Analysis

Time Dependent Density Functional Theory

Introduction and Applications

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Tegernsee, 22 July 2009







- Motivations
- Linear Response Approach





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Motivations

Linear Response Approach



Analysis

Why Density Functional

$$\Psi(\mathbf{r}_1, \mathbf{r}_2, ..., \mathbf{r}_N)$$

$$\downarrow$$

$$G(\mathbf{r}_1, \mathbf{r}_2)$$

$$\downarrow$$

$$\rho(\mathbf{r})$$

Time Dependent Density Functional Theory

Why Density Functional: an old strategy

$$\Psi(\mathbf{r}_1, \mathbf{r}_2, ..., \mathbf{r}_N)$$

$$\downarrow$$

$$G(\mathbf{r}_1, \mathbf{r}_2)$$

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$$\rho(\mathbf{r})$$



Density Functional ... Successfull ?

TABLE I: Physical Review articles with more than 1000 citations through June 2003. PR, Physical Review; PRB, Physical Review B; PRD, Physical Review D; PRL, Physical Review Letters; RMP, Reviews of Modern Physics.



	#	Av.		
Publication	cites	Age	Title	Author(s)
PR 140, A1133 (1965)	3227	26.7	Self-Consistent Equations Including Exchange and	W. Kohn, L. J. Sham
			Correlation Effects	
PR 136 , B864 (1964)	2460	28.7	Inhomogeneous Electron Gas	P. Hohenberg, W. Kohn
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S. Redner http://arxiv.org/abs/physics/0407137

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Large field of research concerned with many-electron systems in time-dependent fields

- absorption spectra
- energy loss spectra
- X scattering
- high-harmonic generation
- photo-emission

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System submitted to an external perturbation







Time Dependent Density Functional Theory

System submitted to an external perturbation



$$egin{aligned} V_{tot} &= arepsilon^{-1} V_{ext} \ V_{tot} &= V_{ext} + V_{inc} \end{aligned}$$

$$\mathbf{E} = \varepsilon^{-1} \mathbf{D}$$



Time Dependent Density Functional Theory



$$V_{tot} = arepsilon^{-1} V_{ext}$$

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Definition of polarizability

 $\begin{array}{rcl} \text{not polarizable} & \Rightarrow & V_{tot} = V_{\text{ext}} & \Rightarrow & \varepsilon^{-1} = 1 \\ polarizable & \Rightarrow & V_{tot} \neq V_{\text{ext}} & \Rightarrow & \varepsilon^{-1} \neq 1 \end{array}$

Time Dependent Density Functional Theory

Francesco Sottile

Definition of polarizability

not polarizable \Rightarrow polarizable

 \Rightarrow

$$V_{tot} = V_{ext} \implies \varepsilon^{-1} = 1$$
$$V_{tot} \neq V_{ext} \implies \varepsilon^{-1} \neq 1$$

Definition of polarizability

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$$egin{array}{lll} V_{tot} = V_{ext} &\Rightarrow & arepsilon^{-1} = 1 \ V_{tot}
eq V_{ext} &\Rightarrow & arepsilon^{-1}
eq 1 \ arepsilon^{-1} = 1 + v \chi \end{array}$$

χ is the polarizability of the system

Definition of polarizability

 \Rightarrow

$$V_{tot} = V_{ext} \Rightarrow \varepsilon^{-1} = 1$$

$$V_{tot} \neq V_{ext} \Rightarrow \varepsilon^{-1} \neq 1$$

$$\varepsilon^{-1} = 1 + 1$$

χ is the polarizability of the system

Vχ

Polarizability

interacting system $\delta n = \chi \delta V_{ext}$ non-interacting system $\delta n_{n-i} = \chi^0 \delta V_{tot}$

Polarizability

interacting system $\delta n = \chi \delta V_{ext}$ non-interacting system $\delta n_{n-i} = \chi^0 \delta V_{tot}$ Single-particle polarizability

$$\chi^{0} = \sum_{ij} \frac{\phi_{i}(\mathbf{r})\phi_{j}^{*}(\mathbf{r})\phi_{i}^{*}(\mathbf{r}')\phi_{j}(\mathbf{r}')}{\omega - (\epsilon_{i} - \epsilon_{j})}$$

hartree, hartree-fock, dft, etc.

🦠 G.D. Mahan Many Particle Physics (Plenum, New York, 1990)

Polarizability



Polarizability

interacting system $\delta n = \chi \delta V_{ext}$ non-interacting system $\delta n_{n-i} = \chi^0 \delta V_{tot}$ \uparrow

Density Functional Formalism

 $\delta n = \delta n_{n-i}$

$$\delta V_{tot} = \delta V_{ext} + \delta V_{H} + \delta V_{xc}$$

Polarizability

$$\chi \delta V_{\text{ext}} = \chi^0 \left(\delta V_{\text{ext}} + \delta V_H + \delta V_{\text{xc}} \right)$$
$$\chi = \chi^0 \left(1 + \frac{\delta V_H}{\delta V_{\text{ext}}} + \frac{\delta V_{\text{xc}}}{\delta V_{\text{ext}}} \right)$$
$$\frac{\delta V_H}{\delta V_{\text{ext}}} = \frac{\delta V_H}{\delta n} \frac{\delta n}{\delta V_{\text{ext}}} = v\chi$$
$$\frac{\delta V_{\text{xc}}}{\delta V_{\text{ext}}} = \frac{\delta V_{xc}}{\delta n} \frac{\delta n}{\delta V_{\text{ext}}} = f_{\text{xc}}\chi$$

with **f_{xc} = exchange-correlation kernel**

Time Dependent Density Functional Theory

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$$\chi = \chi^{0} + \chi^{0} \left(\mathbf{v} + f_{xc} \right) \chi$$

with **f_{xc} = exchange-correlation kernel**

Polarizability

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$$\chi = \left[1 - \chi^{0} \left(v + f_{xc} \right) \right]^{-1} \chi^{0}$$
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Polarizability

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Polarizability χ in TDDFT

• DFT ground-state calc. $\rightarrow \phi_i, \epsilon_i \quad [V_{xc}]$ • $\phi_i, \epsilon_i \quad \rightarrow \quad \chi^0 = \sum_{ij} \frac{\phi_i(\mathbf{r})\phi_j^*(\mathbf{r})\phi_i(\mathbf{r}')\phi_j(\mathbf{r}')}{\omega - (\epsilon_i - \epsilon_j)}$ • $\frac{\delta V_{ii}}{\delta n} = v$ • $\frac{\delta V_{xc}}{\delta n} = f_{xc}$ variation of the potentials • $\chi = \chi^0 + \chi^0 (v + f_{xc}) \chi$

Polarizability χ in TDDFT

- DFT ground-state calc. $\rightarrow \phi_i, \epsilon_i \quad [V_{xc}]$
- $\phi_{i}, \epsilon_{i} \rightarrow \chi^{0} = \sum_{ij} \frac{\varphi_{i}(\epsilon_{j}, \varphi_{j}(\epsilon_{j}), \varphi_{j}(\epsilon_{j}), \varphi_{j}(\epsilon_{j}))}{\omega (\epsilon_{i} \epsilon_{j})}$
 - variation of the potentials
- $\chi = \chi^0 + \chi^0 (v + f_{xc}) \chi$

Polarizability χ in TDDFT

- **1** DFT ground-state calc. $\rightarrow \phi_i, \epsilon_i \quad [V_{xc}]$ **2** $\phi_i, \epsilon_i \quad \rightarrow \quad \chi^0 = \sum_{ij} \frac{\phi_i(\mathbf{r})\phi_j^*(\mathbf{r})\phi_i^*(\mathbf{r}')\phi_j(\mathbf{r}')}{\omega (\epsilon_i \epsilon_j)}$ **3** $\frac{\delta V_H}{\delta \alpha} = V$ **4** $\frac{\delta V_{xc}}{\delta \alpha} = f_{xc}$ **4** variation of the potentials
- $\chi = \chi^{0} + \chi^{0} (v + f_{xc}) \chi$

Polarizability χ in TDDFT

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Analysis

Theoretical Spectroscopy

$$\chi(\mathbf{r},\mathbf{r}',\omega) \to \chi_{\mathbf{c},\mathbf{c}'}(\mathbf{q},\omega) \to \varepsilon_{00}^{-1}(\mathbf{q},\omega)$$

Energy Loss Function = $-\text{Im}\left\{\varepsilon_{00}^{-1}(\mathbf{q},\omega)\right\}$

Absorption = Im
$$\left\{ \frac{1}{\varepsilon_{00}^{-1}(\mathbf{q},\omega)} \right\}$$

🚺 TDDFT

- Motivations
- Linear Response Approach

2 Applications (ELS)

3 Analysis

EELS of Graphite

Some good results ... (graphite)



🚺 A.Marinopoulos *et al.* Phys.Rev.Lett **89**, 76402 (2002)

Time Dependent Density Functional Theory

Inelastic X-ray Scattering

TD-LDA on IXS of Silicon



Inelastic X-ray Scattering

TD-LDA on IXS of Silicon



H-C. Weissker et al., Physical Review Letters 97, 237602 (2006)

Semi-core states

L-edge of Silicon



Luppi *et al.* Phys. Rev. B **78**, 245124 (2008)

EELS of LiF : many-body effects (beyond TDLDA)

$\mathbf{q} = 0.5\Gamma X$



A.Marini *et al.*, PRL **91**, 256402 (2003).

Prediction

ELS of Hafnium Oxide



Zobelli and Sottile, work in progress.

TDDFT

- Motivations
- Linear Response Approach



Why the numerical approach is important



Analysis

ELS of Hafnium Oxide



Zobelli and Sottile, work in progress

Analysis

ELS of Nanotubes via Graphene analysis



Kramberger *et al.*, Phys. Rev. Lett. **100**, 196803 (2008)

EELS of nanotubes: plasmon dispersion



Kramberger *et al.*, Phys. Rev. Lett. **100**, 196803 (2008)

EELS of nanotubes: plasmon dispersion



Analysis

Graphic tools: 'see' the plasmons or the excitons





E=9eV

E=30eV



ELNES of BN

Nitrogen Edge





ELNES of Cu





Numerical simulations

ab-initio calculations

- DFT ground-state calculations (LDA)
- Independant Particles polarizability: χ^0
- RPA Full polarisability: $\chi = \left[1 \chi^0 v\right]^{-1} \chi^0$
- Dielectric function $\varepsilon^{-1} = 1 + v\chi$
- energy loss function $-Im\{\varepsilon^{-1}(\mathbf{q},\omega)\}$



Analysis

Independent particle picture



 \implies given by χ^0 : interpretation in terms of **band-transitions**



Independent particle picture



Μ

Analysis

RPA: random phase approx.



- given by χ:
 no interpretation by
 band-transitions
- contributions from K
- mixing of transitions

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RPA: random phase approx.



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Plasmon dispersion













