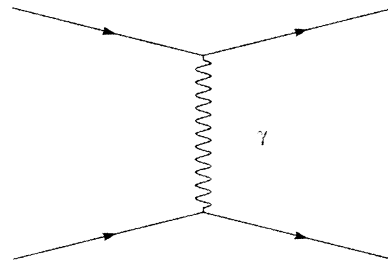


# Introduction to Feynman Diagrams



Igor Reshetnyak



# Before we start...

- Friday 13th
- Programmers' day (256th day of the year)
  - Some numerical applications
  - Less derivations, more discussions

# Plan

- Classical dynamics
  - Trajectories
  - Least action principle
- Quantum mechanics (Feynmann formulation)
  - Functional integral
  - Applications
- Quantum field theory & diagrams
  - What changes?
  - Feynmann diagrams

# Classical Dynamics

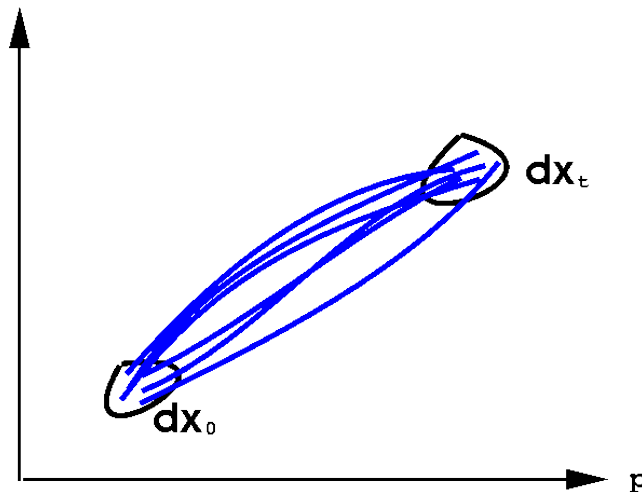
- Deterministic picture
- One can track the full evolution
- Non-trivial things:
  - Initial conditions
  - Instability



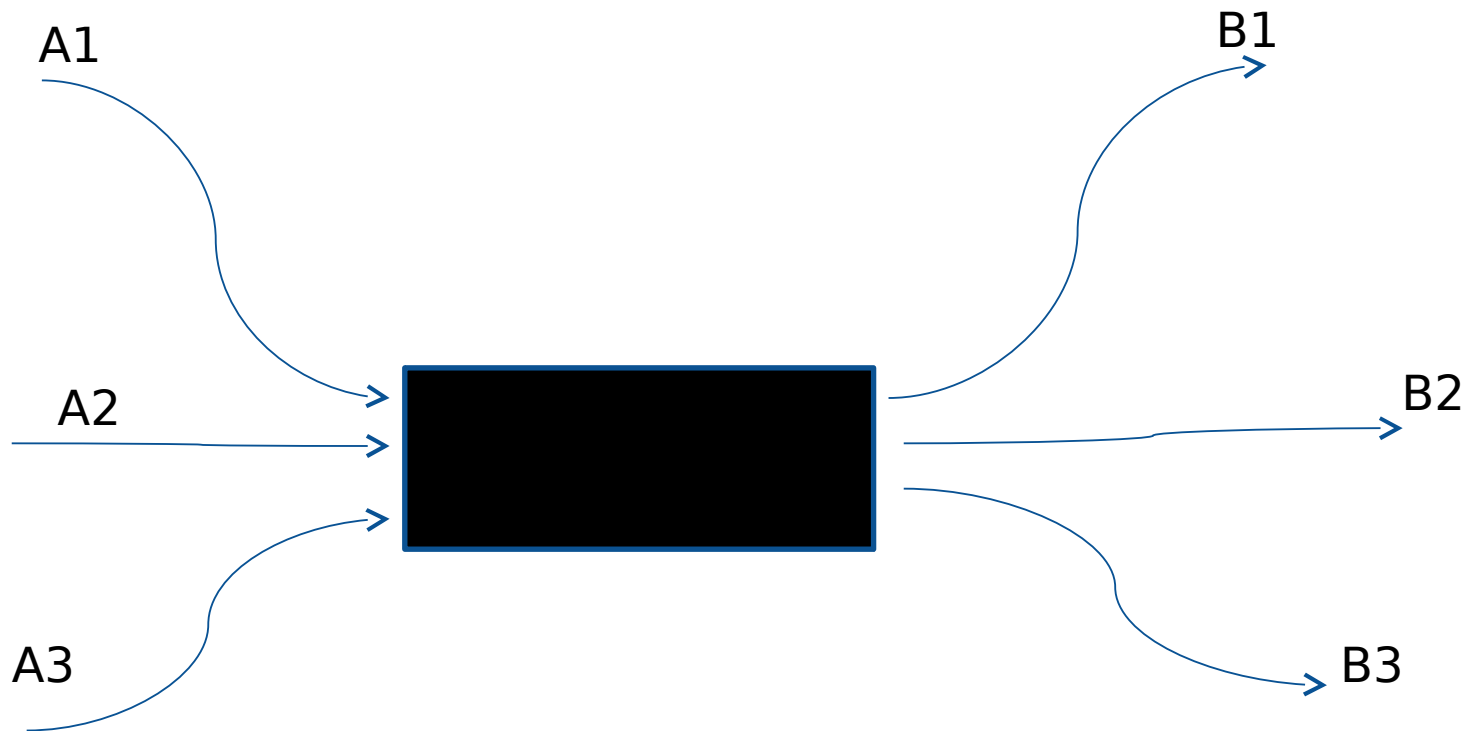
# Classical Dynamics

**Trajectory** - path taken by the system from point A to point B in phase space ( $3N$ )

Liouville Theorem



# Classical Dynamics



Inverse scattering problem!

# Least action principle

- One trajectory
- We know the initial and final state of the system

$$\mathcal{S}[\mathbf{q}(t)] = \int_{t_1}^{t_2} L(\mathbf{q}(t), \dot{\mathbf{q}}(t), t) dt$$

$$\delta\mathcal{S} = 0$$

From optics... Fermat's principle:  
Light takes the shortest path

$$\frac{\partial L}{\partial \mathbf{q}} - \frac{d}{dt} \frac{\partial L}{\partial \dot{\mathbf{q}}} = 0$$


# Classical Statistical Physics

- Paths still deterministic
- Initial conditions (Distribution)
- Nothing really new...

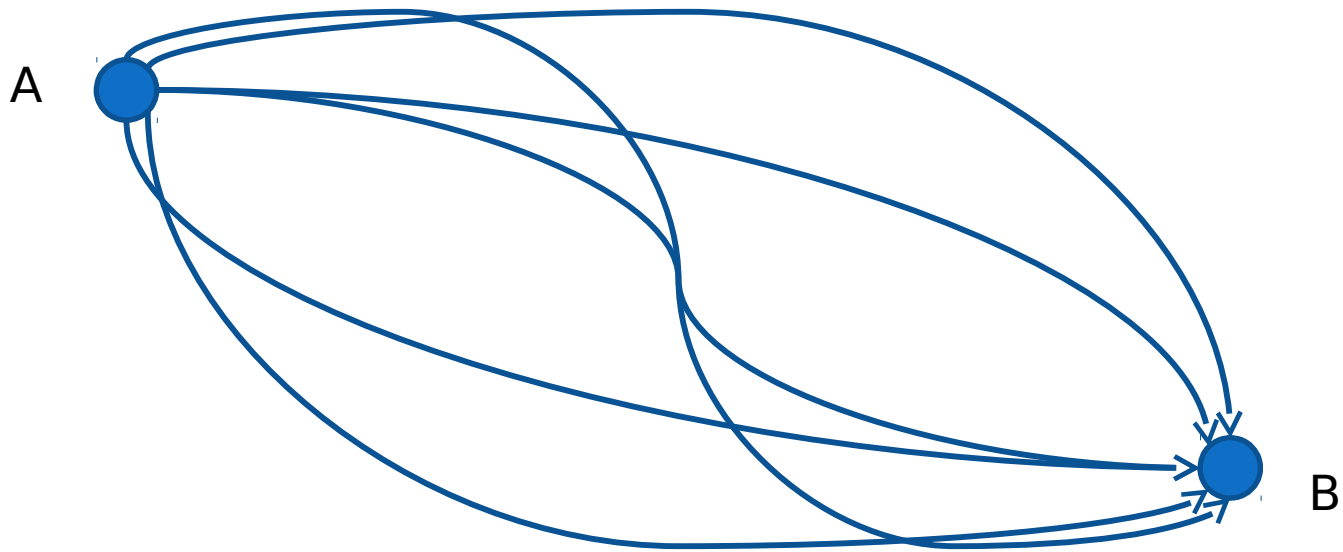


# Quantum Physics

- **Quantum physics is a theory of transformation**

- We CANNOT track full the  path of the system
- Continuous measurement          quantum Zeno effect

# Quantum Physics



Multiple trajectories!

# Quantum Mechanics

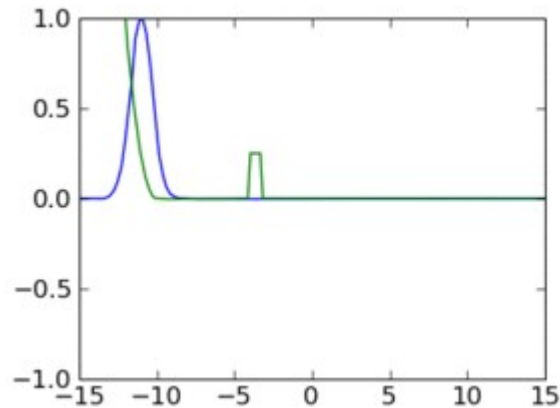
- Multiple trajectories
- Weight  $\sim \exp(i S)$
- **Not always the classical S** e.g.  $L = f(q) (dq/dt)^2$

$$\psi_t(y) = \int \psi_0(x) K(x - y; t) dx = \int \psi_0(x) \int_{x(0)=x}^{x(t)=y} e^{iS} Dx$$

# Application 1

Do things actually work?

$$\exp \left[ -i \Delta t \left( \frac{p^2}{2} + V(x) \right) \right] \simeq \exp \left[ -i V(x) \Delta t / 2 \right] \exp \left[ -i \Delta t \frac{p^2}{2} \right] \exp \left[ -i V(x) \Delta t / 2 \right]$$



# Quantum Statistics

- Rather similar

$$Z = \sum_j \langle j | e^{-\beta H} | j \rangle = \text{Tr} e^{-\beta H}.$$

- Changes:
  - Imaginary time
  - Periodical
  - $L[q, i \, dq/dt]$

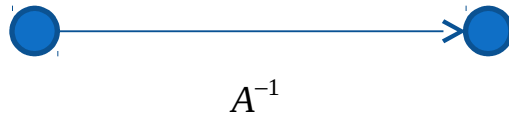
# Field Theory

- Fields i.e. functions of  $x$  and  $t$
- Identical particles...
- But in the end we still have our functional integral...

# Gaussian integrals

$$\int_{-\infty}^{\infty} \exp\left(-\frac{1}{2} \sum_{i,j=1}^n A_{ij} x_i x_j\right) d^n x = \int_{-\infty}^{\infty} \exp\left(-\frac{1}{2} x^T A x\right) d^n x = \sqrt{\frac{(2\pi)^n}{\det A}}$$

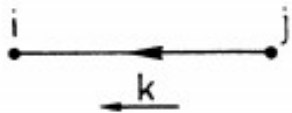
$$\int x^{k_1} \dots x^{k_{2N}} \exp\left(-\frac{1}{2} \sum_{i,j=1}^n A_{ij} x_i x_j\right) d^n x = \sqrt{\frac{(2\pi)^n}{\det A}} \frac{1}{2^N N!} \sum_{\sigma \in S_{2N}} (A^{-1})^{k_{\sigma(1)} k_{\sigma(2)}} \dots (A^{-1})^{k_{\sigma(2N-1)} k_{\sigma(2N)}}$$



# Feynman rules

(Almost the same thing □)

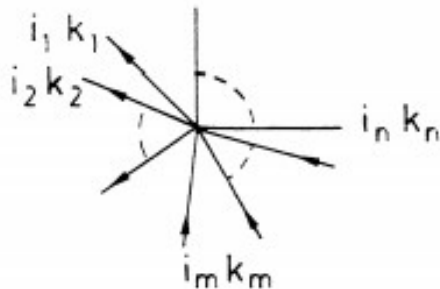
$$\mathcal{L}(x) = \psi_i^*(x) V_{ij} \psi_j(x) + \frac{1}{2} \phi_i(x) W_{ij} \phi_j(x) + \mathcal{L}_I(\psi^*, \psi, \phi)$$



$$\Delta_{Fij}(k) = -\frac{1}{(2\pi)^4 i} \left[ \tilde{V}^{-1}(k) \right]_{ij} ,$$



$$\Delta_{Fij}(k) = -\frac{1}{(2\pi)^4 i} \left[ \frac{1}{2} \tilde{W}(k) + \frac{1}{2} \tilde{W}(-k) \right]_{ij}^{-1} .$$



$$(2\pi)^4 i \sum_{\{1 \dots m-1\}} \sum_{\{m \dots n-1\}} \sum_{\{n \dots\}} (-1)^P \times \\ \times \alpha_{i_1 \dots} (k, k_1, k_2, \dots) \delta_4(k + k_1 + \dots)$$





Thank you for your attention!

# Bibliography

- Main texts:

- Weinberg "Quantum Theory of Fields" (particularly vol 1 chapter 9)
- Abrikosov, Gor'kov, Dzjaloshinskii "Methods of quantum field theory in condensed matter physics"
- Feynmann, Hibbs "Quantum mechanics and path integrals"

- Philosophical things:

- (Russian) <http://mezhpr.fizteh.ru/biblio/q-ivanov/quant-1-0-arphanncajc.pdf>

- Texts online:

- <http://arxiv.org/pdf/quant-ph/0004090v1.pdf>
- <http://www.pa.msu.edu/~pratts/phy831/lectures/lectures.pdf>
- <http://cds.cern.ch/record/186259/files/CERN-73-09.pdf?version=1>
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