Salpeter Equation

$$L(1, 2, 3, 4) = L^0(1, 2, 3, 4) + L^0(1, 2, 5, 6) \left[v(5, 7) \delta(5, 6) \delta(7, 8) + i \frac{\delta \Sigma(5, 6)}{\delta G(7, 8)} \right] L(7, 8, 3, 4)$$

$$L^0(1, 2, 3, 4) = G(1, 2)G(3, 4)$$



- Excitons via Green's functions many-body theory

- Challenges

- Results and accuracy

1st challenge :: G

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starting G
LDA, GGA, HF

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$$G^0(r, r', \omega) = \sum_s \frac{\phi_s^*(r)\phi_s(r')}{\omega - \epsilon_s \pm i\eta}$$

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$$\Sigma \approx iGW$$

$$W^{RPA} = \frac{v}{1 - vGG}$$

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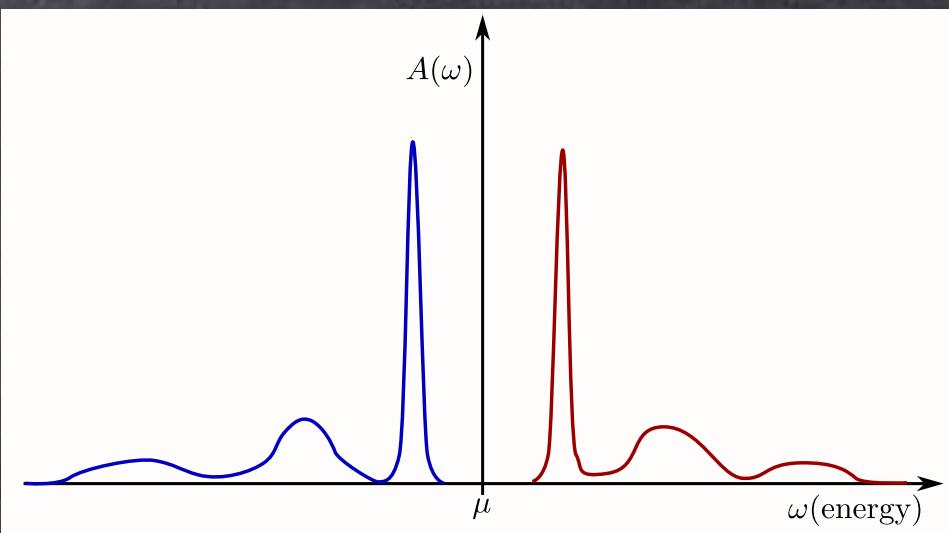
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$$A(\omega) = \frac{1}{\pi} \operatorname{Im} G$$



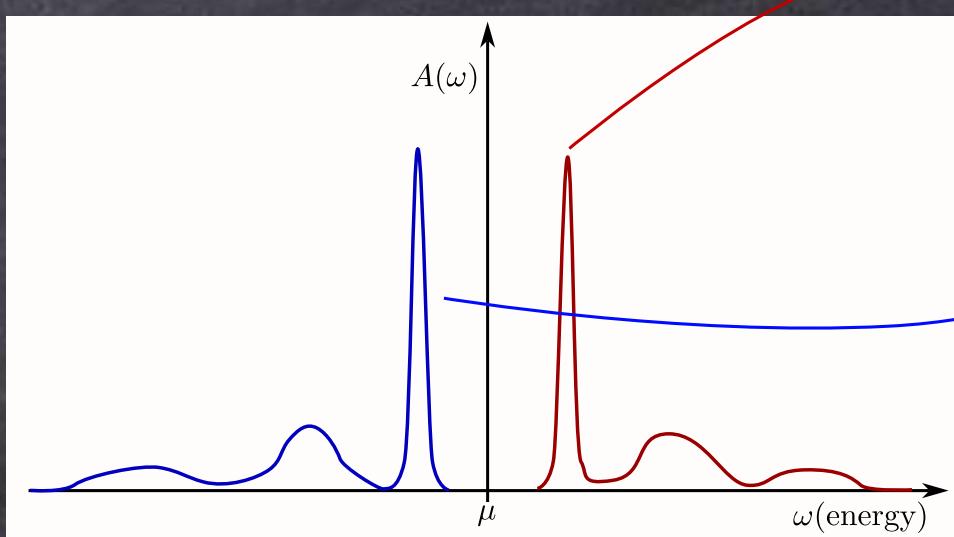
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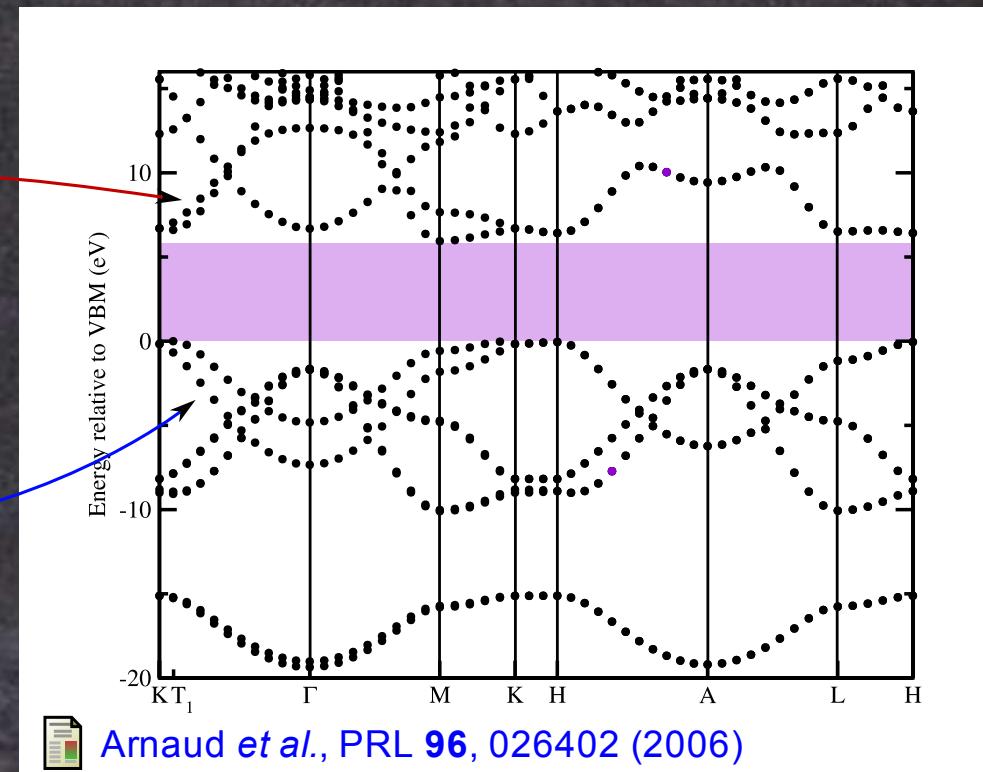
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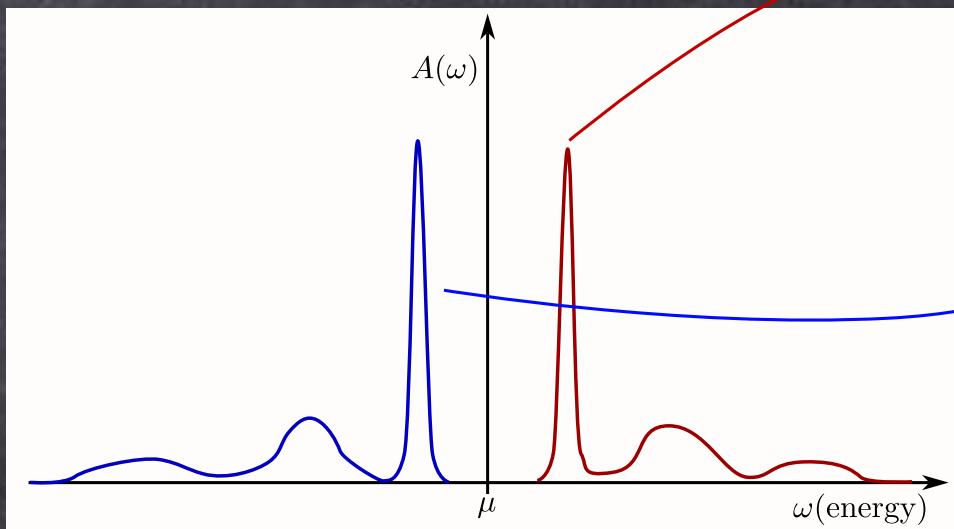
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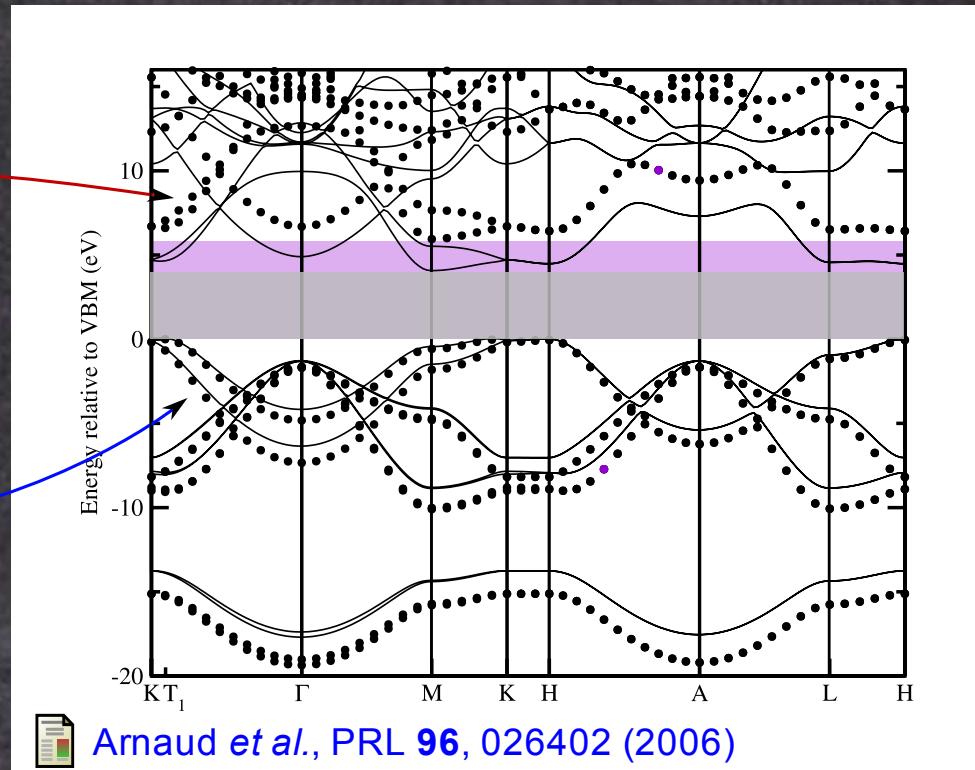
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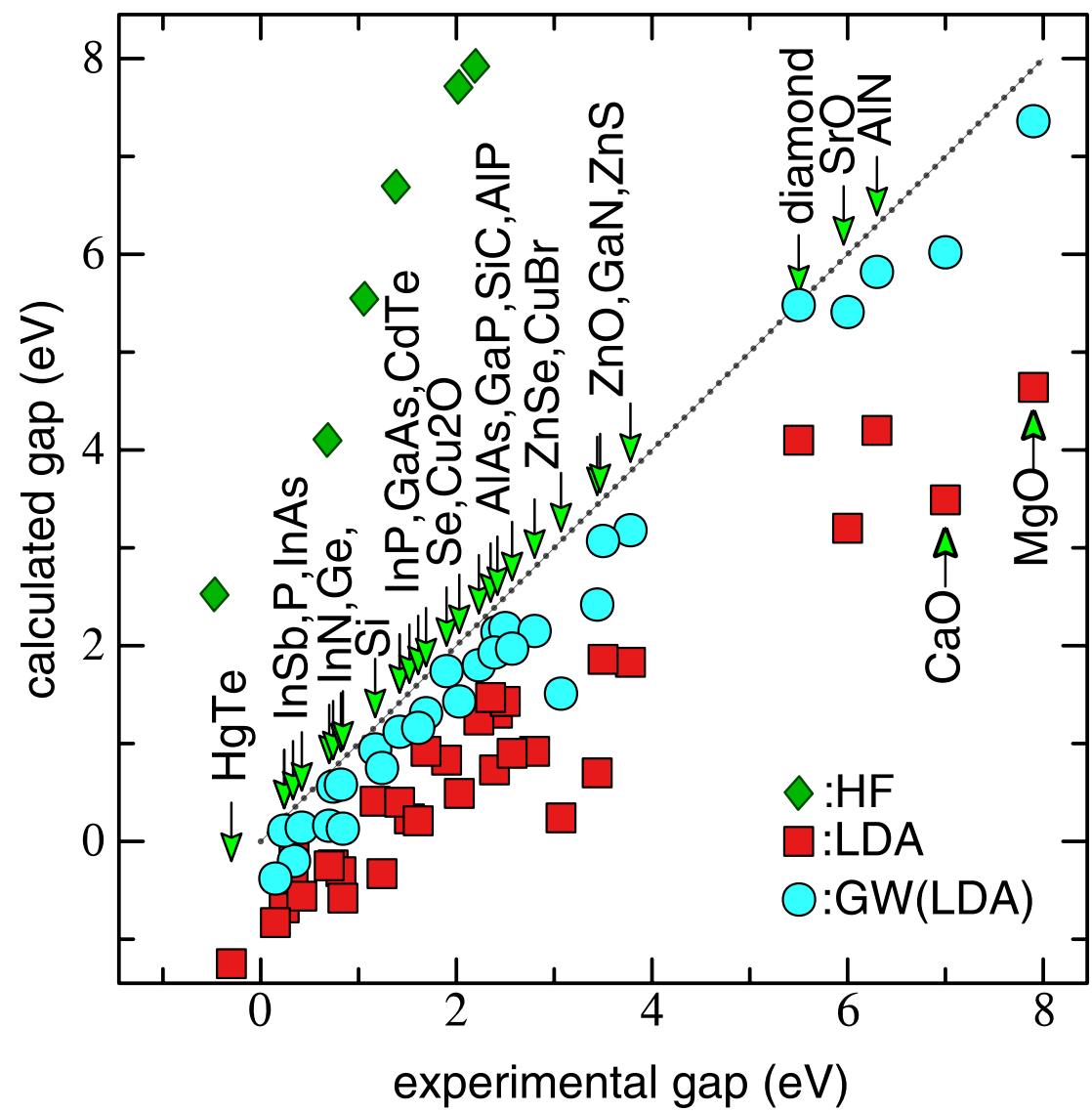


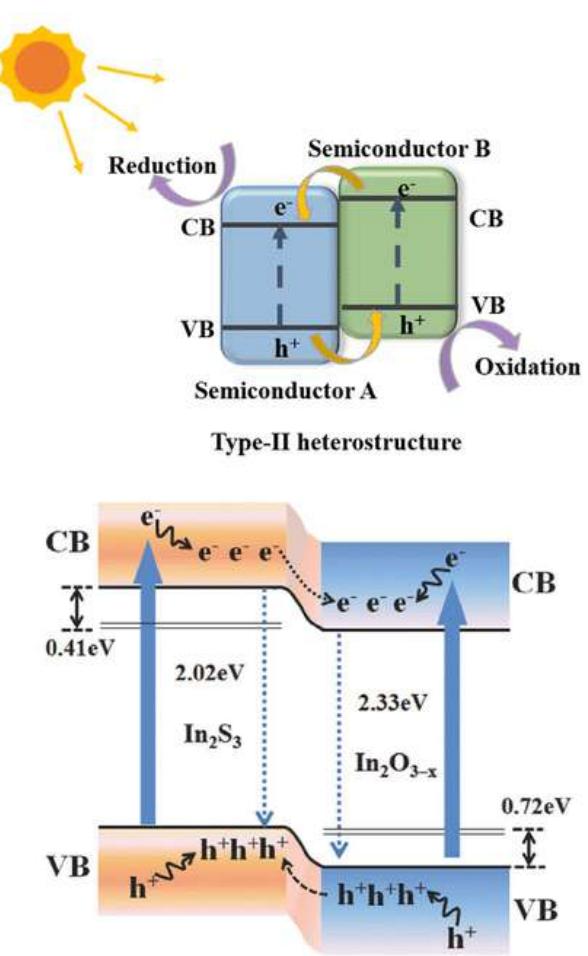
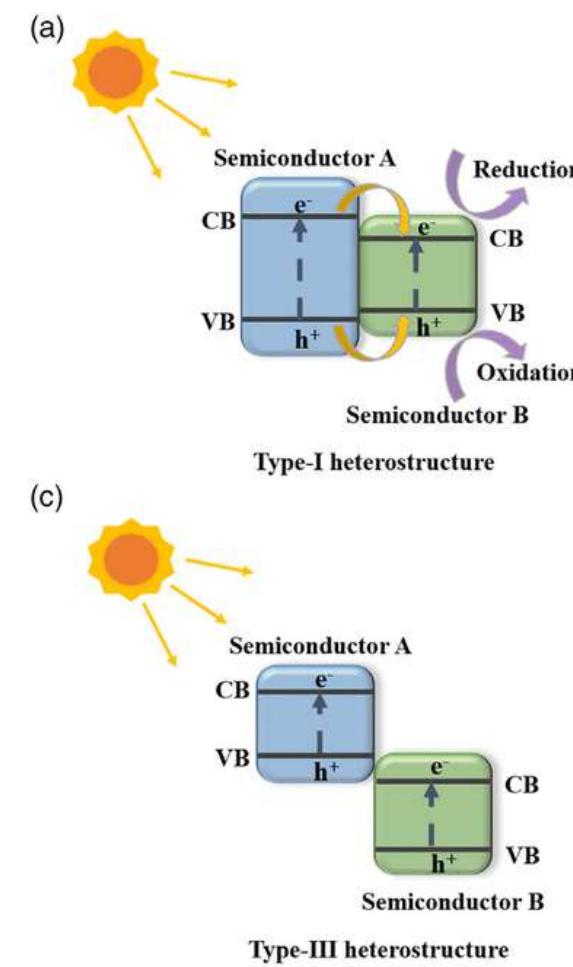
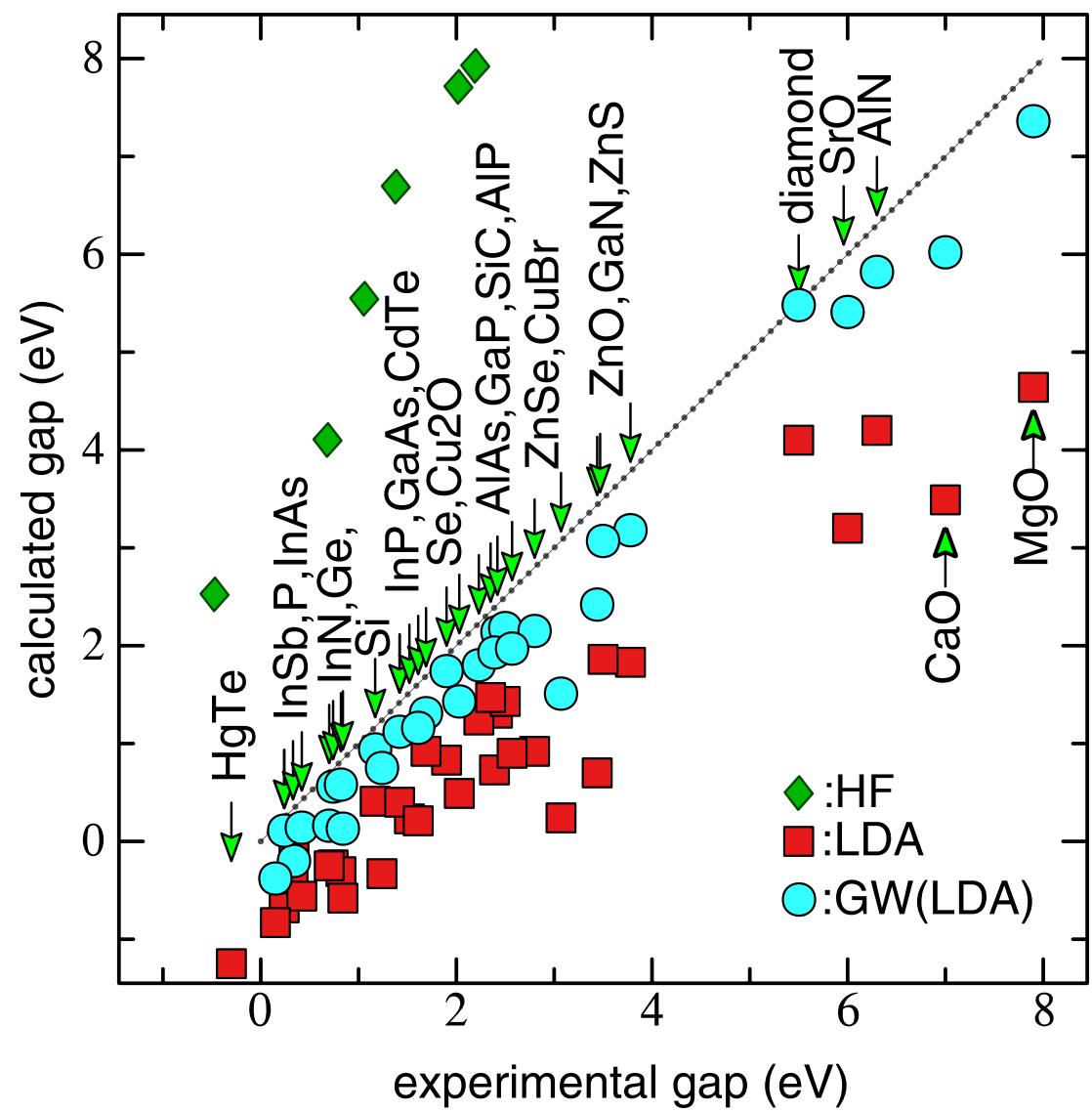
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$$W^{RPA} = \frac{v}{1 - vGG}$$



Arnaud et al., PRL 96, 026402 (2006)





2nd challenge :: solving the BSE

$$L(1, 2, 3, 4) = L^0(1, 2, 3, 4) + L^0(1, 2, 5, 6) \left[v(5, 7) \delta(5, 6) \delta(7, 8) + i \frac{\delta \Sigma(5, 6)}{\delta G(7, 8)} \right] L(7, 8, 3, 4)$$

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$$L(1, 2, 3, 4) = GG + GG [v - W] L$$

$$L(r_1, r_2, r_3, r_4, \omega) \Rightarrow L_{vc}^{v'c'}(\omega)$$

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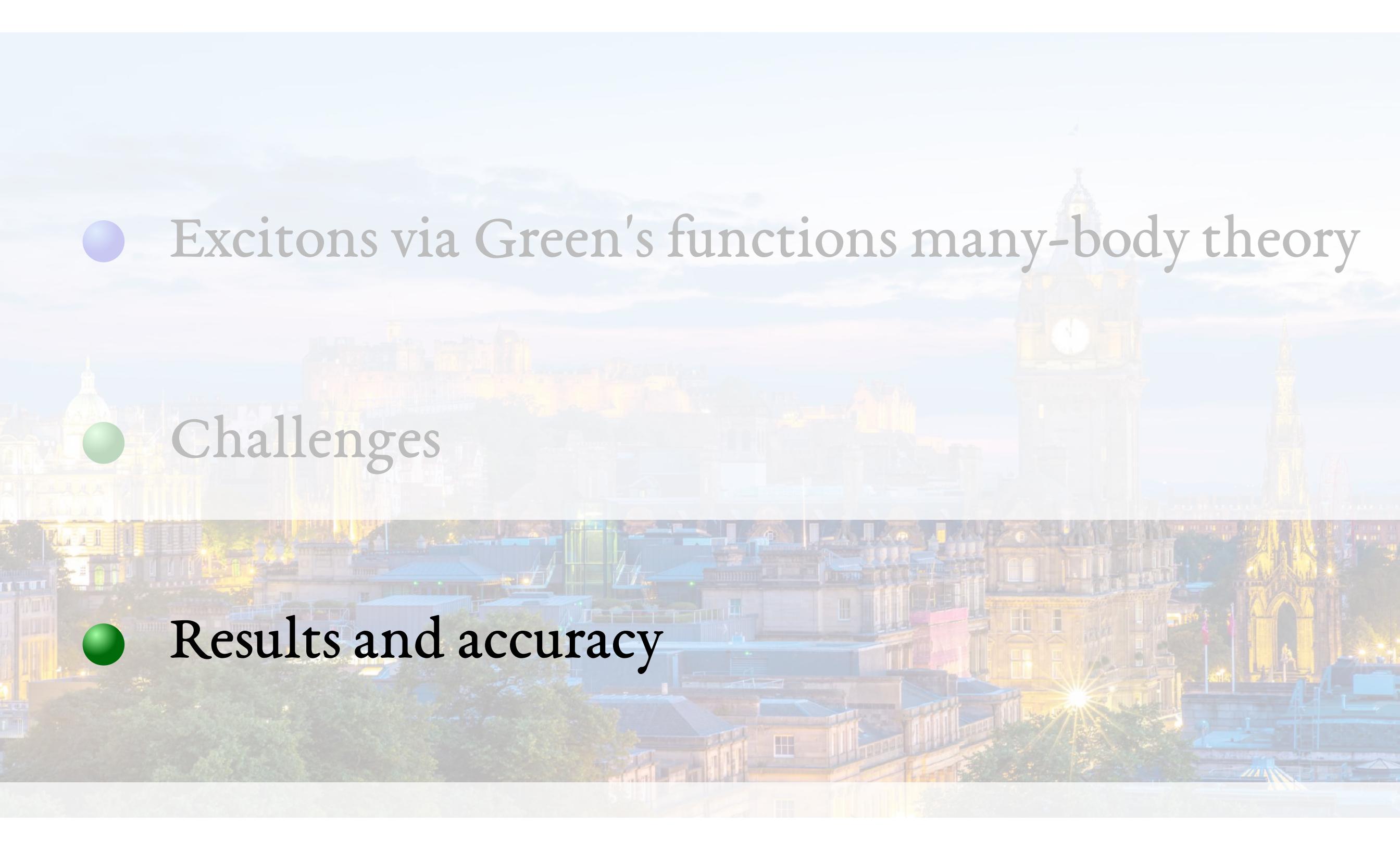
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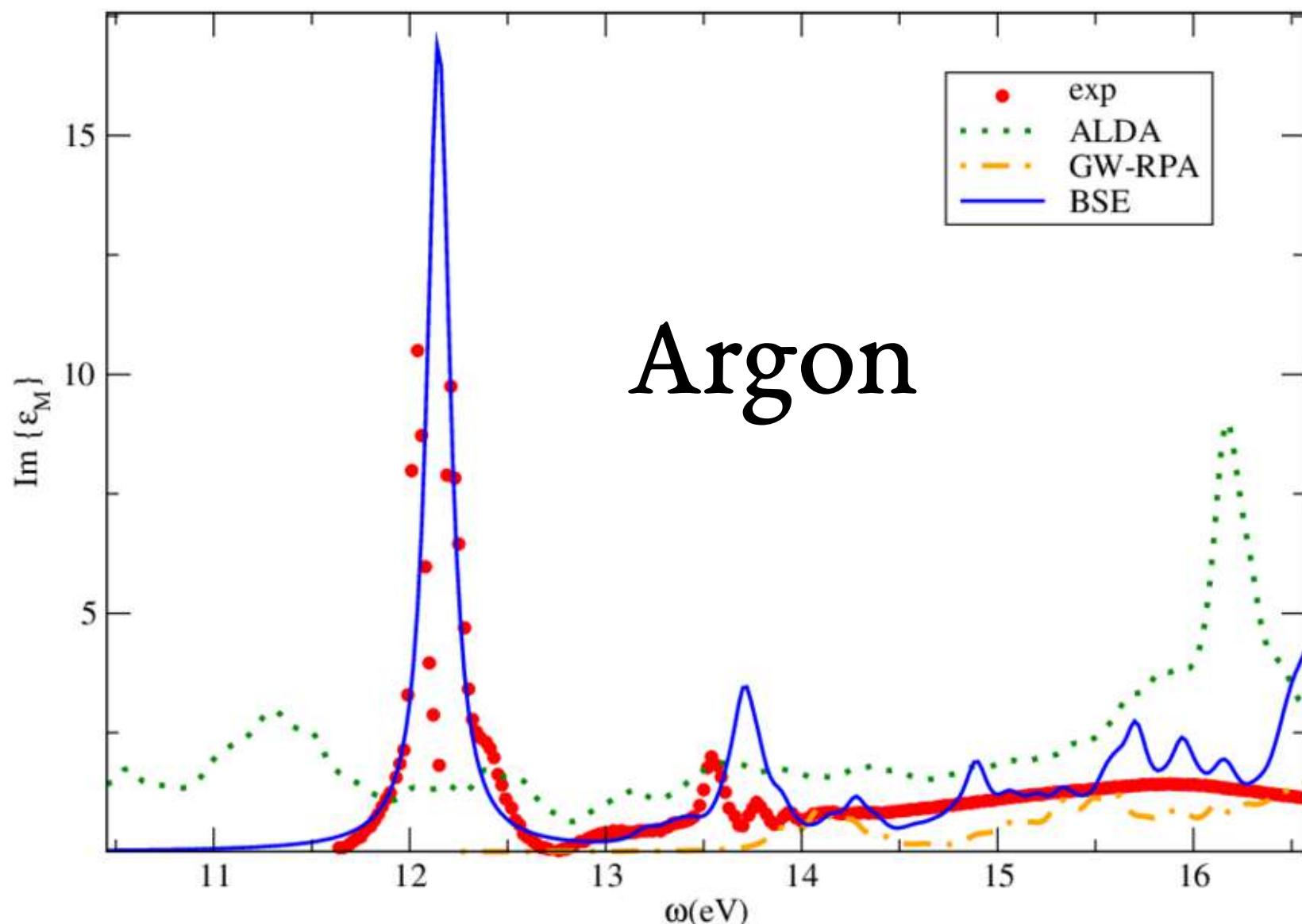
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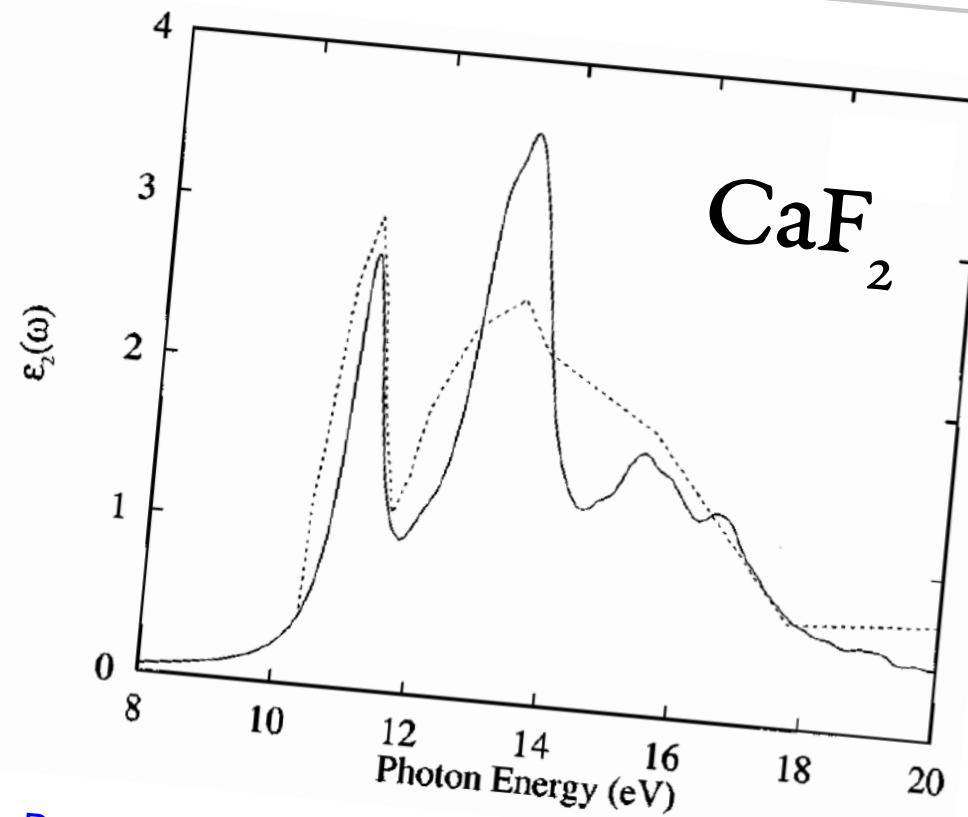
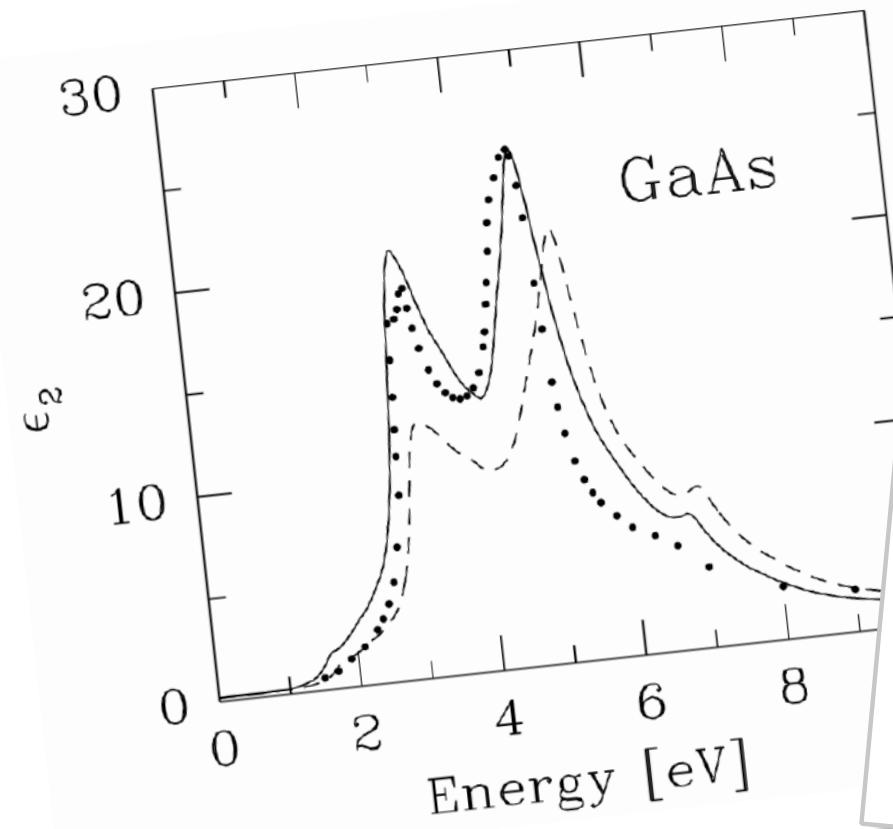
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scaling N_{at}^{4-6}

- 
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Phys. Rev. B 76 161103 (2007)

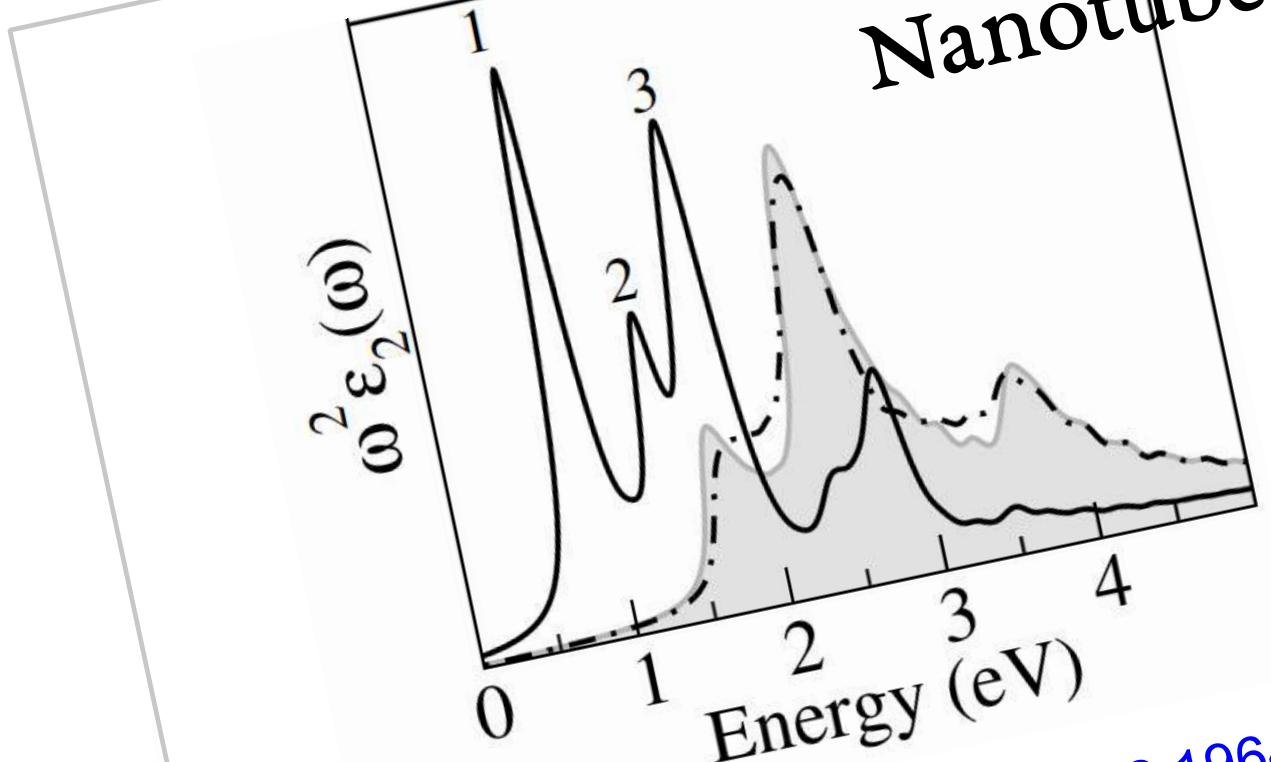


Benedict and Shirley Phys. Rev. B **59**, 5441 (1999)

Rohlfing and Louie Phys. Rev. Lett. **81**, 2312 (1998)



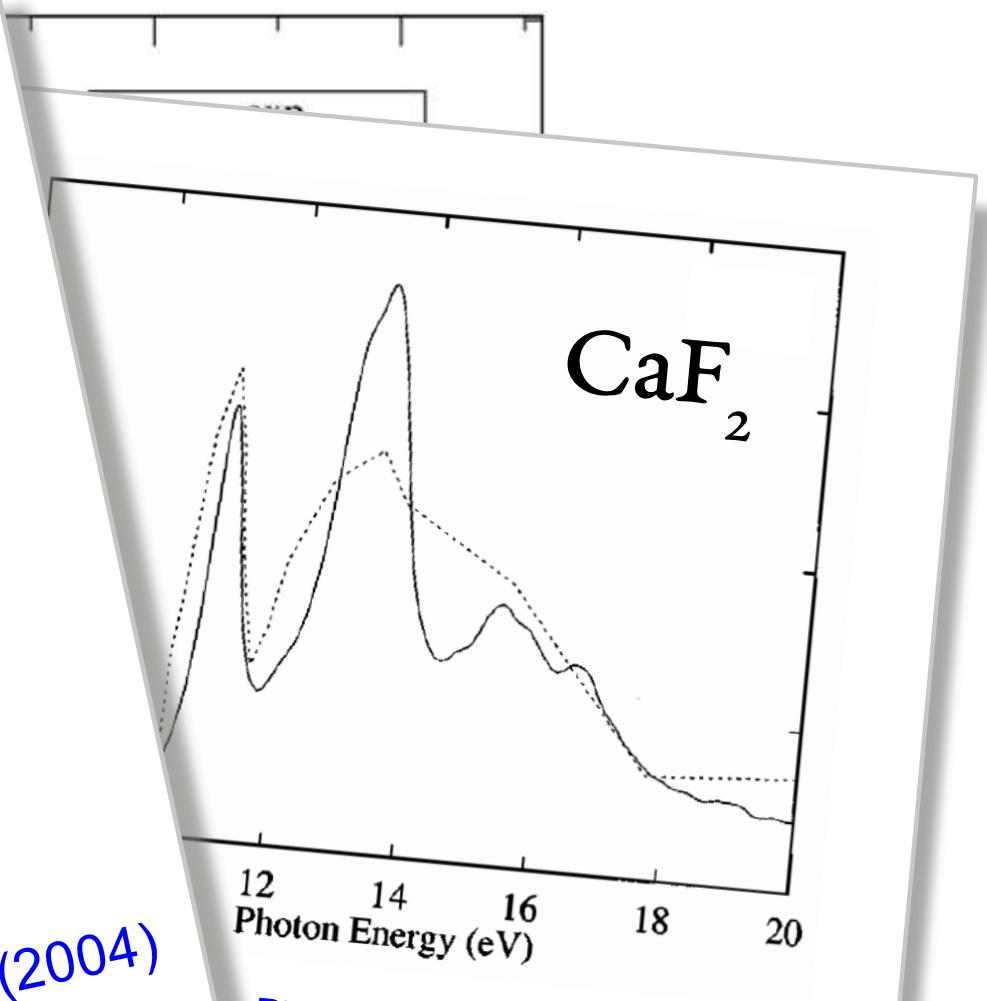
Phys. Rev. B **76** 161103 (2007)



Rohli

Chang et al., Phys. Rev. Lett. 92 196401 (2004)

Phys. Rev. B 76 161103 (2007)



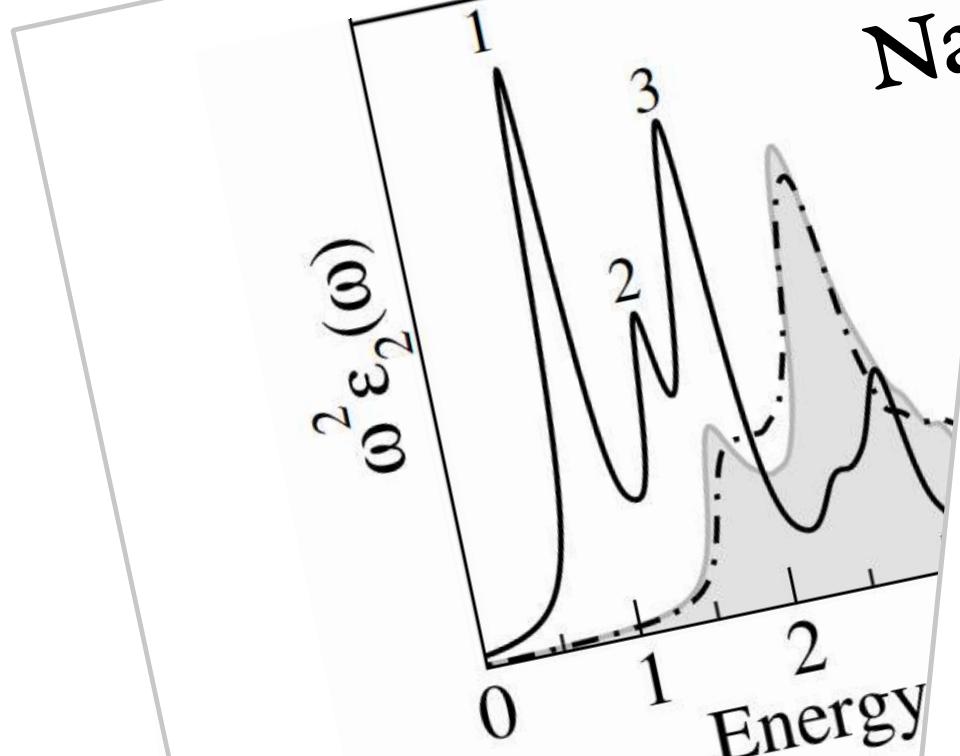
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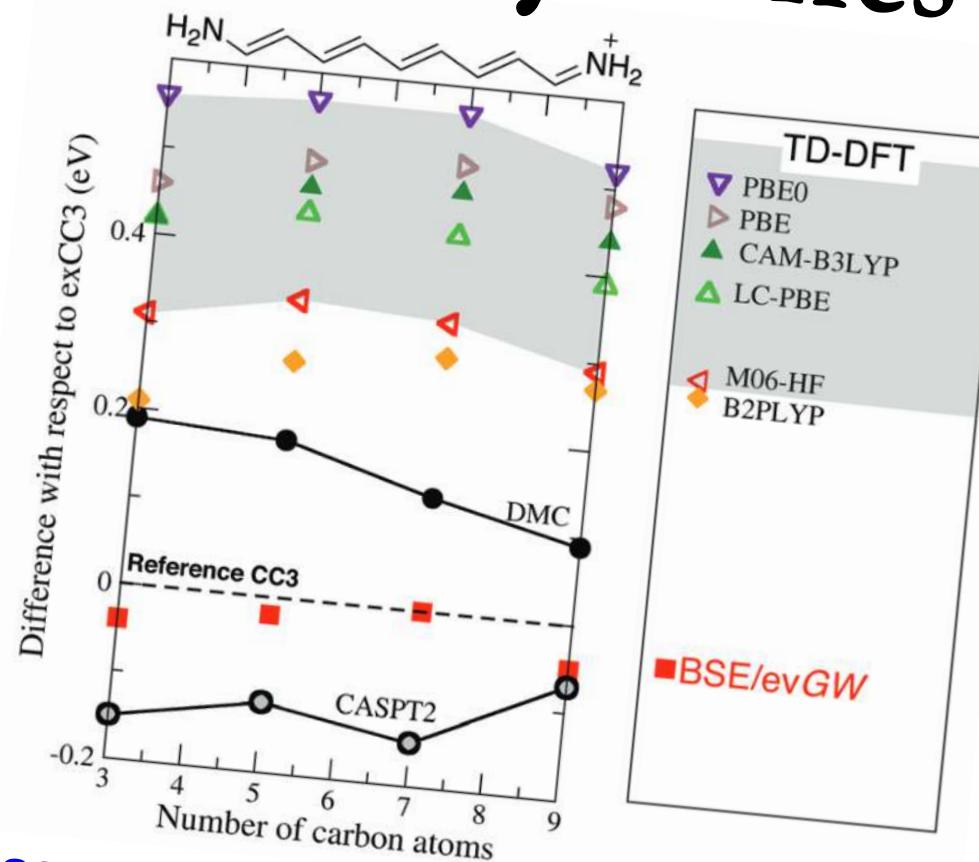


Na⁺

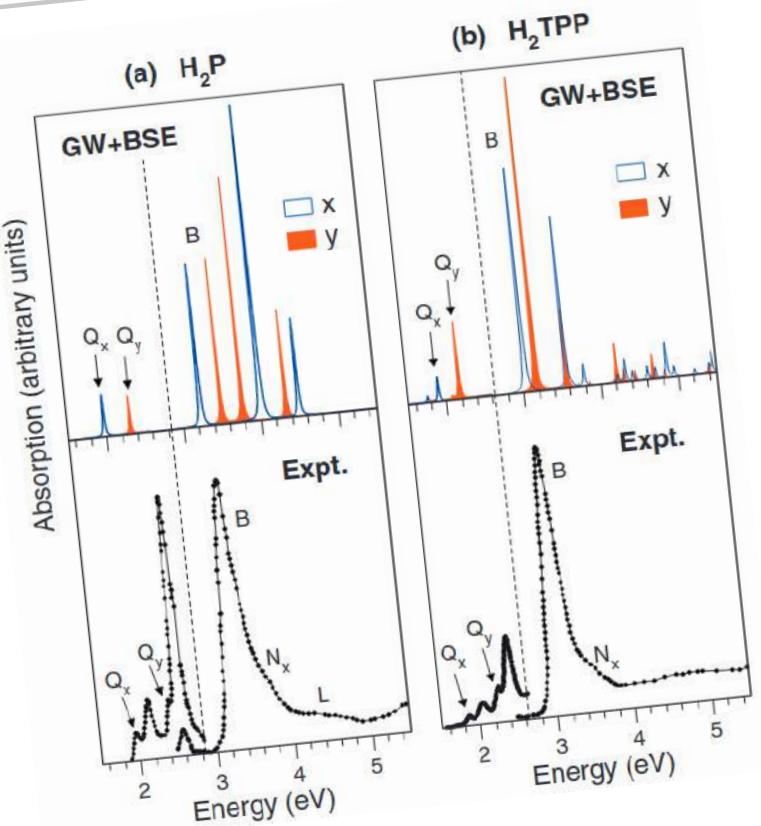


Blase et al. Chem. Soc. Rev. 47, 1022 (2018)

streptocyanines



Porphyrins

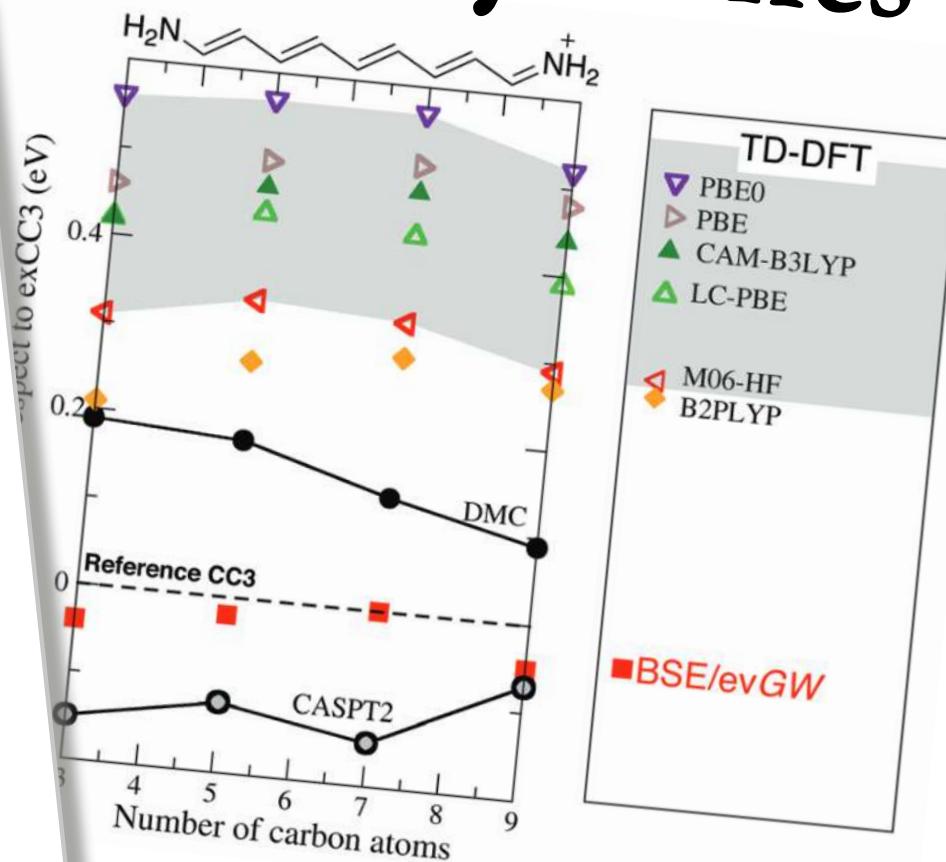


Palummo et al., J. Chem. Phys. 131 084102 (2009)



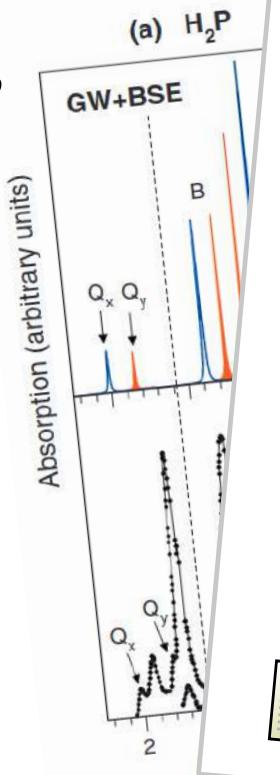
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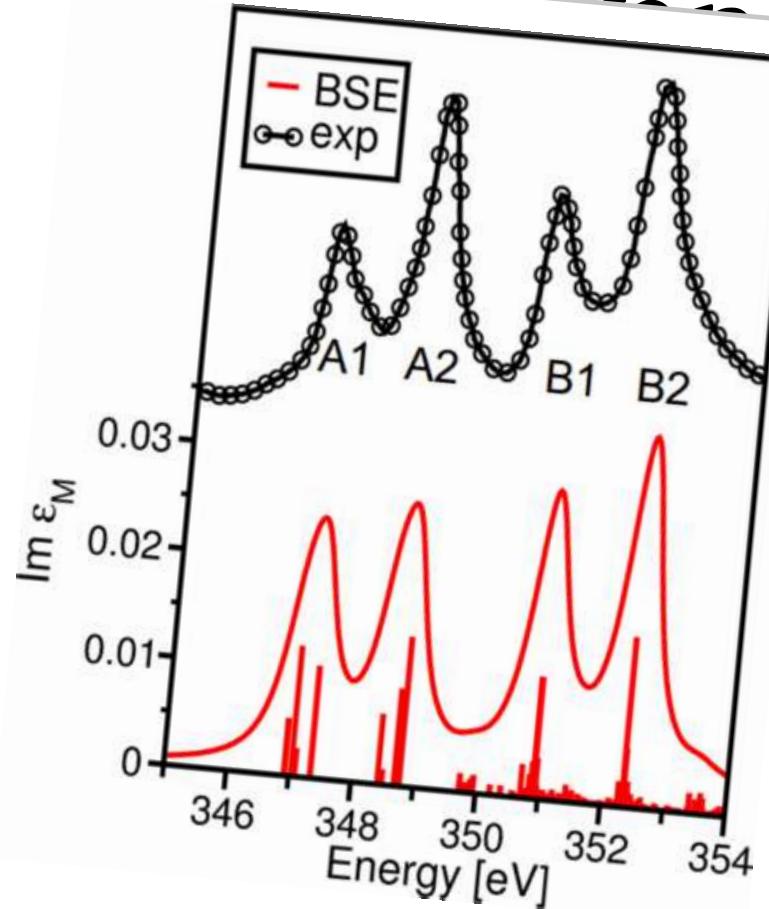


et al. Chem. Soc. Rev. 47, 1022 (2018)

Porphyrins



CaO Ca L-edge



es



Vorwerk *et al.*, Phys. Rev. B 95, 155121 (2017)



Palummo *et al.*, J. Chem. Phys. 131 084102



Phys. Rev. B 76 161103 (2007)

, 1022 (2018)

BSE :: accurate for absorption spectra (and excitation energies)



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- it captures the physics of the electron-hole interaction

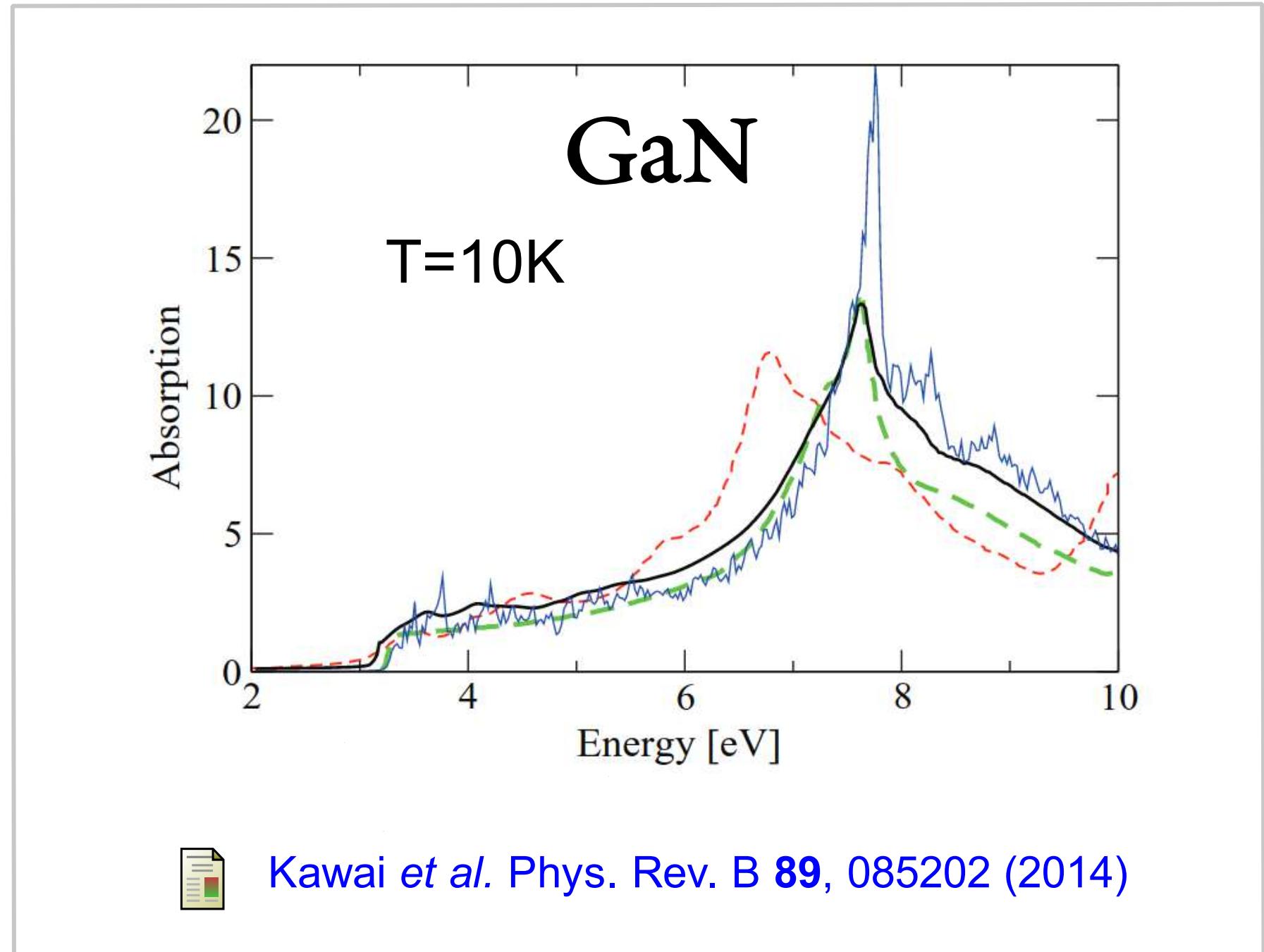


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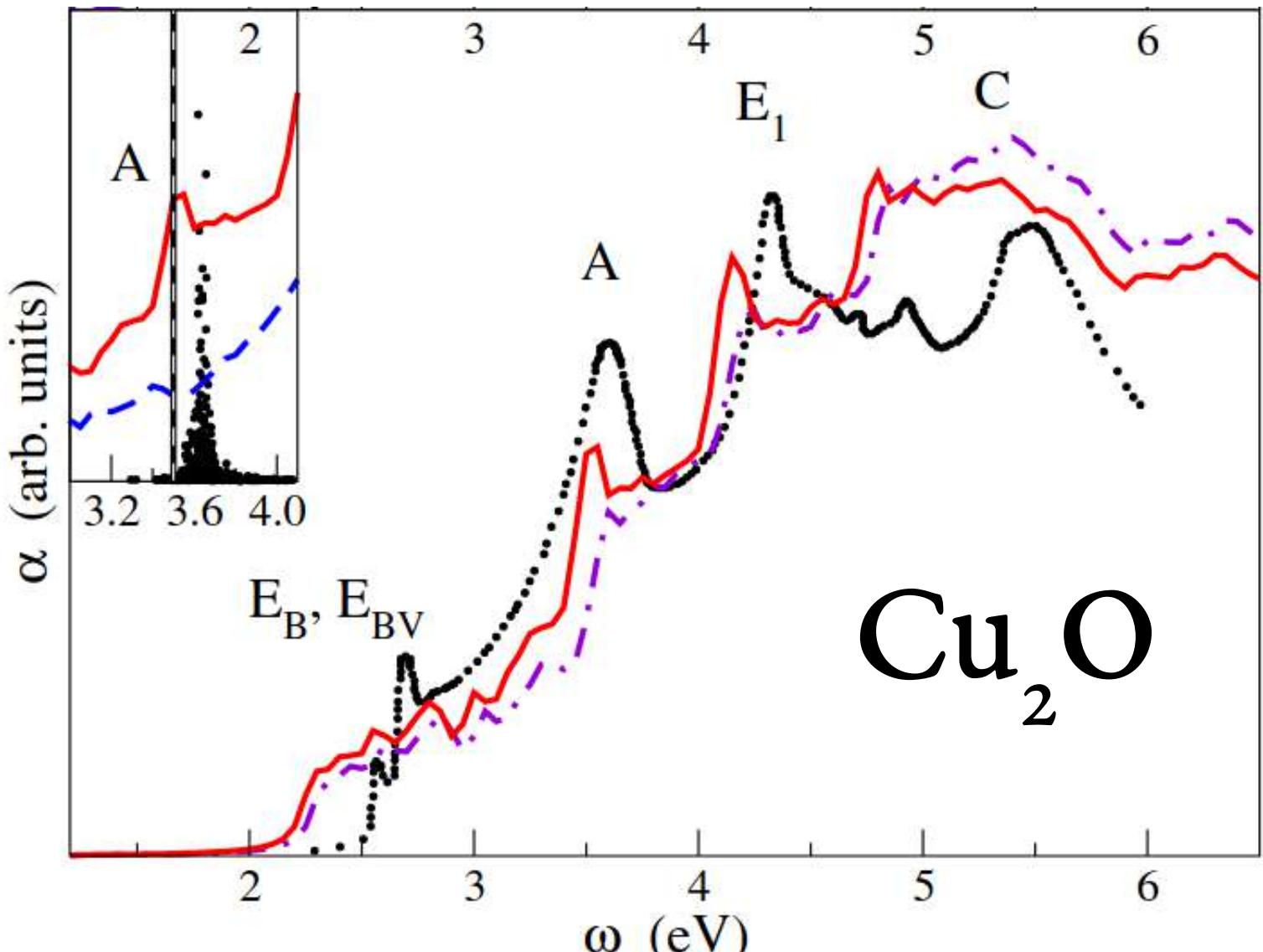
- it captures the physics of the electron-hole interaction
- it can (automatically) profit from extensions



Temperature electron-phonon



SC-GW band-structure



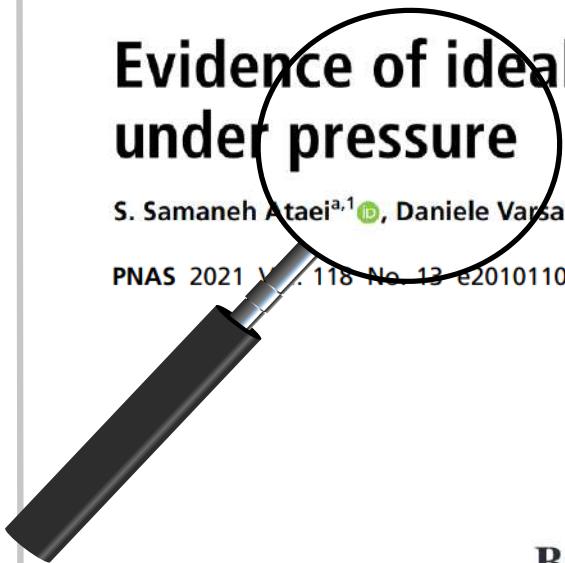
Bruneval *et al.* Phys. Rev. Lett. **97**, 267601 (2006)

Evidence of ideal excitonic insulator in bulk MoS₂ under pressure

S. Samaneh Ataei^{a,1}, Daniele Vassano^{a,1}, Elisa Molinari^{a,b}, and Massimo Rontani^{a,2}

PNAS 2021, Vol. 118, No. 13, e2010110118

<https://doi.org/10.1073/pnas.2010110118>



PHYSICAL REVIEW B, VOLUME 65, 155332

Bethe-Salpeter equation for magnetoexcitons in quantum wells

Z. G. Koinov*

Department of Physics & Astronomy, University of Texas at San Antonio, San Antonio, Texas 78249

(Received 10 December 2001; published 11 April 2002)

PRL 116, 196804 (2016)

PHYSICAL REVIEW LETTERS

week ending
13 MAY 2016

Three-particle correlation from a Many-Body Perspective: Trions in a Carbon Nanotube

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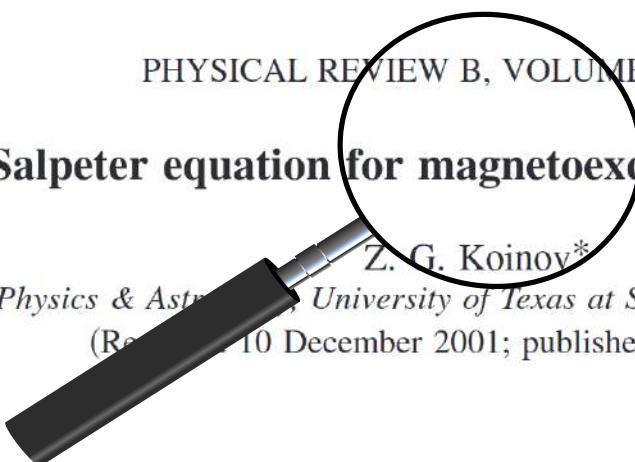
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Bethe-Salpeter Equation - finite momentum transfer

$$S(\mathbf{q}, \omega) \propto \chi_M(\mathbf{q}, \omega) = \sum_{\lambda} \frac{\left| \sum_{vc} A_{\lambda}^{vc, \mathbf{q}} \langle c | e^{i\mathbf{q} \cdot \mathbf{r}} | v \rangle \right|^2}{\omega - E_{\lambda}(\mathbf{q}) + i\eta}$$



Caliebe *et al.* Phys. Rev. Lett. **84**, 3907 (2000)



Soininen and Shirley, Phys. Rev. B **61**, 16423 (2000)



Vinson *et al.* Phys. Rev. B **83**, 115106 (2011)



Gatti and Sottile, Phys. Rev. B **98**, 155113 (2013)

BSE :: accurate for absorption spectra (and excitation energies)

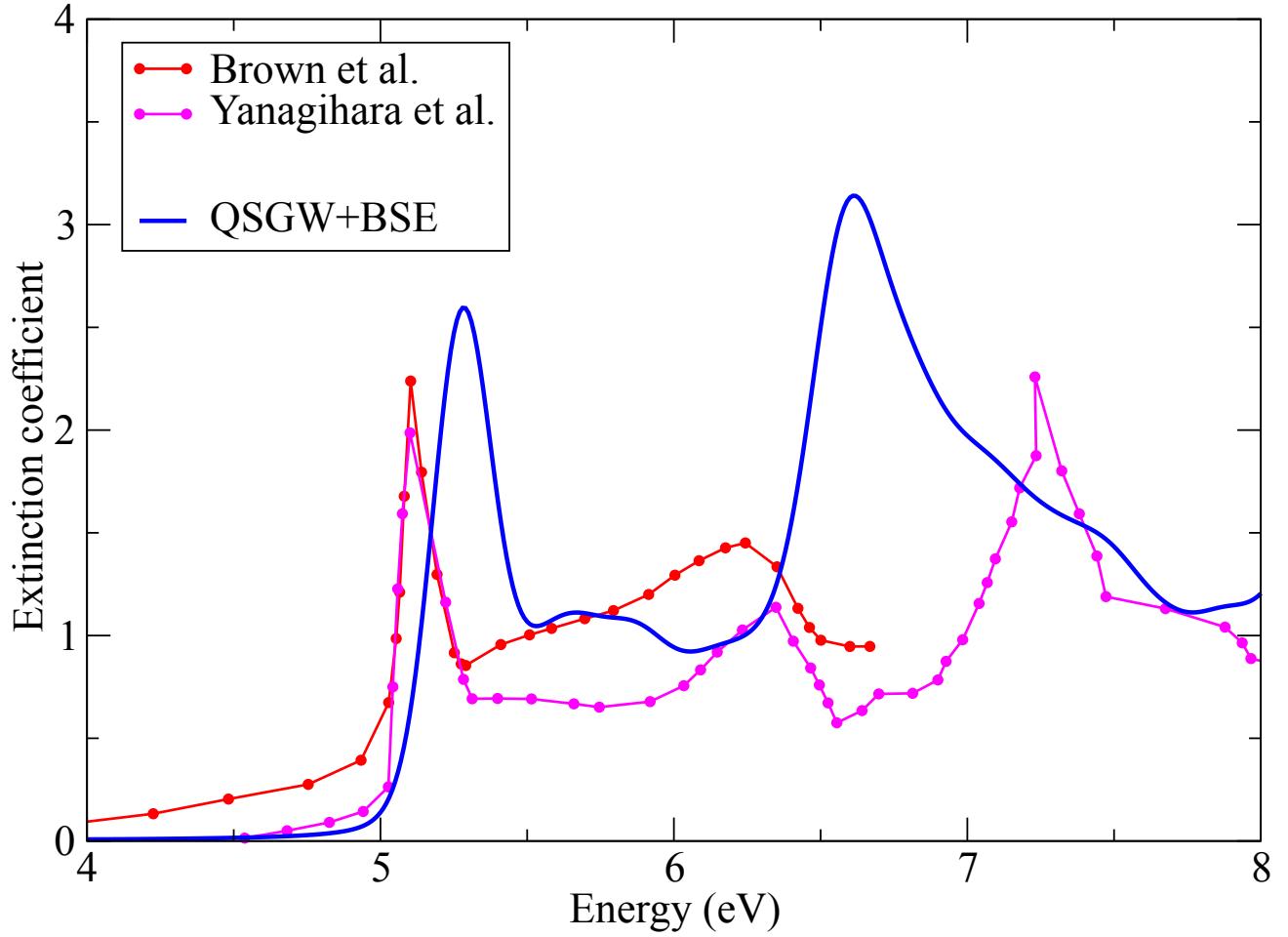
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BSE :: accurate for absorption spectra (and excitation energies)

- it captures the physics of the electron-hole interaction
- it can (automatically) profit from extensions
- *ab initio* → predictions
- analysis tools (why? how? who is responsible?)





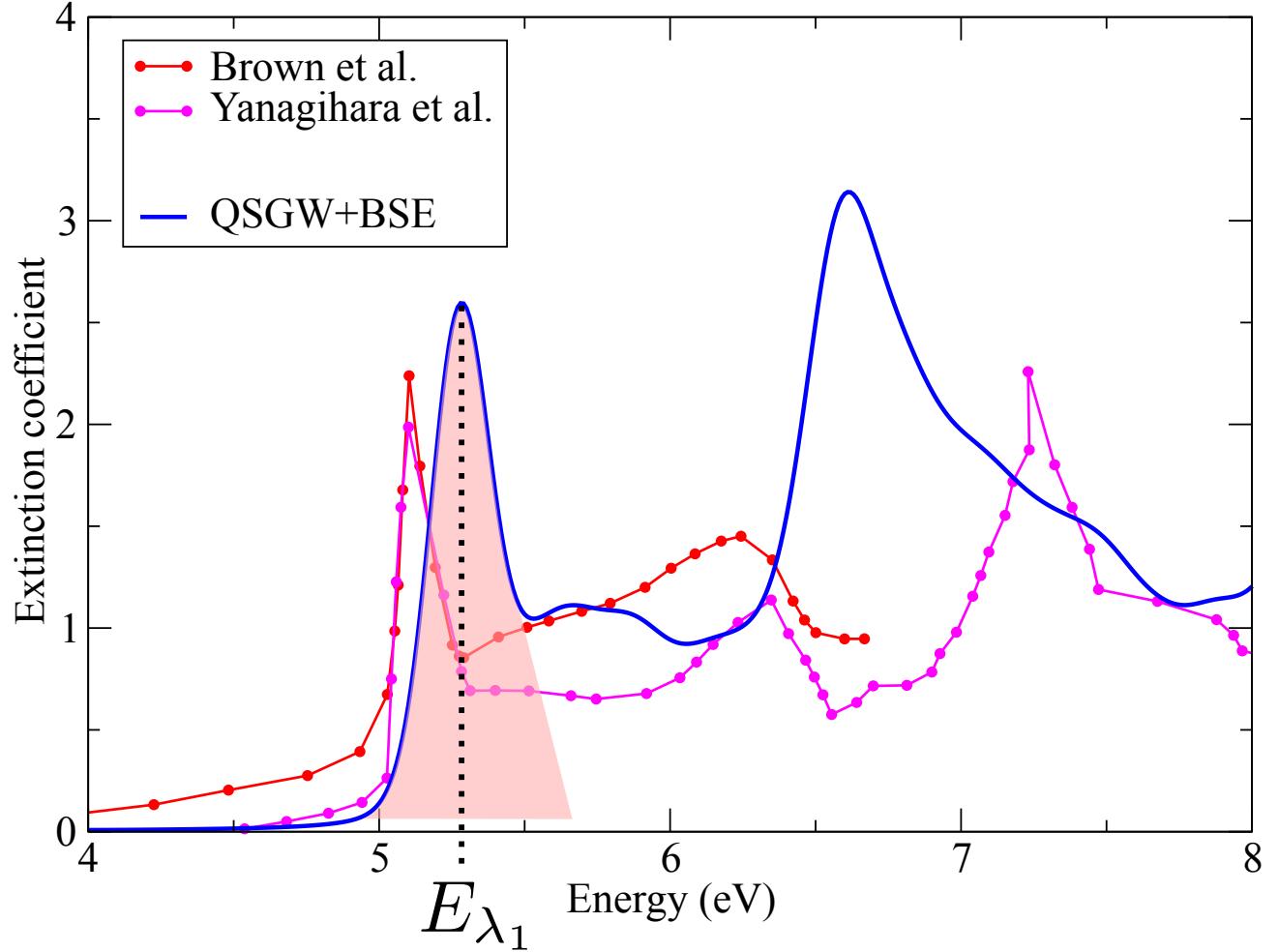
AgCl absorption

$$\chi_M = \sum_{\lambda} \frac{\left| \sum_{vck} A_{\lambda}^{vc\mathbf{k}} \langle c\mathbf{k} | \hat{d} | v\mathbf{k} \rangle \right|^2}{\omega - E_{\lambda} + i\eta}$$

$$\kappa = \text{Im} \sqrt{\frac{1}{1 + v_0 \chi_M}}$$

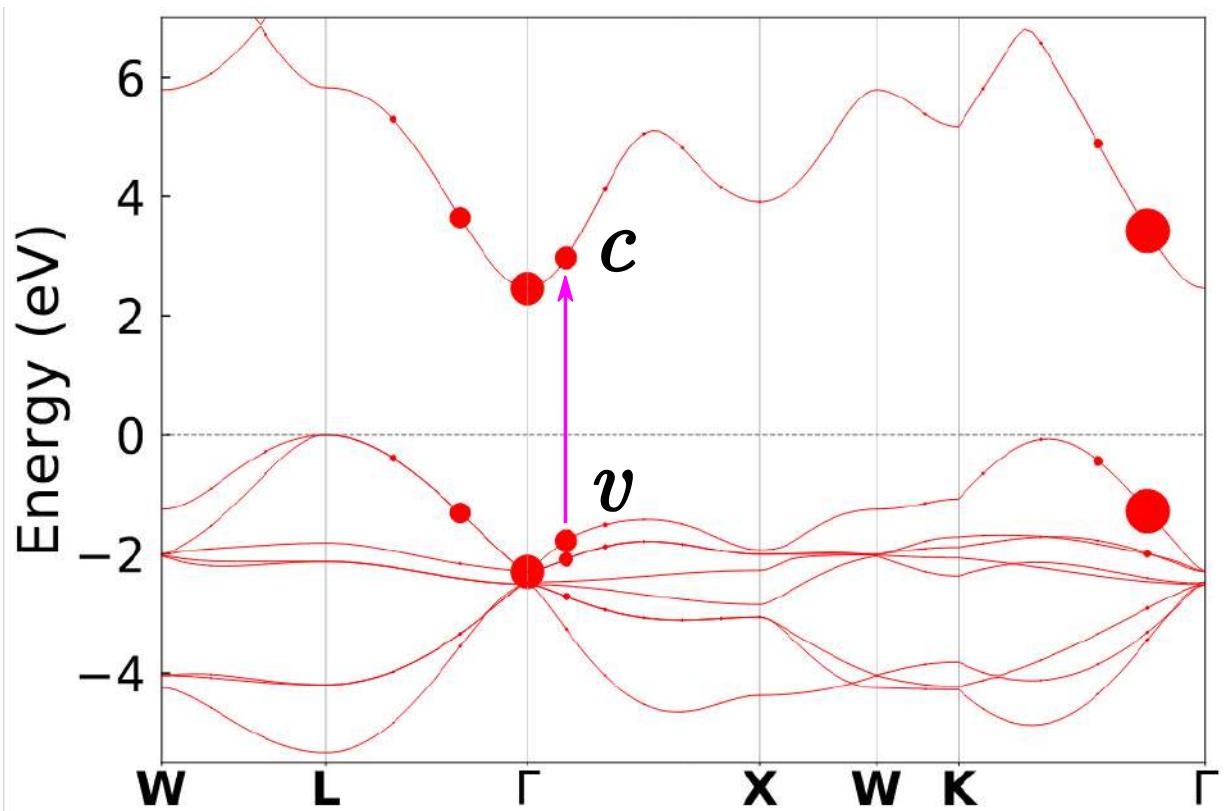


Lorin et al. Phys. Rev. B **104**, 235149 (2021)



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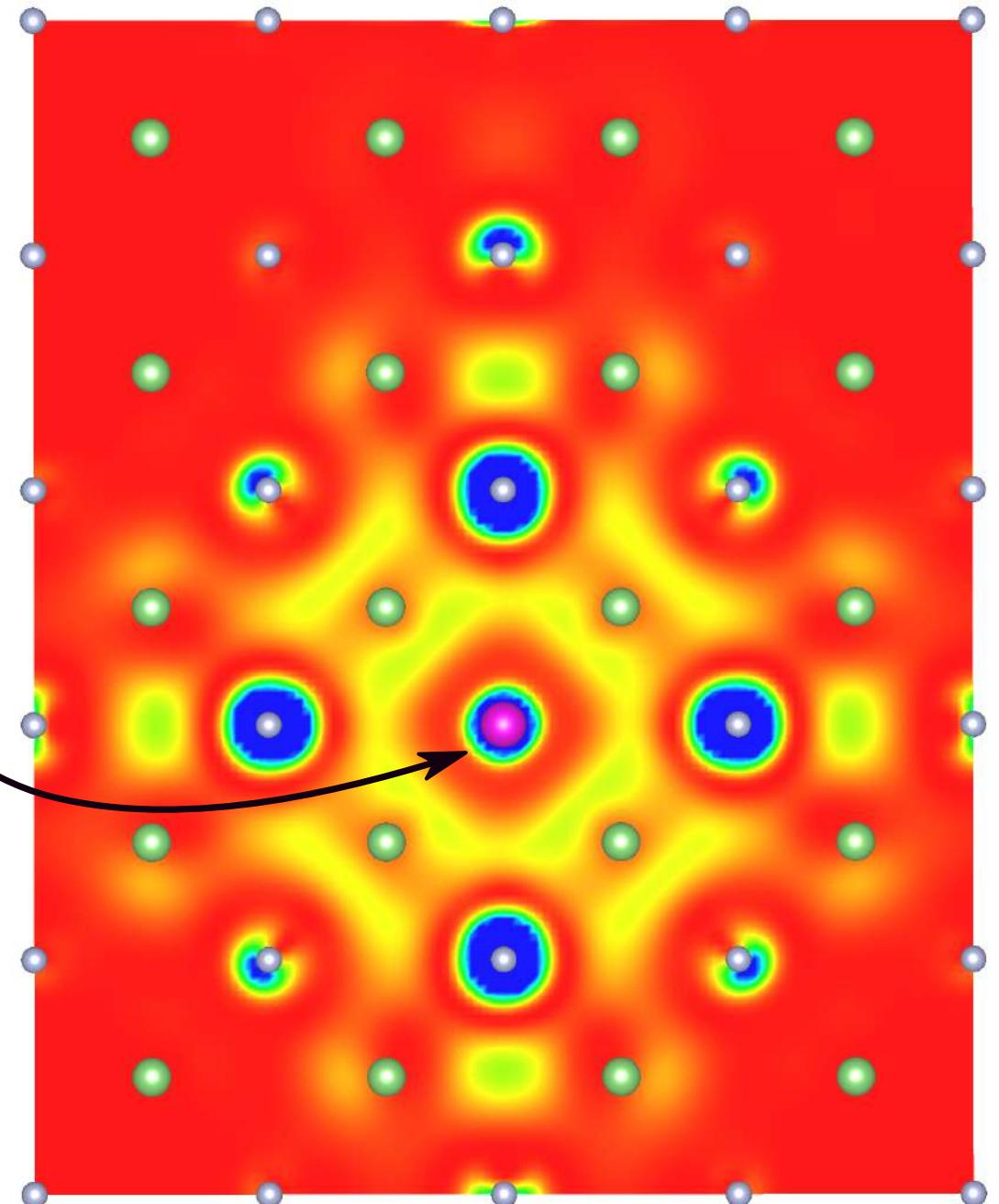


Excitonic wavefunction of LiF

$$\Psi_\lambda(\mathbf{r}_e, \mathbf{r}_h) = \sum_{vck} A_\lambda^{vck} \psi_{c\mathbf{k}}^*(\mathbf{r}_e) \psi_{v\mathbf{k}}(\mathbf{r}_h)$$

Excitonic wavefunction of LiF

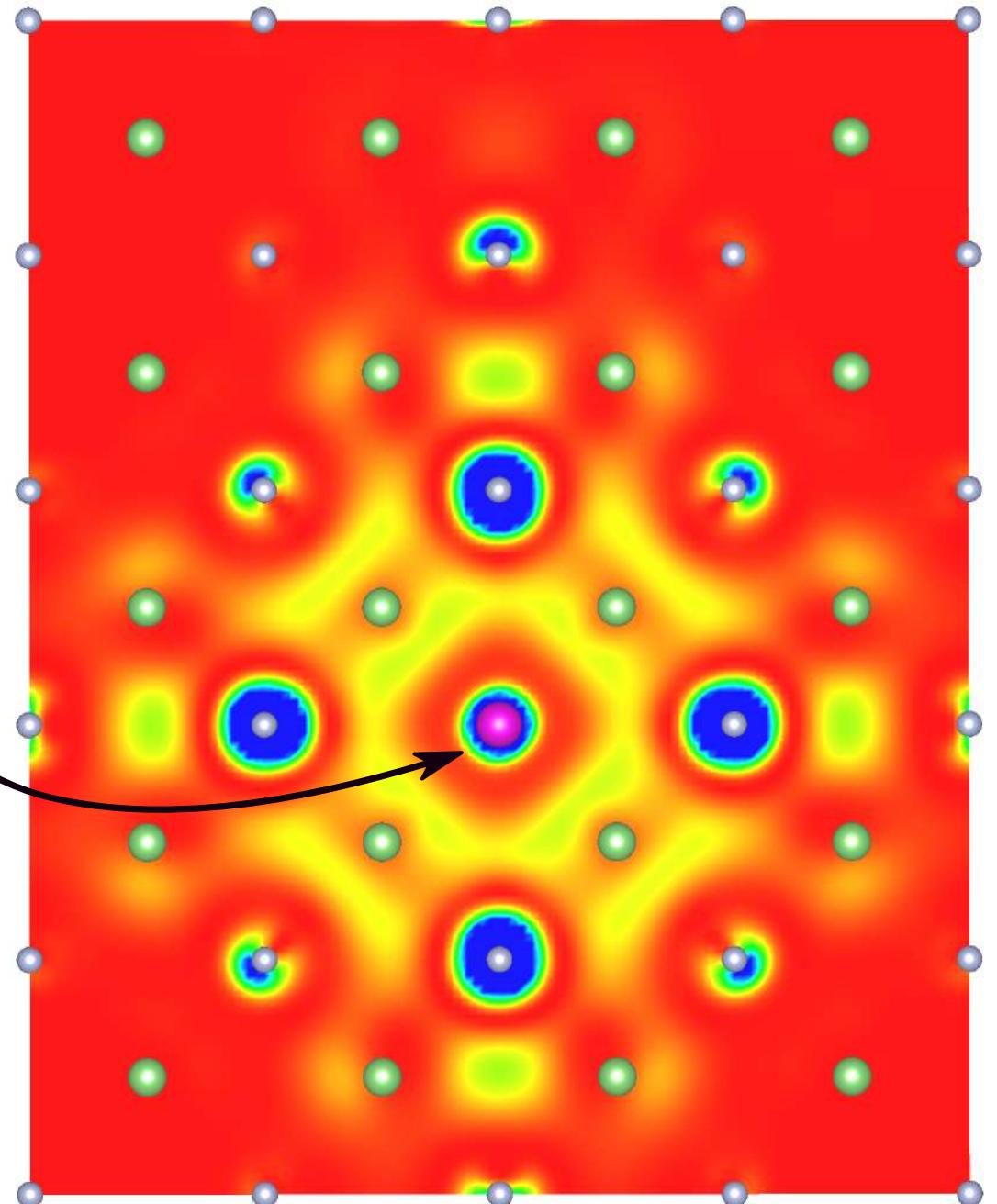
$$|\Psi_\lambda(\mathbf{r}_e, \mathbf{r}_h)|^2 = \left| \sum_{vck} A_{\lambda}^{vck} \psi_{c\mathbf{k}}^*(\mathbf{r}_e) \psi_{v\mathbf{k}}(\mathbf{r}_h) \right|^2$$



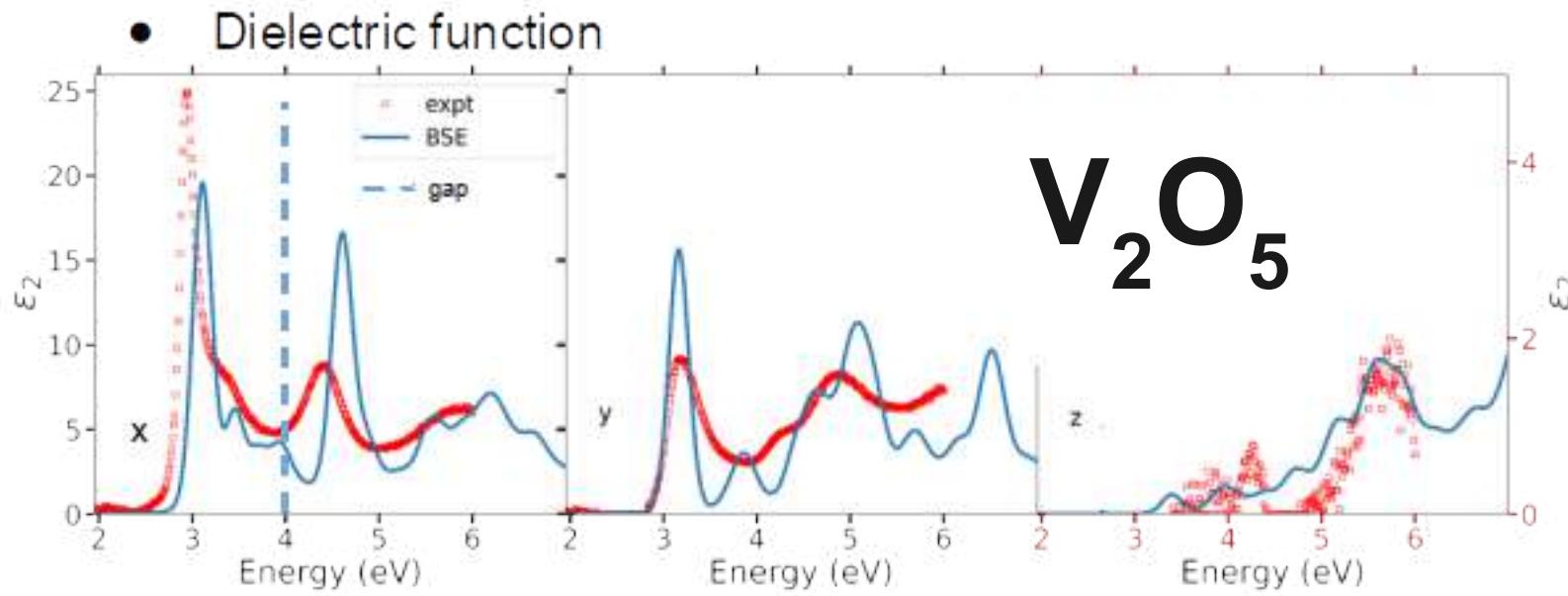
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- where is the exciton localised ?
- how much ?

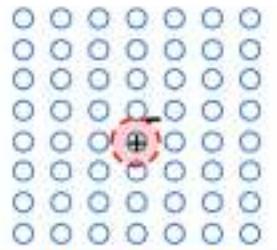


Gatti and Sottile, Phys. Rev. B **98**, 155113 (2013)

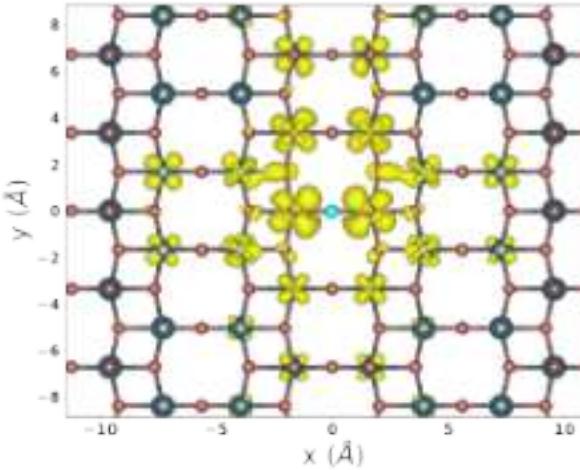


• Exciton wave function

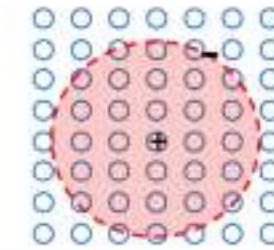
Textbook



Frenkel exciton
Binding energy ~1 eV



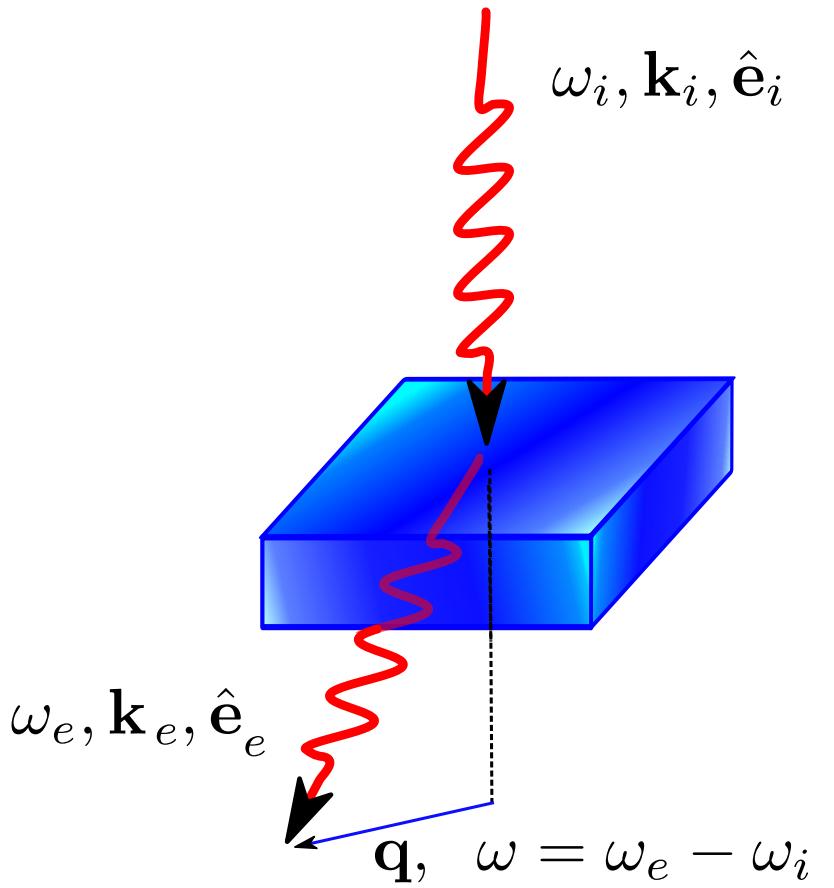
Textbook



Wannier-Mott exciton
Binding energy ~10 meV



X-ray scattering

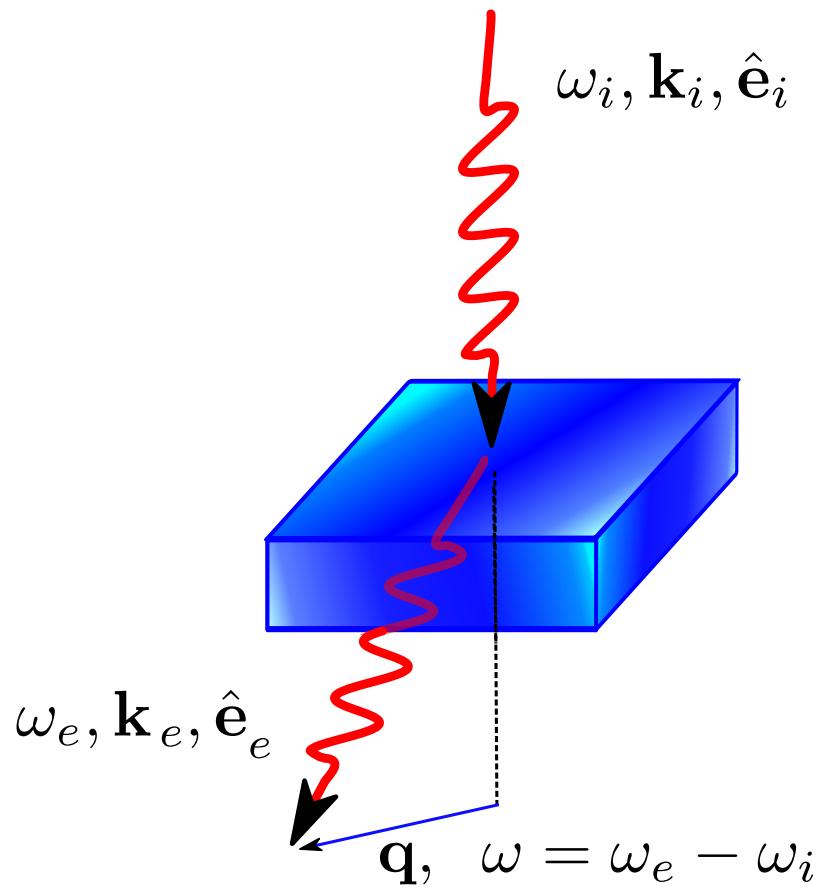


$$\frac{d^2\sigma}{d\Omega_2 d\omega_e} \propto \sum_f \left| \langle f | e^{i\mathbf{q} \cdot \mathbf{r}} | 0 \rangle - \frac{i\omega_{i/e}}{2mc^2} \sum_n \frac{\langle f | e^{-i\mathbf{k}_f \cdot \mathbf{r}} \nabla | n \rangle \langle n | e^{i\mathbf{k}_i \cdot \mathbf{r}} \nabla | 0 \rangle}{\omega_i - (E_n - E_0)} \right|^2 \times \delta(\omega - (E_f - E_0))$$

X-ray scattering

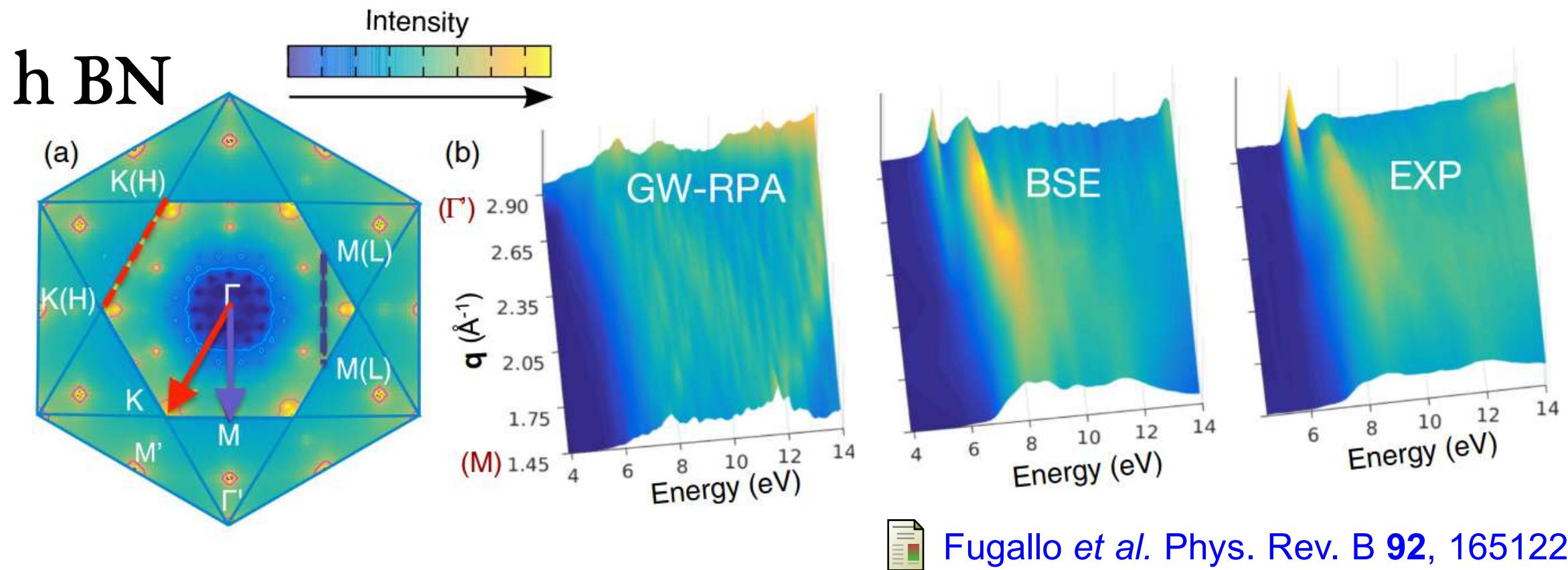
non-Resonant IXS

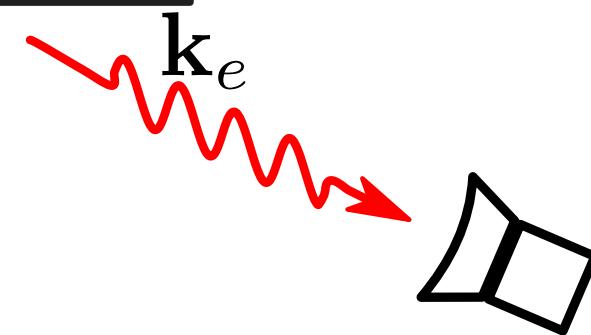
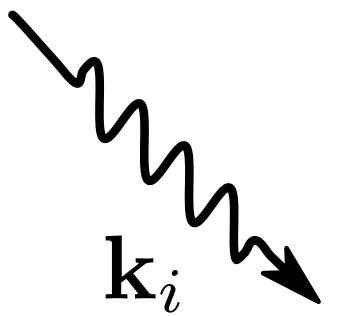
$$\frac{d^2\sigma}{d\Omega_2 d\omega_e} \propto \sum_f \left| \langle f | e^{i\mathbf{q} \cdot \mathbf{r}} | 0 \rangle - \frac{i\omega_{i/e}}{2mc^2} \sum_n \frac{\langle f | e^{-i\mathbf{k}_f \cdot \mathbf{r}} \nabla | n \rangle \langle n | e^{i\mathbf{k}_i \cdot \mathbf{r}} \nabla | 0 \rangle}{\omega_i - (E_n - E_0)} \right|^2 \times \delta(\omega - (E_f - E_0))$$

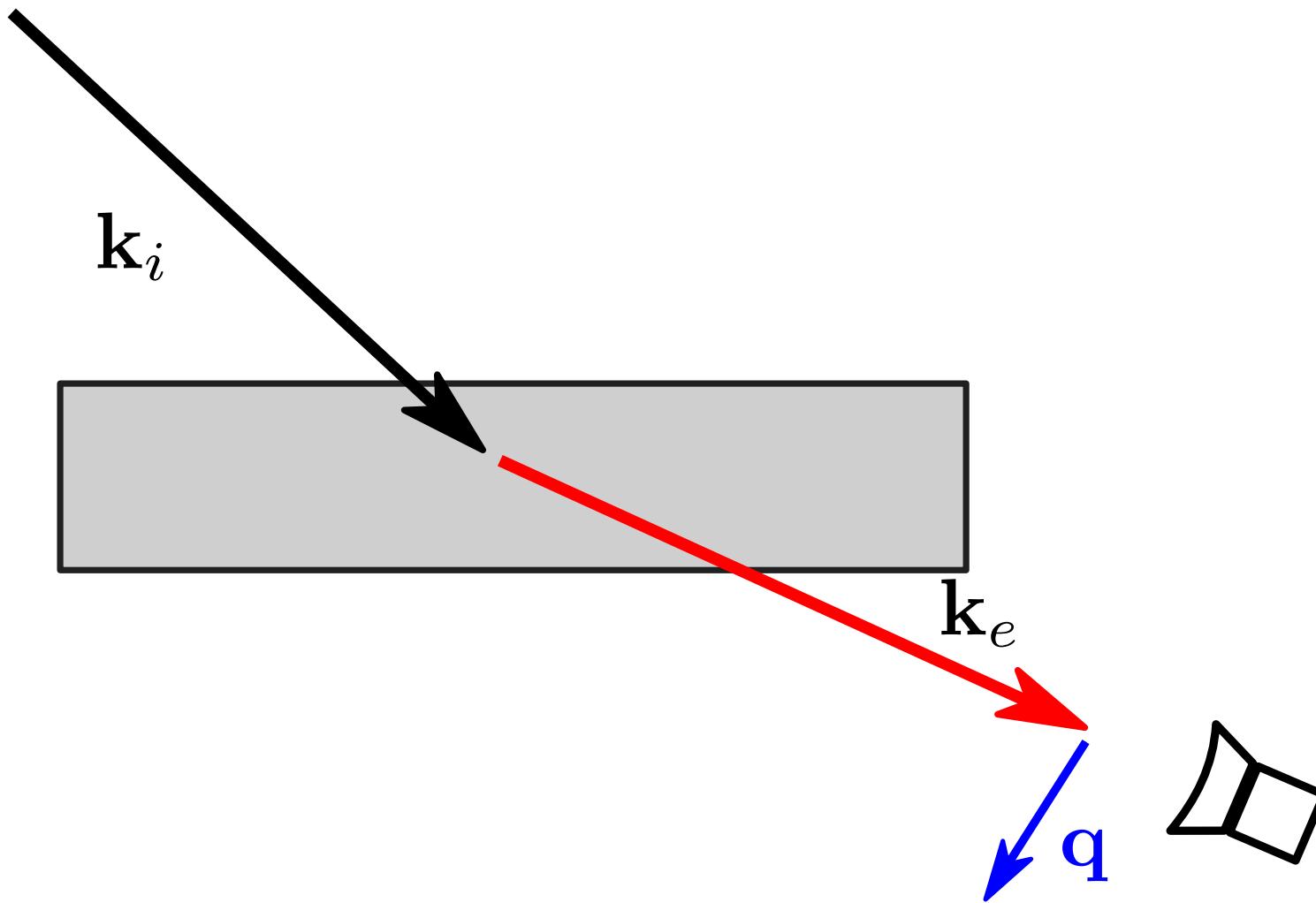


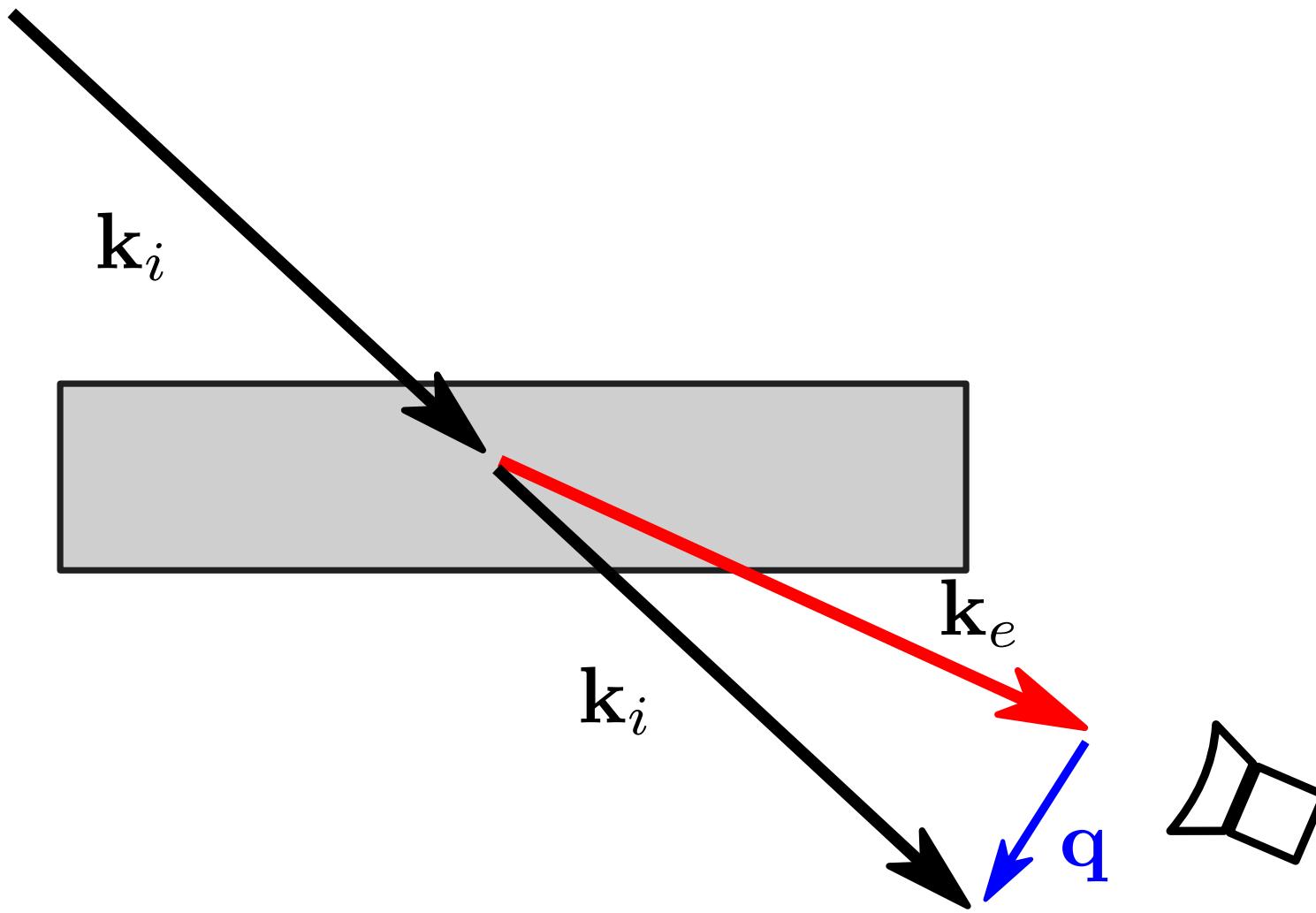
Bethe-Salpeter Equation - finite momentum transfer

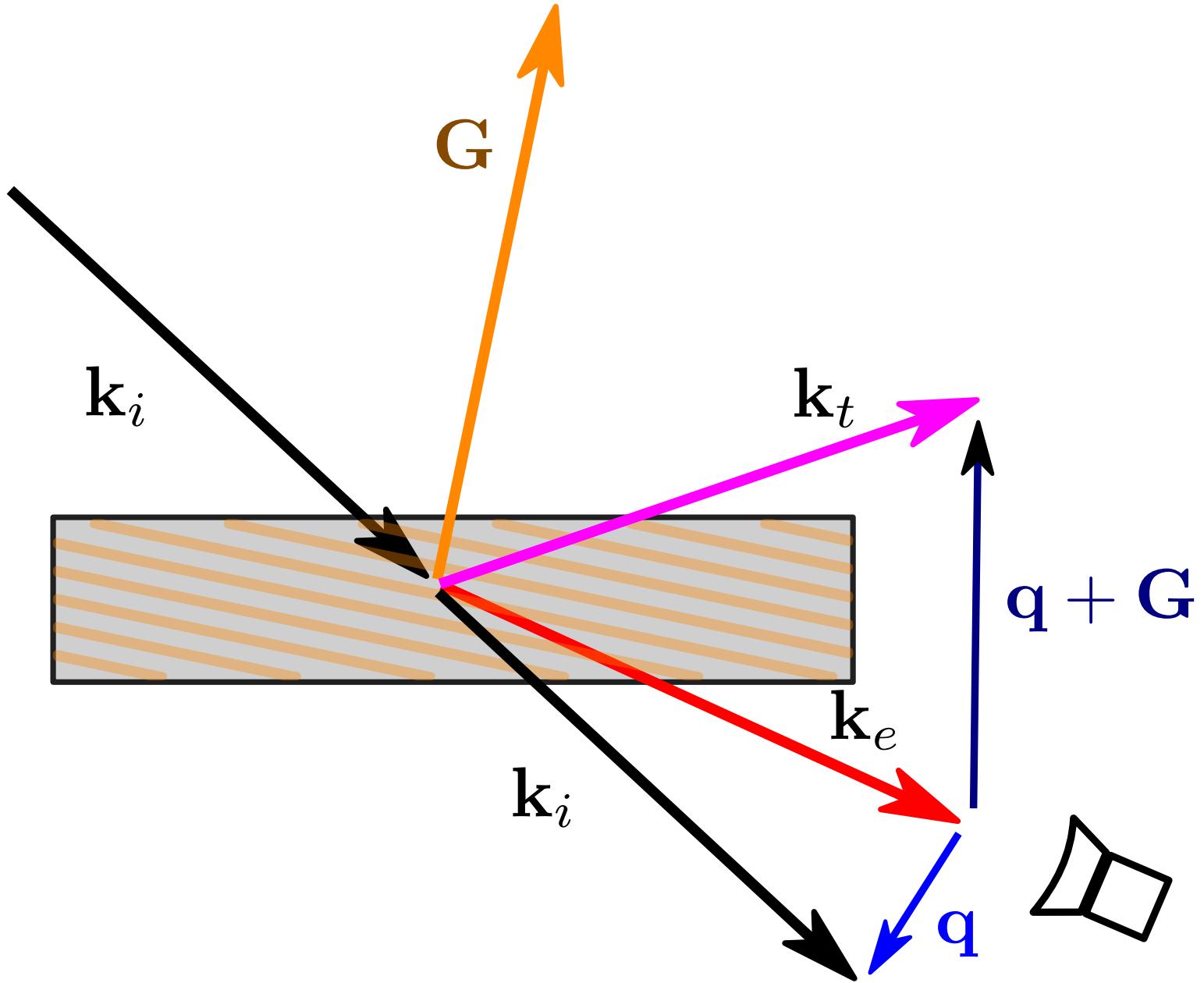
$$S(\mathbf{q}, \omega) \propto \chi_M(\mathbf{q}, \omega) = \sum_{\lambda} \frac{\left| \sum_{vc} A_{\lambda}^{vc, \mathbf{q}} \langle c | e^{i\mathbf{q} \cdot \mathbf{r}} | v \rangle \right|^2}{\omega - E_{\lambda}(\mathbf{q}) + i\eta}$$

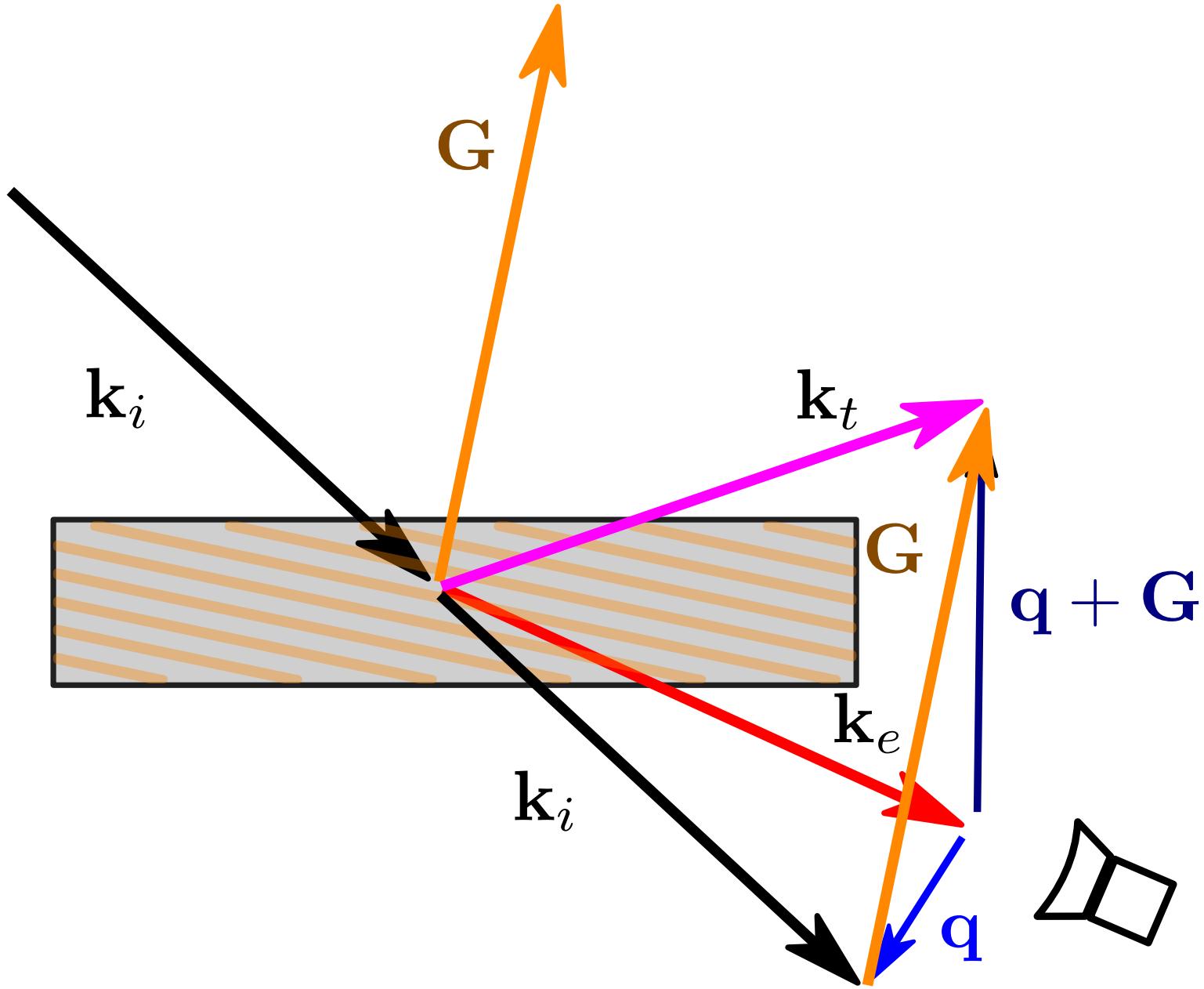




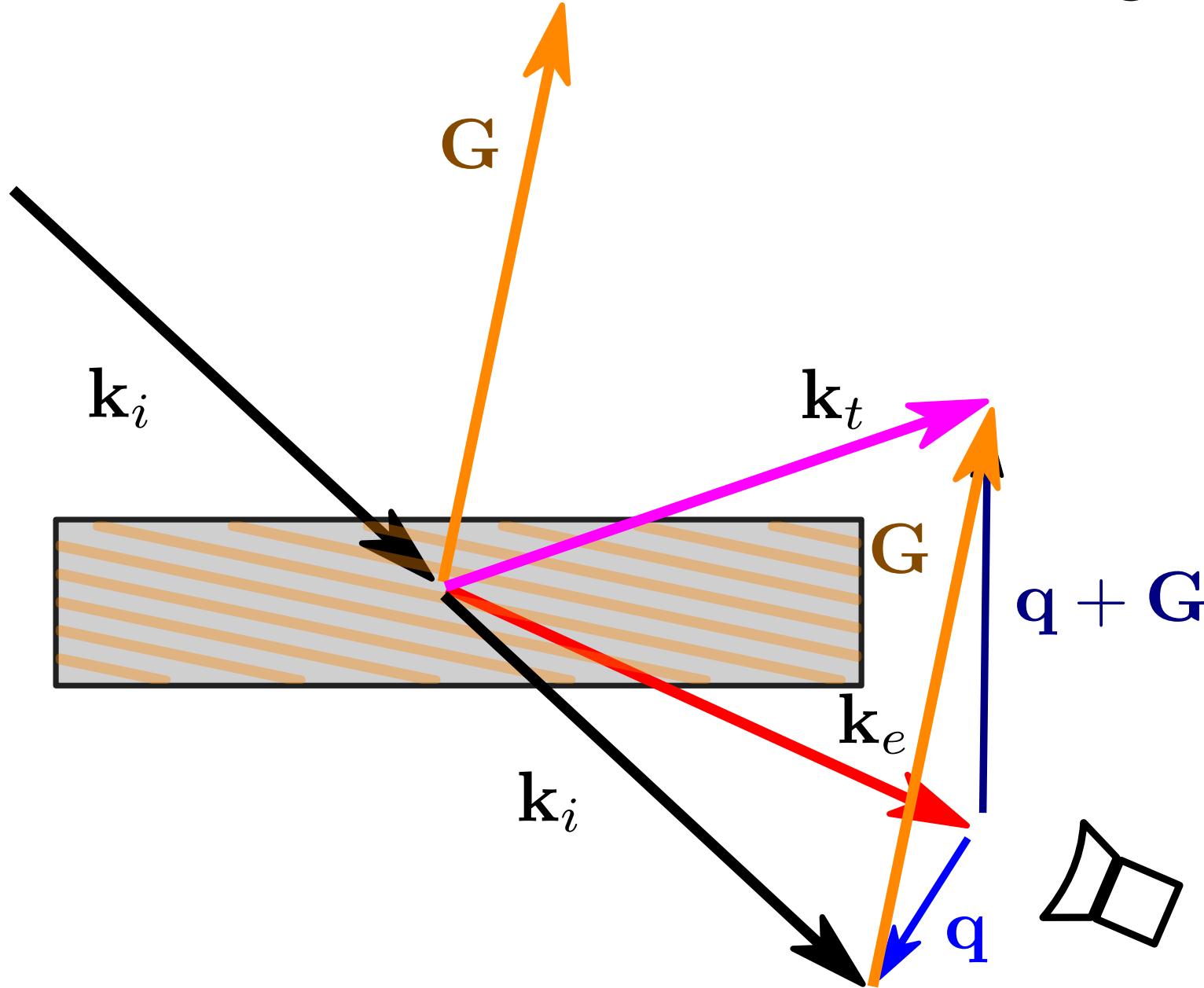








Coherent Inelastic X-ray scattering



Coherent Inelastic X-ray scattering

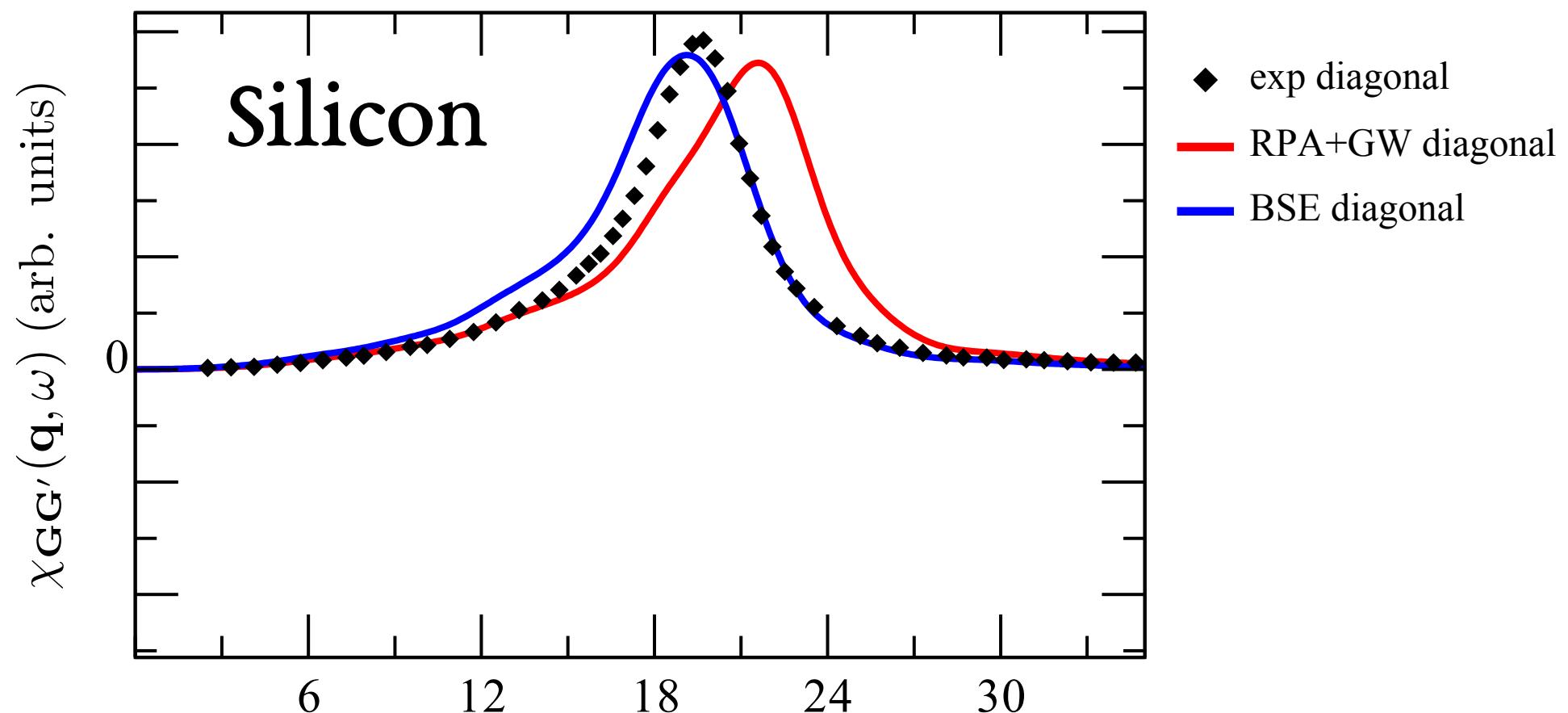
$$\chi(\mathbf{q}, \mathbf{q} + \mathbf{G}, \omega) = \sum_{\lambda\lambda'} \frac{\sum_{vc} A_{\lambda}^{vc, \mathbf{q}} \langle c | e^{-i\mathbf{q}\cdot\mathbf{r}} | v \rangle S_{\lambda\lambda'}^{-1} \sum_{v'c'} A_{\lambda}^{*, v'c', \mathbf{q}} \langle v' | e^{i(\mathbf{q}+\mathbf{G})\cdot\mathbf{r}} | c' \rangle}{\omega - E_{\lambda}(\mathbf{q}) + i\eta}$$

Coherent Inelastic X-ray scattering

$$\chi(\mathbf{q}, \mathbf{q} + \mathbf{G}, \omega) = \sum_{\lambda\lambda'} \frac{\sum_{vc} A_{\lambda}^{vc, \mathbf{q}} \langle c | e^{-i\mathbf{q}\cdot\mathbf{r}} | v \rangle S_{\lambda\lambda'}^{-1} \sum_{v'c'} A_{\lambda}^{*, v'c', \mathbf{q}} \langle v' | e^{i(\mathbf{q}+\mathbf{G})\cdot\mathbf{r}} | c' \rangle}{\omega - E_{\lambda}(\mathbf{q}) + i\eta}$$



Igor Reshetnyak *et al.*
Phys. Rev. Research **1**,
032010(R) (2019)



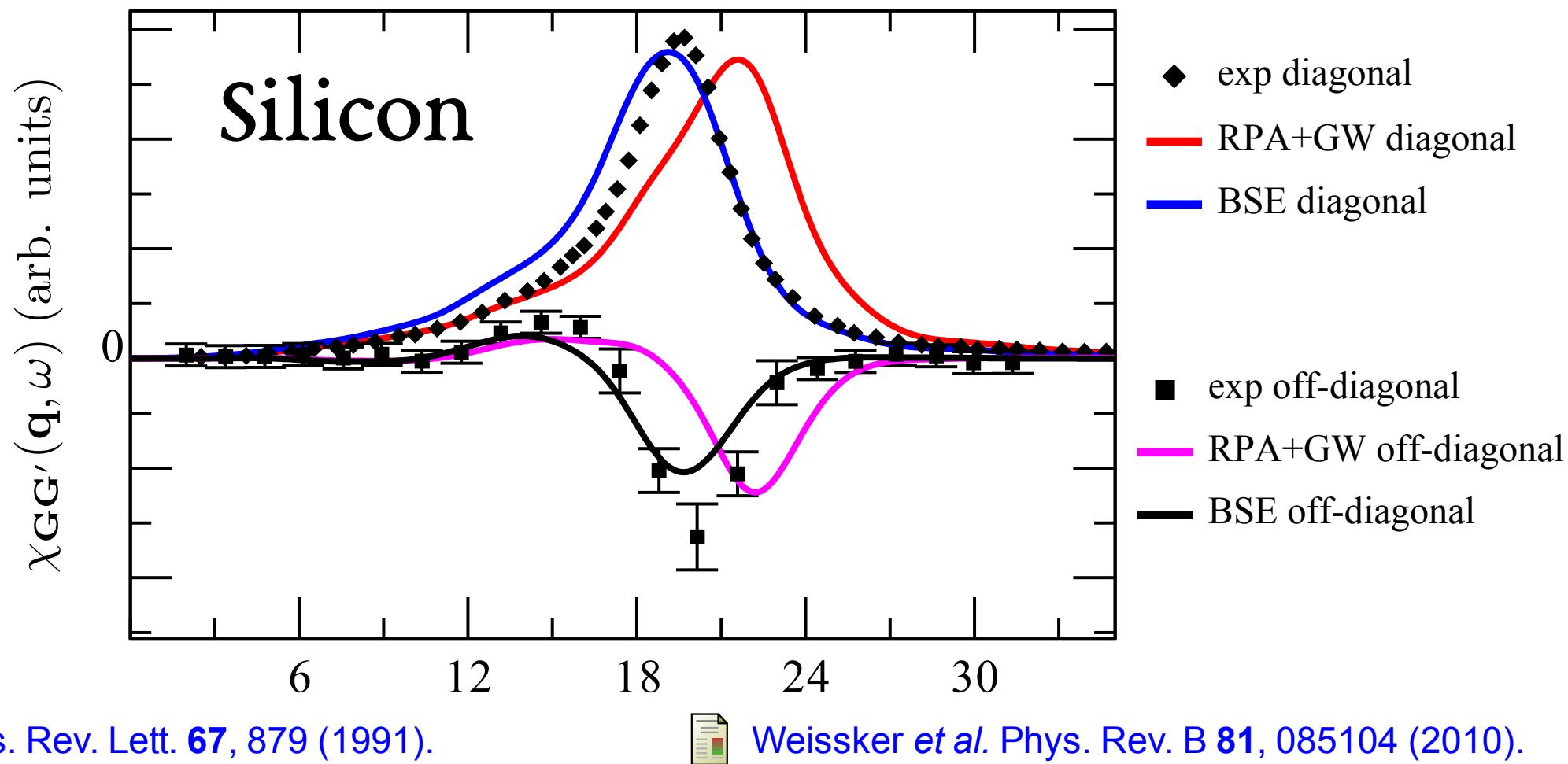
Weissker *et al.* Phys. Rev. B **81**, 085104 (2010).

Coherent Inelastic X-ray scattering

$$\chi(\mathbf{q}, \mathbf{q} + \mathbf{G}, \omega) = \sum_{\lambda\lambda'} \frac{\sum_{vc} A_{\lambda}^{vc, \mathbf{q}} \langle c | e^{-i\mathbf{q}\cdot\mathbf{r}} | v \rangle S_{\lambda\lambda'}^{-1} \sum_{v'c'} A_{\lambda}^{*, v'c', \mathbf{q}} \langle v' | e^{i(\mathbf{q}+\mathbf{G})\cdot\mathbf{r}} | c' \rangle}{\omega - E_{\lambda}(\mathbf{q}) + i\eta}$$



Igor Reshetnyak *et al.*
Phys. Rev. Research **1**,
032010(R) (2019)



Schülke and Kaprolat, Phys. Rev. Lett. **67**, 879 (1991).



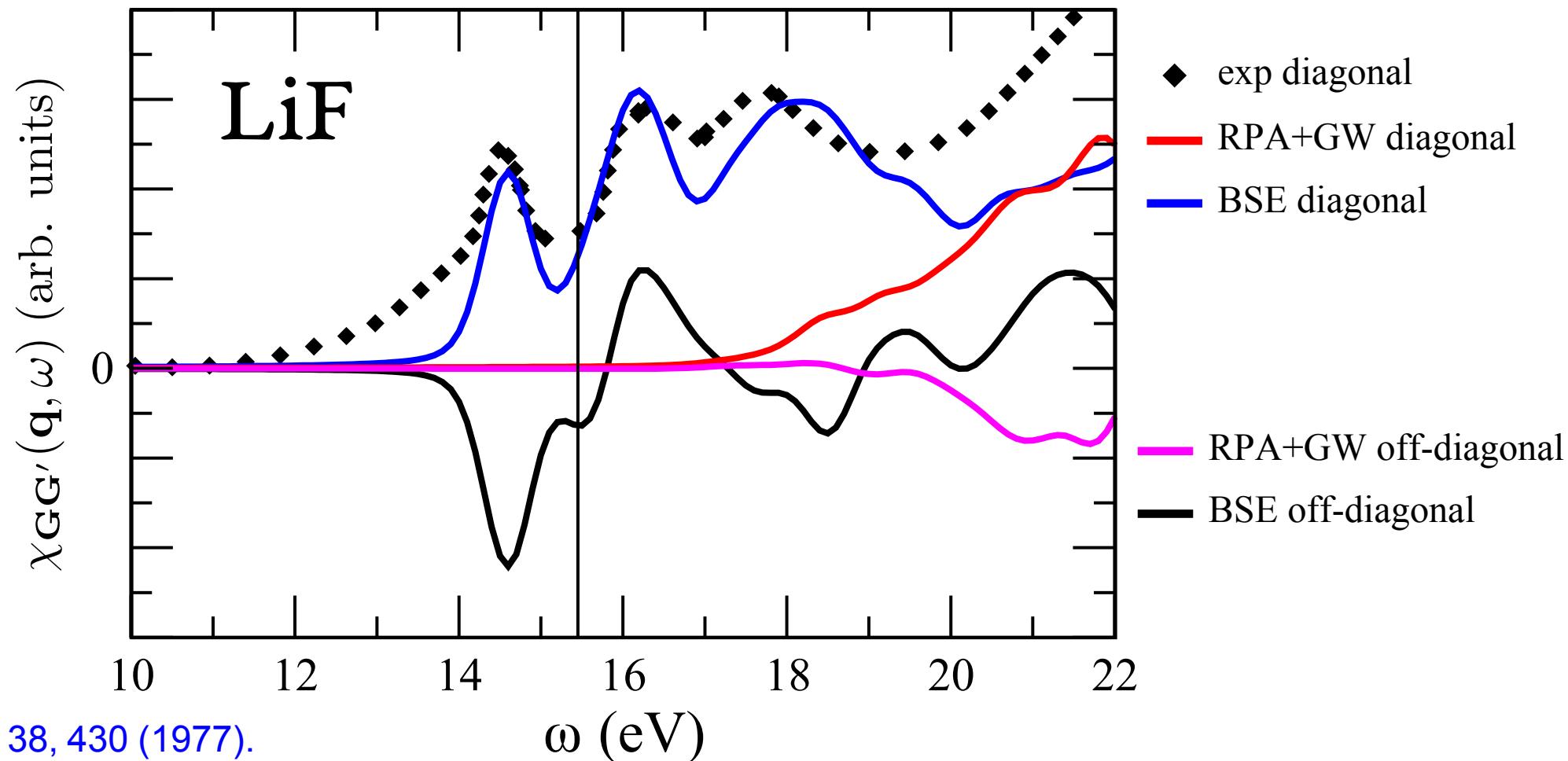
Weissker *et al.* Phys. Rev. B **81**, 085104 (2010).

Coherent Inelastic X-ray scattering

$$\chi(\mathbf{q}, \mathbf{q} + \mathbf{G}, \omega) = \sum_{\lambda\lambda'} \frac{\sum_{vc} A_{\lambda}^{vc, \mathbf{q}} \langle c | e^{-i\mathbf{q}\cdot\mathbf{r}} | v \rangle S_{\lambda\lambda'}^{-1} \sum_{v'c'} A_{\lambda}^{*, v'c', \mathbf{q}} \langle v' | e^{i(\mathbf{q}+\mathbf{G})\cdot\mathbf{r}} | c' \rangle}{\omega - E_{\lambda}(\mathbf{q}) + i\eta}$$



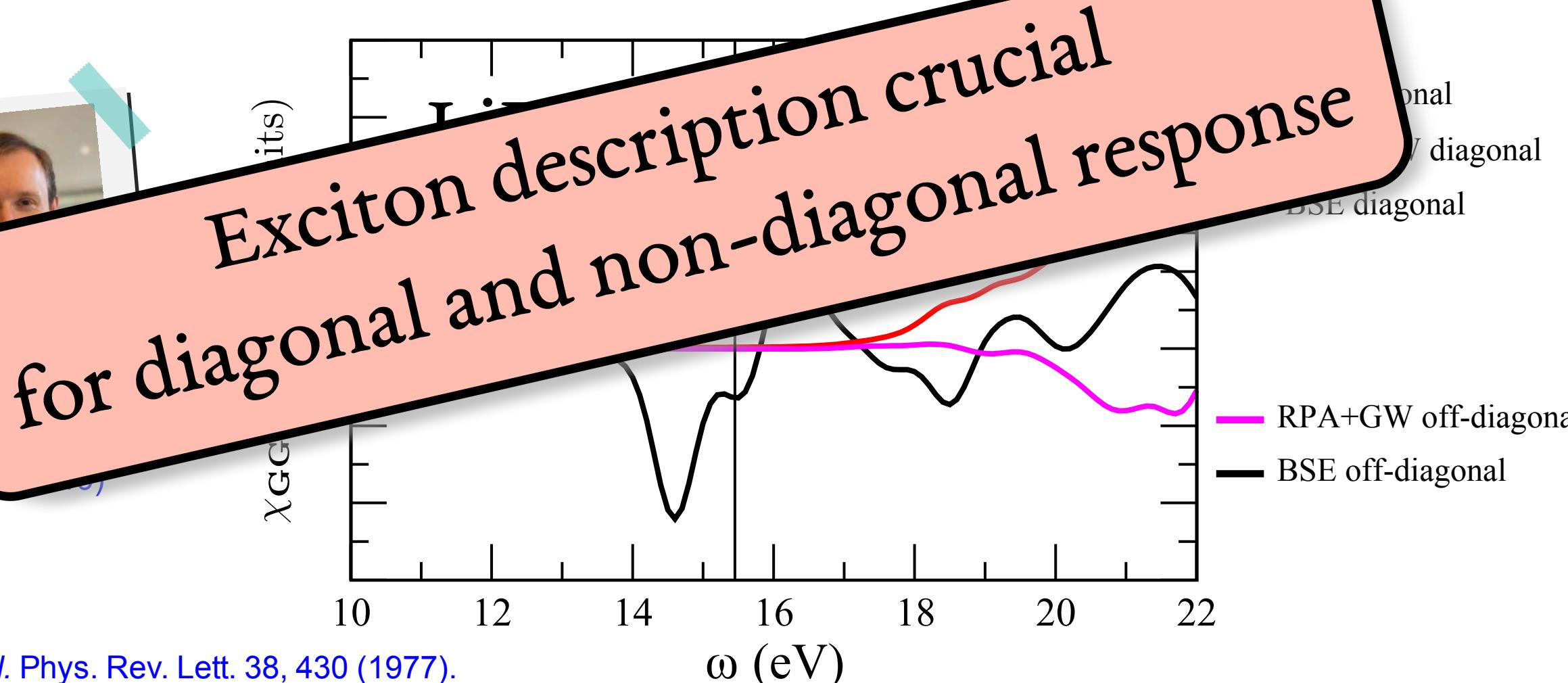
Igor Reshetnyak *et al.*
Phys. Rev. Research 1,
032010(R) (2019)



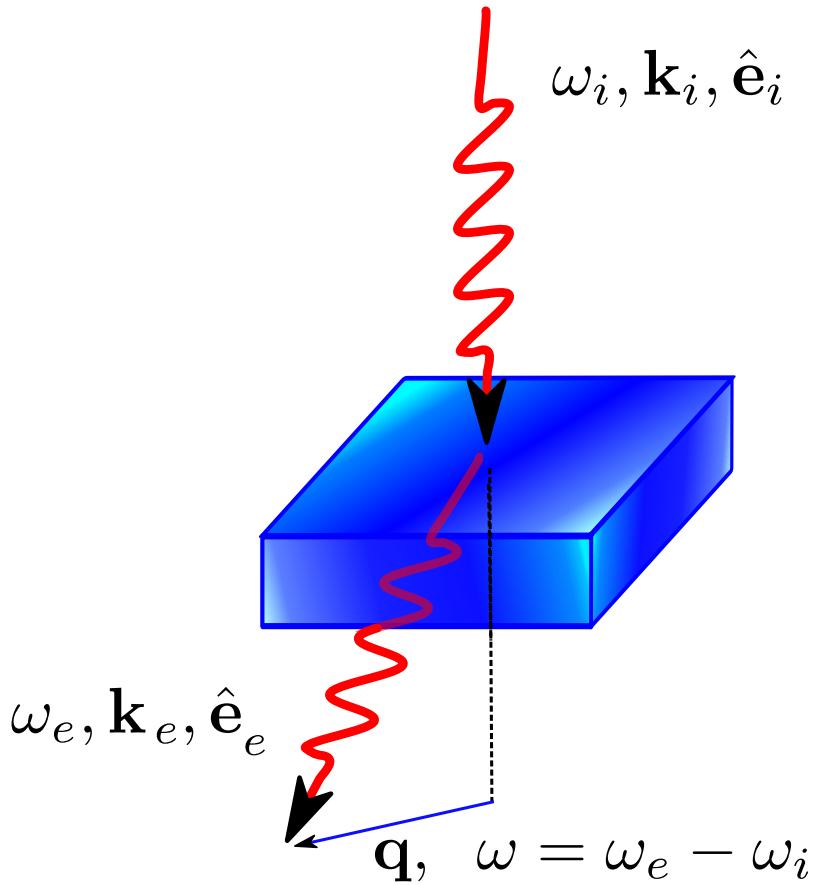
Fields *et al.* Phys. Rev. Lett. 38, 430 (1977).

Coherent Inelastic X-ray scattering

$$\chi(\mathbf{q}, \mathbf{q} + \mathbf{G}, \omega) = \sum_{\lambda\lambda'} \frac{\sum_{vc} A_{\lambda}^{vc, \mathbf{q}} \langle c | e^{-i\mathbf{q}\cdot\mathbf{r}} | v \rangle S_{\lambda\lambda'}^{-1} \sum_{v'c'} A_{\lambda}^{*, v'c', \mathbf{q}} \langle v' | e^{i(\mathbf{q}+\mathbf{G})\cdot\mathbf{r}} | c' \rangle}{\omega - E_{\lambda}(\mathbf{q}) + i\eta}$$



X-ray scattering



non-Resonant IXS

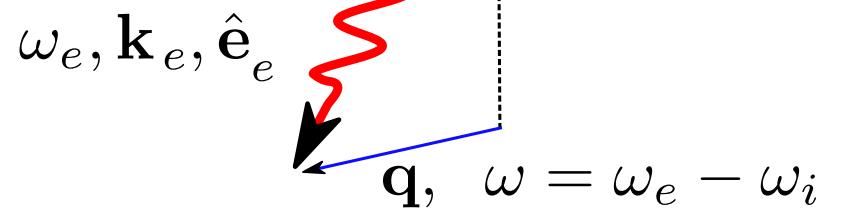
$$\frac{d^2\sigma}{d\Omega_2 d\omega_e} \propto \sum_f \left| \langle f | e^{i\mathbf{q} \cdot \mathbf{r}} | 0 \rangle - \frac{i\omega_{i/e}}{2mc^2} \sum_n \frac{\langle f | e^{-i\mathbf{k}_f \cdot \mathbf{r}} \nabla | n \rangle \langle n | e^{i\mathbf{k}_i \cdot \mathbf{r}} \nabla | 0 \rangle}{\omega_i - (E_n - E_0)} \right|^2 \times \delta(\omega - (E_f - E_0))$$

X-ray scattering

non-Resonant IXS

$$\frac{d^2\sigma}{d\Omega_2 d\omega_e} \propto \sum_f \left| \langle f | e^{i\mathbf{q} \cdot \mathbf{r}} | 0 \rangle - \frac{i\omega_{i/e}}{2mc^2} \sum_n \frac{\langle f | e^{-i\mathbf{k}_f \cdot \mathbf{r}} \nabla | n \rangle \langle n | e^{i\mathbf{k}_i \cdot \mathbf{r}} \nabla | 0 \rangle}{\omega_i - (E_n - E_0)} \right|^2 \times \delta(\omega - (E_f - E_0))$$

Resonant IXS

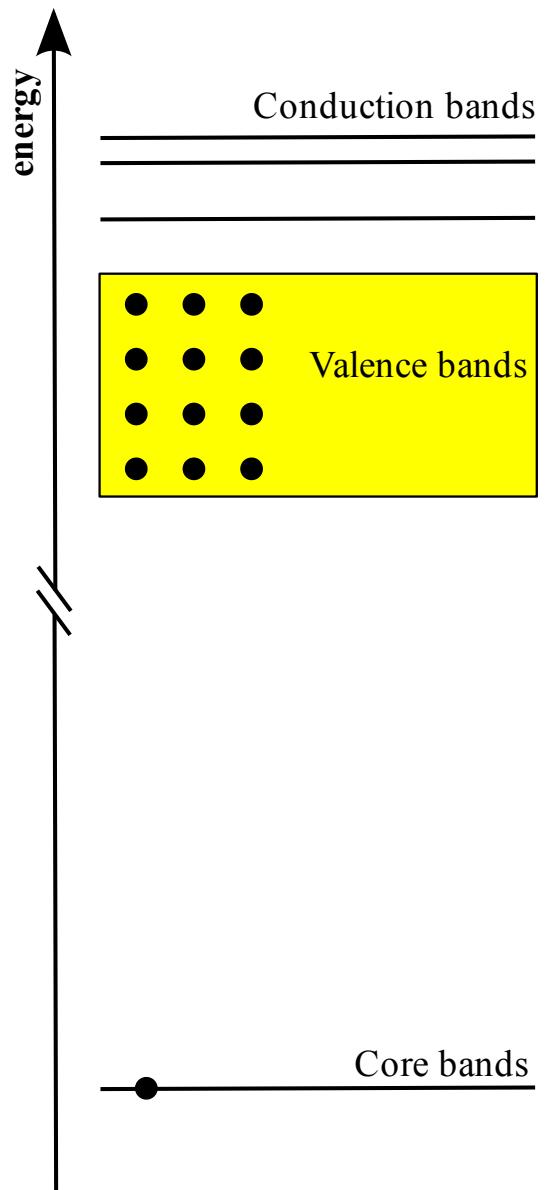


Resonant IXS

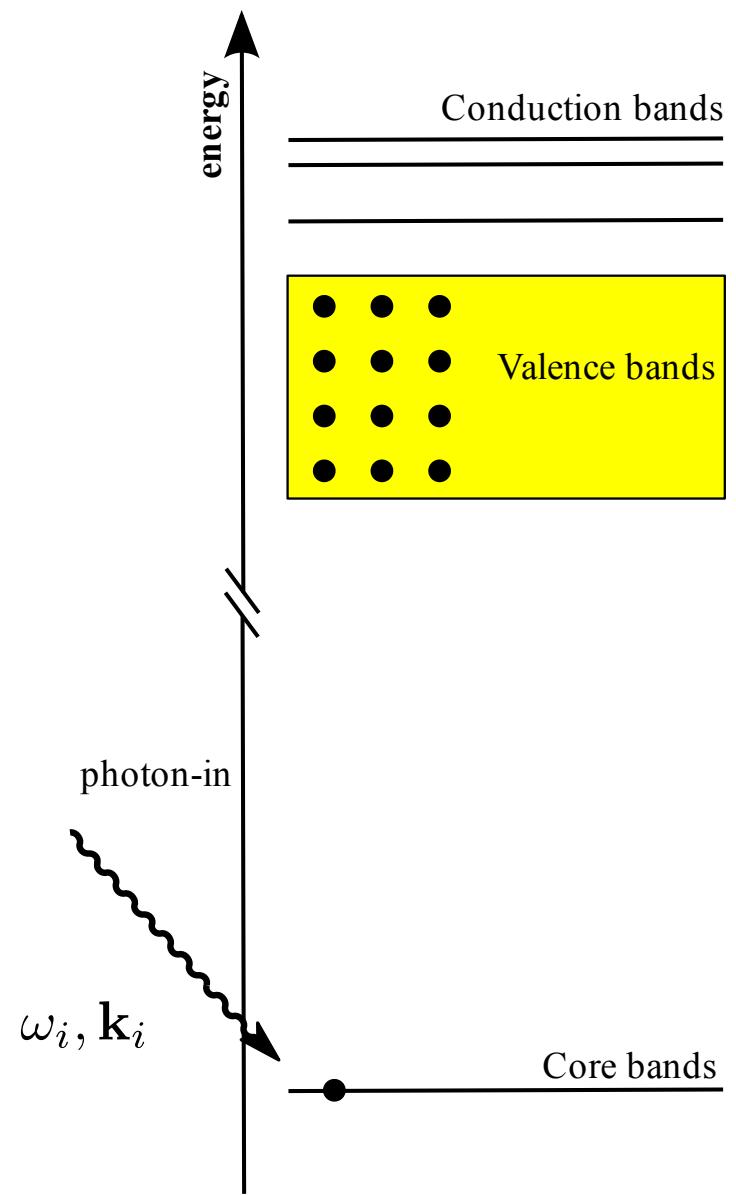
$$\frac{d^2\sigma}{d\Omega_2 d\omega_e} \propto \sum_f \left| \sum_n \frac{\langle f | e^{-i\mathbf{k}_f \cdot \mathbf{r}} \nabla | n \rangle \langle n | e^{i\mathbf{k}_i \cdot \mathbf{r}} \nabla | 0 \rangle}{\omega_i - (E_n - E_0)} \right|^2 \times \delta(\omega - (E_f - E_0))$$



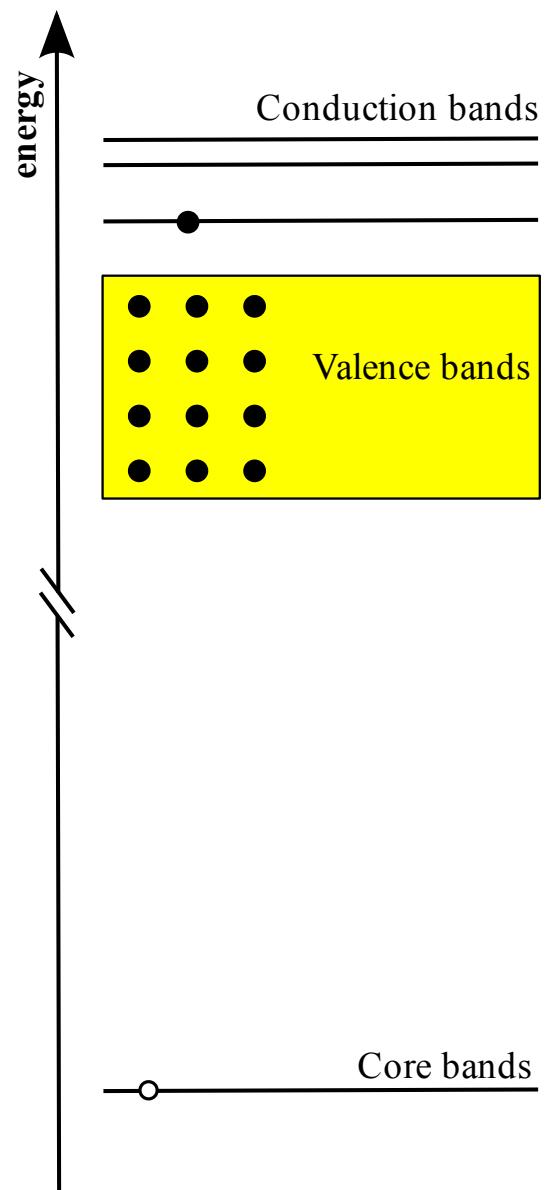
Initial state



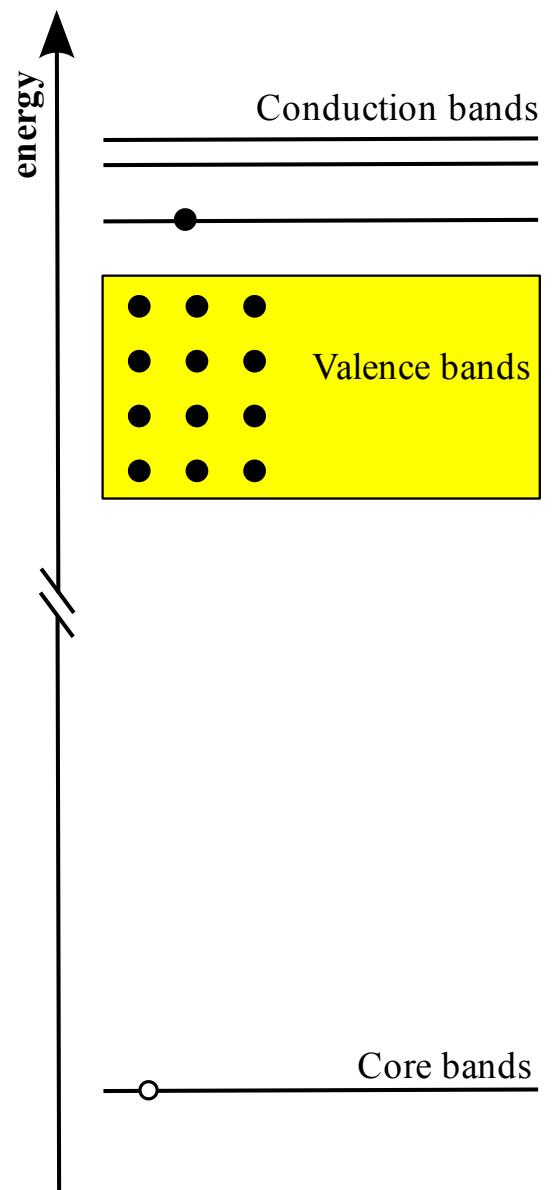
Initial state



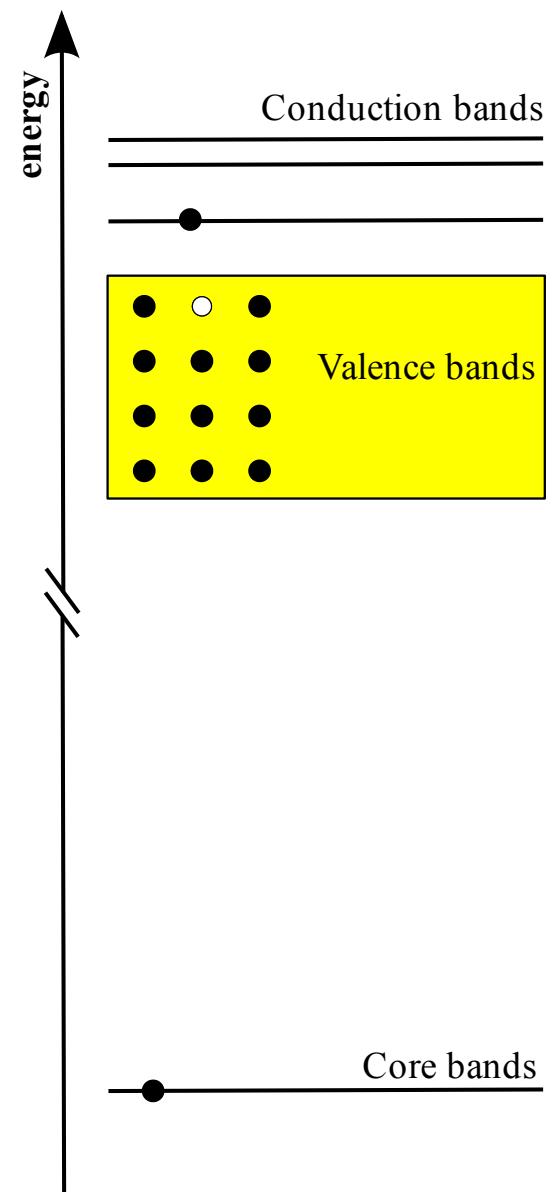
Intermediate state

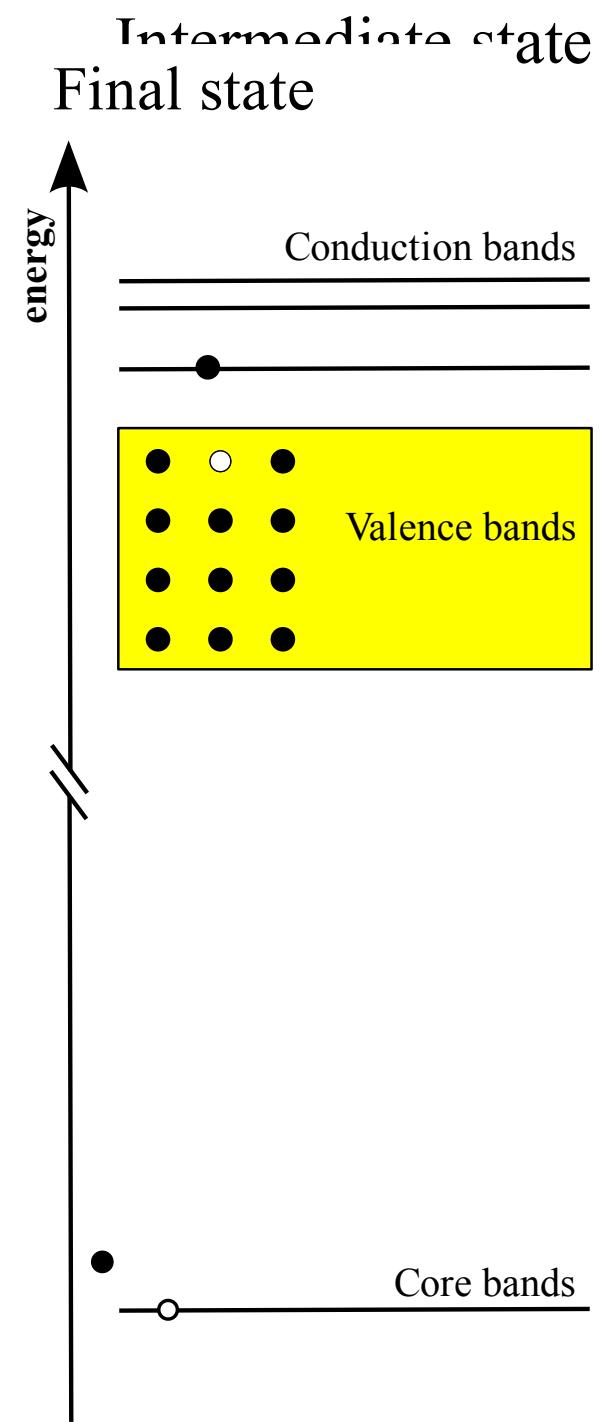
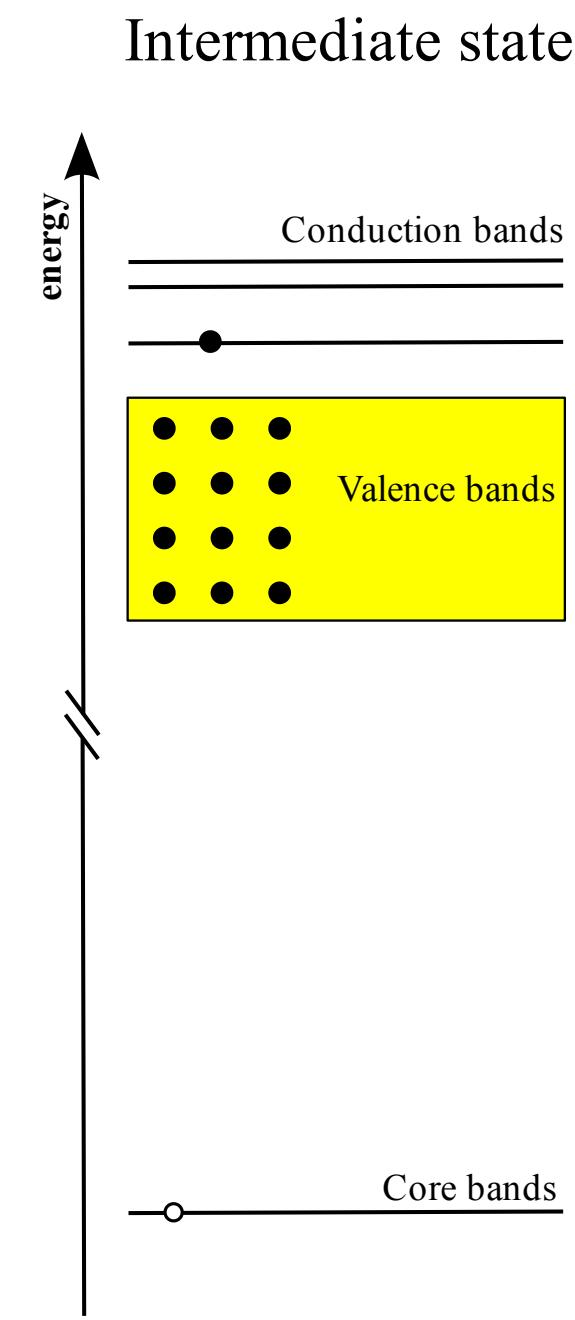
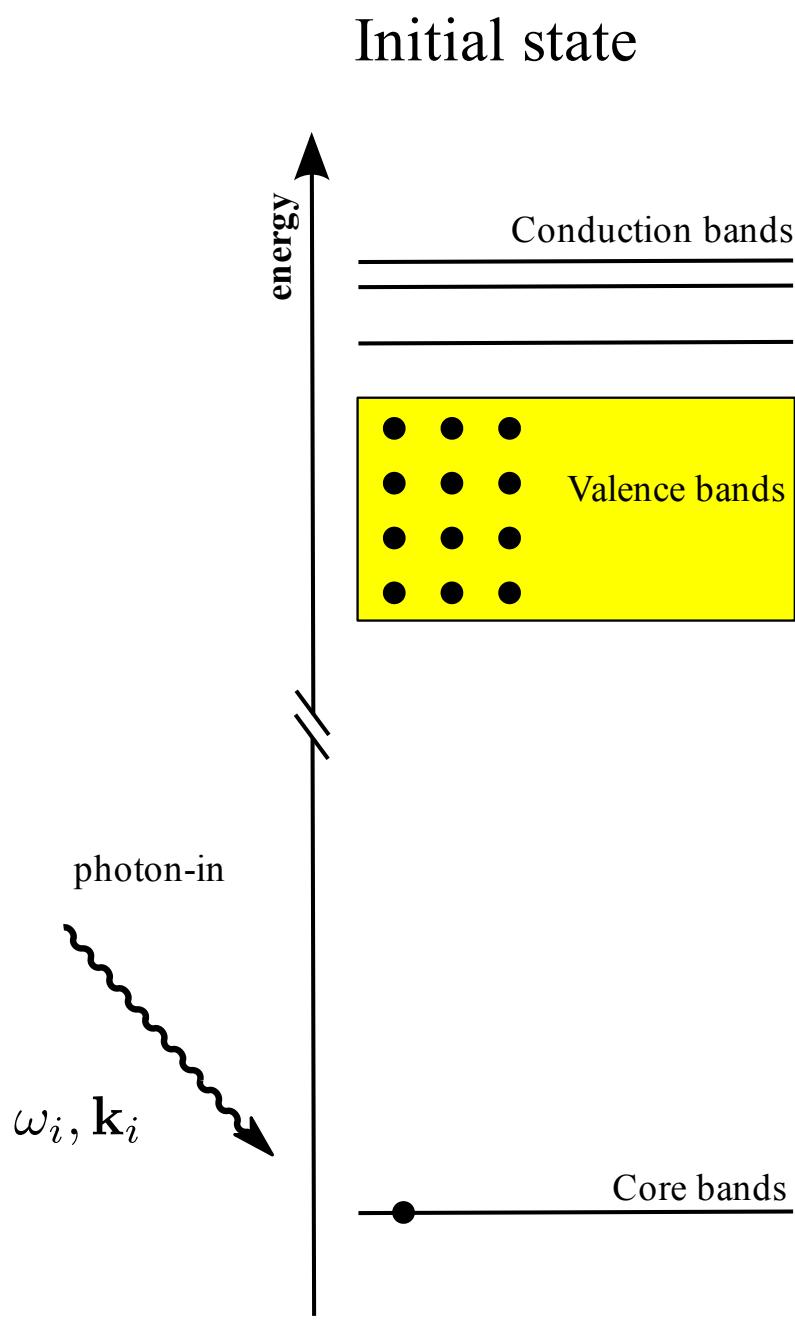


Intermediate state



Final state





$$\frac{d^2\sigma}{d\Omega_2 d\omega_e} \propto \sum_f \left| \sum_n \frac{\langle f | e^{-i\mathbf{k}_f \cdot \mathbf{r}} \nabla | n \rangle \langle n | e^{i\mathbf{k}_i \cdot \mathbf{r}} \nabla | 0 \rangle}{\omega_i - (E_n - E_0)} \right|^2 \times \delta(\omega - (E_f - E_0))$$



[Shirley, Phys. Rev. Lett. **80**, 794 \(1998\)](#)



[Vinson et al., Phys. Rev. B **94**, 035163 \(2016\)](#)

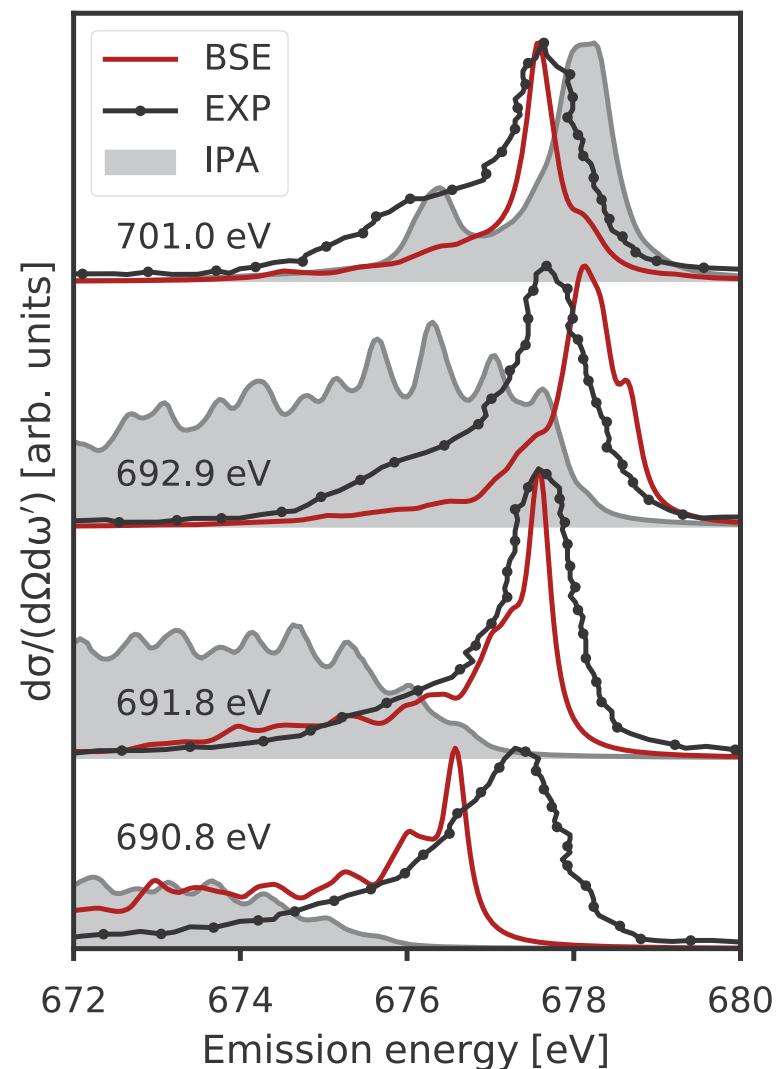


[Geondzhian and Gilmore, Phys. Rev. B **98**, 214305 \(2018\)](#)

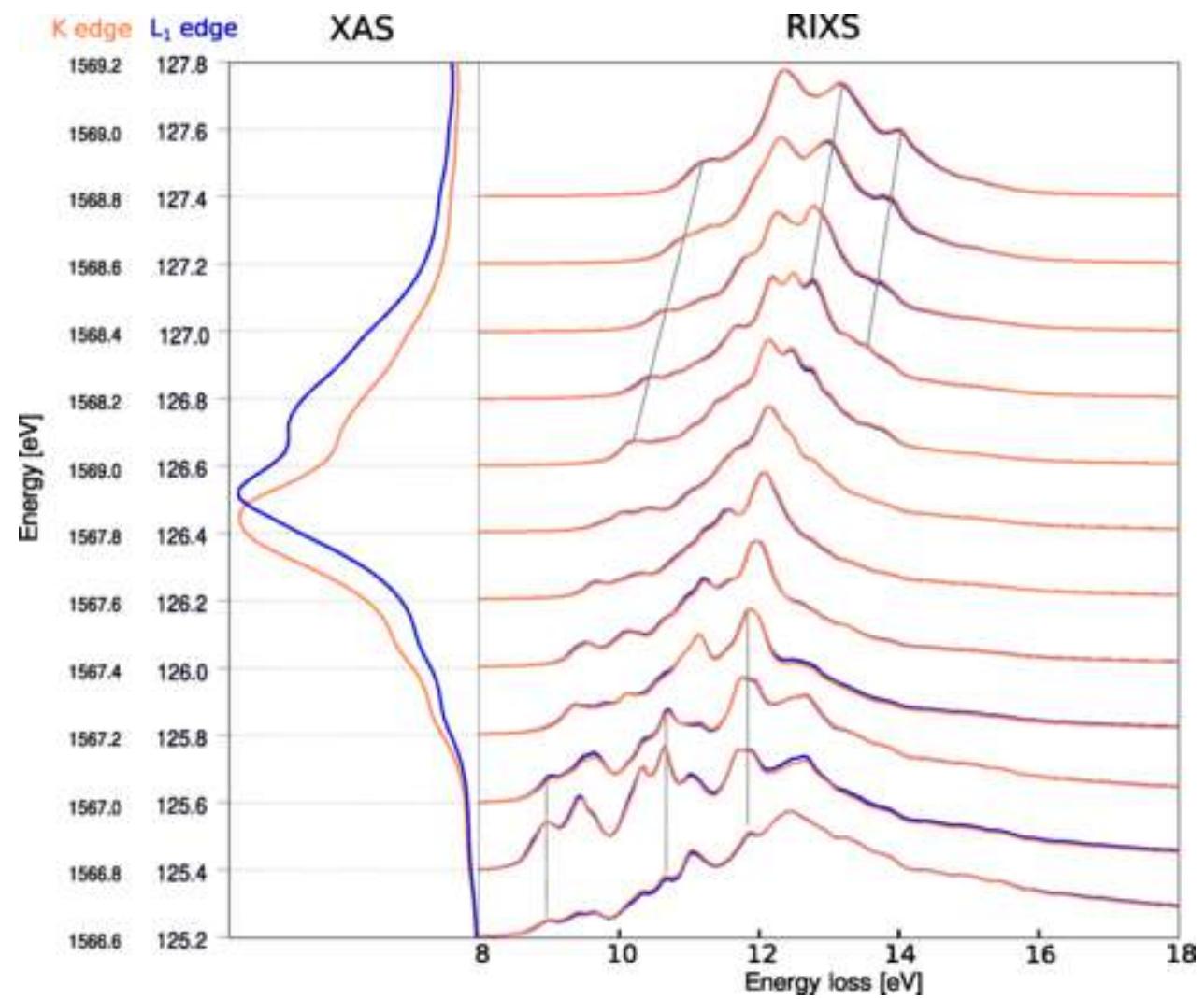


$$\frac{d^2\sigma}{d\Omega_2 d\omega_e} \propto \text{Im} \sum_{\substack{c, c', c'', c''' \\ \mu, \mu', \mu'', \mu'''}} \sum_{\substack{v, v' \\ \mathbf{k}, \mathbf{k}', \mathbf{k}'', \mathbf{k}'''}} \left[\tilde{\rho}_{\mu v \mathbf{k}} \cdot \chi_{c \mu \mathbf{k}}^{c' \mu' \mathbf{k}'}(\omega_i) \cdot \tilde{\rho}_{c' \mu' \mathbf{k}'} \right]^* \chi_{cv \mathbf{k}}^{c'' v' \mathbf{k}''}(\omega) \left[\tilde{\rho}_{\mu'' v' \mathbf{k}''} \cdot \chi_{c'' \mu'' \mathbf{k}''}^{c''' \mu''' \mathbf{k}'''(\omega_i)} \cdot \tilde{\rho}_{c''' \mu''' \mathbf{k}''' \mathbf{k}''' \mathbf{k}''' \mathbf{k}'''}$$

LiF



Al_2O_3



Vorwerk *et al.* Phys. Rev. Research **2**, 042003(R) (2020)



Vorwerk *et al.* Phys. Chem. Chem. Phys. **24**, 17439 (2022).



Urquiza *et al.*, Phys. Rev. B **109**, 115157 (2024)



Urquiza *et al.*, Phys. Rev. B **107**, 205148 (2023)

Spectra and excitons via BSE

- Accurate and extendible
- Rely on the description of the initial (ground) state
- Ab initio and predictive
- Linear response (and beyond) spectroscopies
- Cumbersome calculations

Thanks to the Theoretical Spectroscopy Group



and to You