

Excitonic Effects within TDDFT

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Outline

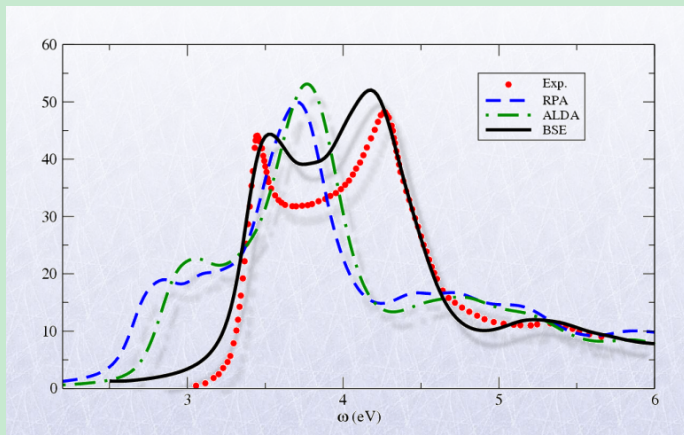
- 1 BSE and TDDFT up to 2002
- 2 The Mapping Theory Kernel
 - Theory
 - Results
- 3 Conclusions and Perspectives

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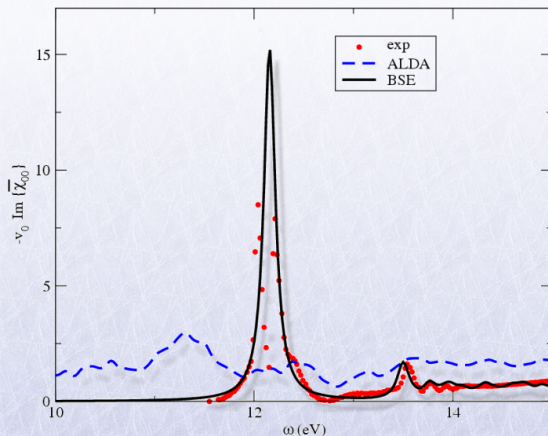
Optical Absorption Spectra of Solids

Semiconductors: Silicon



Optical Absorption Spectra of Solids

Insulators: Argon








Optical Absorption Spectra of Solids

- ALDA bad for any solids!! though quick
- BSE good but cumbersome

Optical Absorption Spectra of Solids






The problem of Abs in solids. Towards a better understanding

-  Reining *et al.* Phys.Rev.Lett. **88**, 66404 (2002)
Long-range kernel
-  de Boeij *et al.* J.Chem.Phys. **115**, 1995 (2002)
Polarization density functional. Long-range.
-  Kim and Görling Phys.Rev.Lett. **89**, 96402 (2002)
Exact-exchange
-  Sottile *et al.* Phys.Rev.B **68**, 205112 (2003)
Long-range and contact exciton.
-  Botti *et al.* Phys. Rev. B **72**, 125203 (2005)
Dynamic long-range component

Parameters to fit to experiments.

Optical Absorption Spectra of Solids

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Parameters to fit to experiments.

Beyond ALDA approximation

Abs in solids. Insights from MBPT

Parameter-free **Ab initio** kernels

 Sottile *et al.* Phys.Rev.Lett. **91**, 56402 (2003)

Full many-body kernel. Mapping Theory.

 Marini *et al.* Phys.Rev.Lett. **91**, 256402 (2003)

Full many-body kernel. Perturbation Theory.

$$f_{xc} = \chi_0^{-1} GGWGG \chi_0^{-1}$$

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The Mapping Theory

The idea

BSE works \Rightarrow $\left\{ \begin{array}{l} \text{we get the ingredients of the BSE} \\ \text{and we put them in TDDFT} \end{array} \right.$

The Mapping Theory

BSE: Excitonic Hamiltonian

4-point

$$H_{(vc)(v'c')}^{\text{BSE}} = \left[(E_c - E_v) \delta_{vv'} \delta_{cc'} + v_{vc}^{v'c'} - W_{vc}^{v'c'} \right]$$

The Mapping Theory

BSE: Excitonic Hamiltonian

4-point

$$H^{\text{BSE}} = \left[(E_c - E_v) + \ll v \gg - \ll W \gg \right]$$

The Mapping Theory

BSE: Excitonic Hamiltonian

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$$H^{\text{BSE}} = \left[\left(\epsilon_c + \Delta_c^{\text{GW}} - \epsilon_v - \Delta_v^{\text{GW}} \right) + \ll v \gg - \ll W \gg \right]$$

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$$H^{\text{BSE}} = \left[\left(\epsilon_c + \Delta_c^{\text{GW}} - \epsilon_v - \Delta_v^{\text{GW}} \right) + \ll v \gg - \ll W \gg \right]$$

TDDFT: Polarizability equation

4-point

$$\chi = \chi_0 + \chi_0 (v + f_{xc}) \chi$$

The Mapping Theory

BSE: Excitonic Hamiltonian

4-point

$$H^{\text{BSE}} = \left[\left(\epsilon_c + \Delta_c^{\text{GW}} - \epsilon_v - \Delta_v^{\text{GW}} \right) + \ll v \gg - \ll W \gg \right]$$

TDDFT: written in transition space

4-point

$$H^{\text{TDDFT}} = \left[(\epsilon_c - \epsilon_v) + \ll v \gg + \ll f_{xc} \gg \right]$$

The Mapping Theory

BSE: Excitonic Hamiltonian

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TDDFT: written in transition space

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$$H^{\text{TDDFT}} = \left[(\epsilon_c - \epsilon_v) + \ll v \gg + \ll f_{xc} \gg \right]$$

The exchange-correlation kernel f_{xc} has to take into account both GW corrections and excitonic effects !!

The Mapping Theory

BSE: Excitonic Hamiltonian

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$$H^{\text{BSE}} = \left[(E_c - E_v) + \ll v \gg - \ll W \gg \right]$$

TDDFT: written in transition space

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$$H^{\text{TDDFT}} = \left[(E_c - E_v) + \ll v \gg + \ll f_{xc} \gg \right]$$

**Same starting point for both BSE and TDDFT:
the GW band-structure.**

The Mapping Theory

BSE: Excitonic Hamiltonian

4-point

$$H^{\text{BSE}} = \left[(E_c - E_v) + \ll v \gg - \ll W \gg \right]$$

TDDFT: written in transition space

4-point

$$H^{\text{TDDFT}} = \left[(E_c - E_v) + \ll v \gg + \ll f_{xc} \gg \right]$$

We concentrate, then, only on the excitonic effects.

The Mapping Theory

BSE: Excitonic Hamiltonian

4-point

$$H^{\text{BSE}} = \left[(E_c - E_v) + \ll v \gg - \ll W \gg \right]$$

TDDFT: written in transition space

4-point

$$H^{\text{TDDFT}} = \left[(E_c - E_v) + \ll v \gg - \ll W \gg \right]$$

We substitute the 'unknown' $\ll f_{xc} \gg$ with $\ll W \gg$.

The Mapping Theory

The idea

We want to use $\ll W \gg$, but in a 2-point equation.

$$\chi(12, \omega) = \chi_0(12, \omega) + \chi_0(13, \omega) (v(34) + f_{xc}(34, \omega)) \chi(42, \omega)$$

The Mapping Theory

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The Mapping Theory

$$\chi = \chi_0 + \chi_0 (\mathbf{v} + \mathbf{f}_{xc}) \chi$$

$$\chi = (\mathbf{1} - \chi_0 \mathbf{v} - \chi_0 \mathbf{f}_{xc})^{-1} \chi_0$$

Let's define an invertible matrix $X(12, \omega) = \sum_{vc} \phi_v(1) \phi_c(1) g_{vc}(2, \omega)$

$$\chi = X X^{-1} (\mathbf{1} - \chi_0 \mathbf{v} - \chi_0 X^{-1} X \mathbf{f}_{xc})^{-1} \chi_0$$

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$$\chi = X (X - \chi_0 v X - \chi_0 X^{-1} T)^{-1} \chi_0$$

$$T(12, \omega) = X(13, \omega) f_{xc}(34, \omega) X(42, \omega) =$$

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$$T(12, \omega) = \sum_{\substack{vc \\ v'c'}} g_{vc}(1, \omega) \ll f_{xc} \gg g_{v'c'}(2, \omega)$$

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The Mapping Theory

TDDFT 2-point equation containing $\ll W \gg$

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$$X(12, \omega) = \sum_{vc} \phi_v(1) \phi_c(1) g_{vc}(2, \omega)$$

What about the application ??

The Mapping Theory

TDDFT 2-point equation containing $\ll W \gg$

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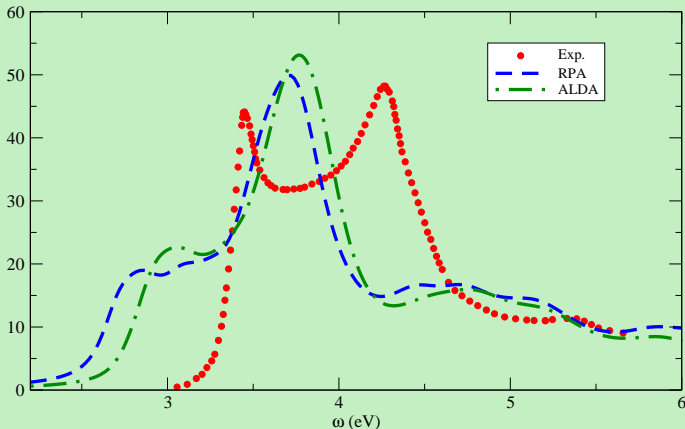
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The Mapping Theory: Results

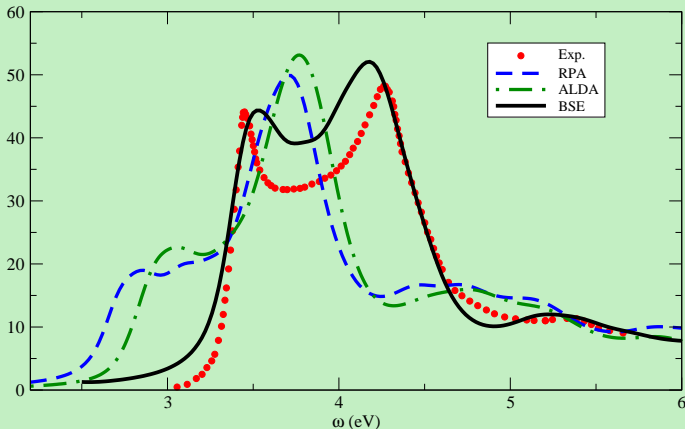
Absorption of Silicon





The Mapping Theory: Results

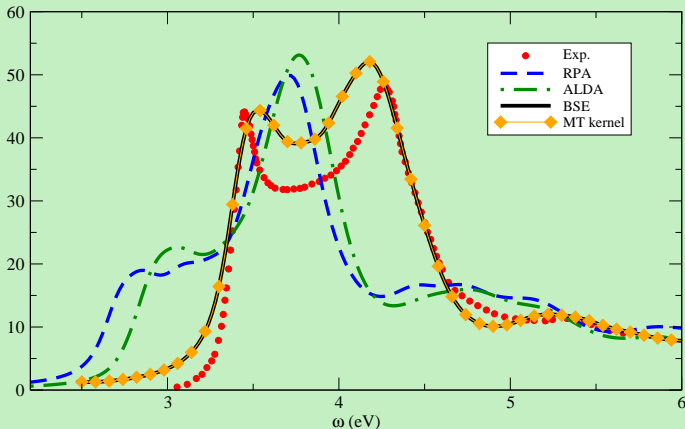
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The Mapping Theory: Results

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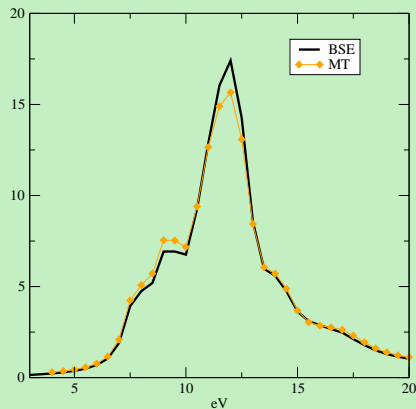
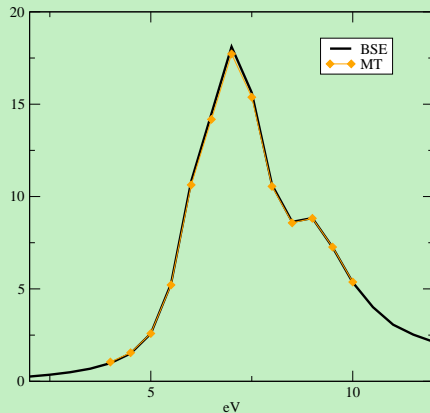


F.Sottile *et al.* Phys.Rev.Lett **91**, 56402 (2003)



The Mapping Theory: Results

Absorption of Silicon Carbide and Diamond

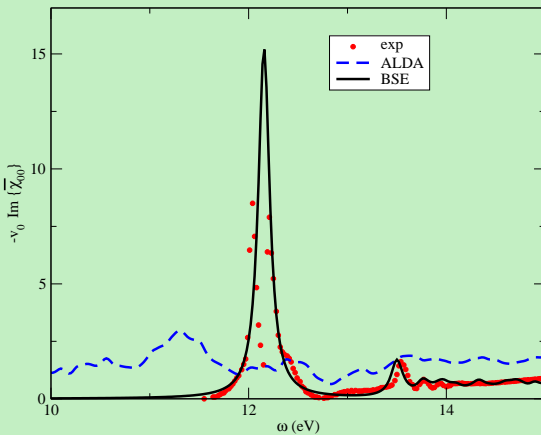


last week preliminary results :-)



The Mapping Theory: Results

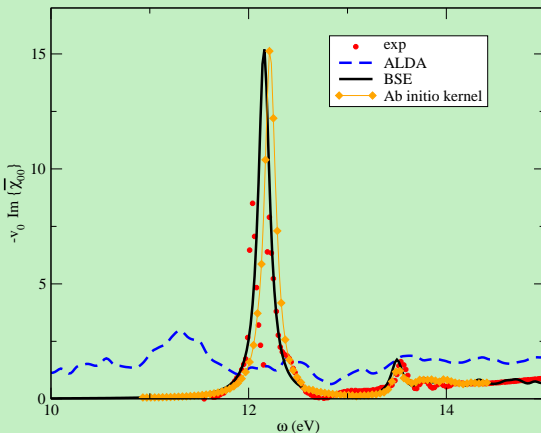
Absorption of Argon





The Mapping Theory: Results

Absorption of Argon







F.Sottile *et al.*, submitted to Phys.Rev.Lett.



The Mapping Theory: Results

Tested also on absorption of SiO_2 , DNA bases, Ge-nanowires, RAS of diamond surface, and EELS of LiF.

-  Marini *et al.* Phys.Rev.Lett. **91**, 256402 (2003).
-  Bruno *et al.* Phys.Rev.B **72** 153310, (2005).
-  Palumbo *et al.* Phys.Rev.Lett. **94** 087404 (2005).
-  Varsano *et al.* J.Phys.Chem.B **110** 7129 (2006).

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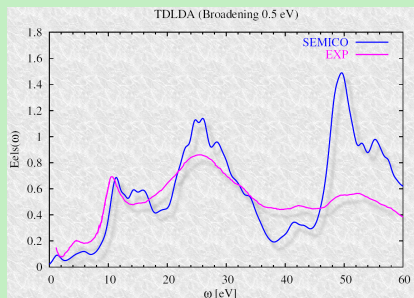
Conclusions

TDDFT is the method of choice

- ✓ Absorption spectra of simple molecules
- ✓ Electron energy loss spectra
- ✓ Inelastic X-ray scattering spectroscopy
- ✓ Absorption of Solids (BSE-like scaling)

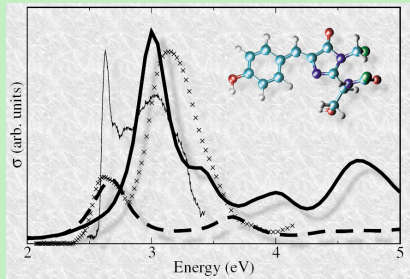
Towards new applications

Strongly correlated systems

EEL spectrum of VO₂

M.Gatti, preliminary results

Biological systems

Abs spectrum of Green
Fluorescent ProteinM.Marques *et al.* Phys.Rev.Lett
90, 258101 (2003)

New Frontiers

Excited-State Dynamics

TDDFT-MD, Ehrenfest dynamics, quantum effects of the ions, non-adiabaticity, etc.



 Sugino and Miyamoto, Phys.Rev.B **59**, 2579 (1999)

New Frontiers

TDDFT concept into MBPT



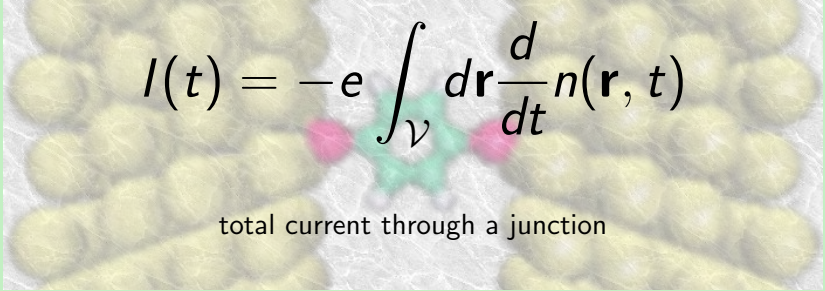
$$\Sigma = GW\Gamma$$

i.e. a promising path to go beyond GW approx through TDDFT

 F.Bruneval *et al.* Phys.Rev.Lett **94**, 186402 (2005)

New Frontiers

Quantum Transport in TDDFT


$$I(t) = -e \int_V d\mathbf{r} \frac{d}{dt} n(\mathbf{r}, t)$$

total current through a junction

 G.Stefanucci *et al.* *Europhys.Lett.* **67**, 14 (2004)

New Frontiers

Let's go back to Ground-State

Total energies calculations via TDDFT

$$E = T_{KS} + V_{ext} + E_H + E_{xc}$$

$$E_{xc} \propto \int d\mathbf{r} d\mathbf{r}' \int_0^1 d\lambda \int_0^\infty du \chi^\lambda(\mathbf{r}, \mathbf{r}', iu)$$

adiabatic connection fluctuation-dissipation theorem



D.C.Langreth *et al.* Solid State Comm. **17**, 1425 (1975)



M.Lein *et al.* **61**, 13431 (2000)

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- ✓ Absorption of Solids (BSE-like scaling)

- ✓ Promising for new applications.

Open problems

- open-shell atoms
- charge-transfer excitations
- *really* efficient calculations of solids
approximation for f_{xc}

Conclusions

TDDFT is the method of choice

- ✓ Absorption spectra of simple molecules
- ✓ Electron energy loss spectra
- ✓ Inelastic X-ray scattering spectroscopy
- ✓ Absorption of Solids (BSE-like scaling)

- ✓ Promising for new applications.

Open problems

- open-shell atoms
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 - *really* efficient calculations of solids
- approximation for f_{xc}**

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**Thank you for your attention
during these three days**

