

# Introduction to Green's functions

Matteo Gatti

ETSF Users' Meeting and Training Day  
Ecole Polytechnique - 22 October 2010

# Outline

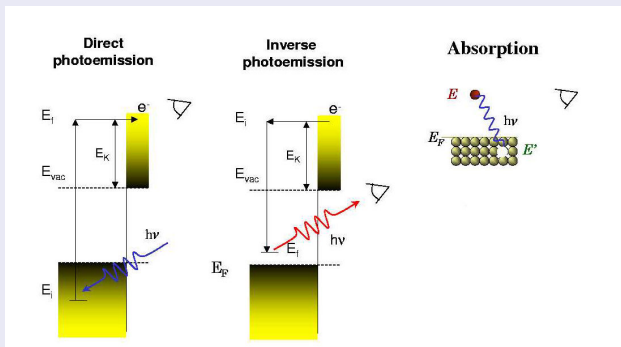
- 1 Motivation
- 2 Green's functions
- 3 The GW Approximation
- 4 The Bethe-Salpeter Equation

# Outline

- 1 Motivation
- 2 Green's functions
- 3 The GW Approximation
- 4 The Bethe-Salpeter Equation

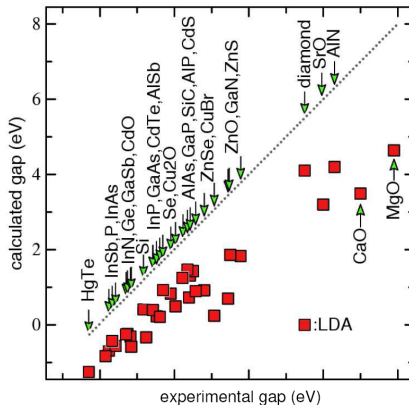
# Electronic Spectroscopy

Here only two categories:



- 1 charged excitations: photoemission and inverse photoemission
- 2 neutral excitations: absorption and electron energy loss

# Why do we have to study more than DFT?



adapted from M. van Schilfgaarde *et al.*, PRL **96** (2006).

# MBPT vs. TDDFT: different worlds, same physics

## (TD)DFT

- based on the density
- response function  $\chi$ : neutral excitations
- moves density around
- is efficient (simple)

## MBPT

- based on Green's functions
- one-particle  $G$ : electron addition and removal - GW  
two-particle  $L$ : electron-hole excitation - BSE
- moves (quasi)particles around
- is intuitive (easy)

# Table of characters

- Density: local in space and time  
 $\rho(\mathbf{r}_1, t_1)$
- Density matrix: non-local in space  
 $\gamma(\mathbf{r}_1, \mathbf{r}_2, t)$
- One-particle Green's function: non-local in space and time:  
 $G(\mathbf{r}_1, \mathbf{r}_2, t_1, t_2)$

- $G(1, 2) \equiv G(\mathbf{r}_1, \mathbf{r}_2, t_1, t_2)$
- $G(\mathbf{r}_1, \mathbf{r}_2, t_1, t_2) = G(\mathbf{r}_1, \mathbf{r}_2, t_1 - t_2) \Rightarrow G(\mathbf{r}_1, \mathbf{r}_2, \omega)$
- $\rho(\mathbf{r}_1, t_1) = -iG(\mathbf{r}_1, \mathbf{r}_1, t_1, t_1^+)$

# Table of characters

- Density: local in space and time

$$\rho(\mathbf{r}_1, t_1)$$

- Density matrix: non-local in space

$$\gamma(\mathbf{r}_1, \mathbf{r}_2, t)$$

- One-particle Green's function: non-local in space and time:

$$G(\mathbf{r}_1, \mathbf{r}_2, t_1, t_2)$$

- $G(1, 2) \equiv G(\mathbf{r}_1, \mathbf{r}_2, t_1, t_2)$

- $G(\mathbf{r}_1, \mathbf{r}_2, t_1, t_2) = G(\mathbf{r}_1, \mathbf{r}_2, t_1 - t_2) \Rightarrow G(\mathbf{r}_1, \mathbf{r}_2, \omega)$

- $\rho(\mathbf{r}_1, t_1) = -iG(\mathbf{r}_1, \mathbf{r}_1, t_1, t_1^+)$

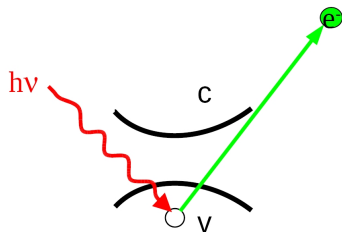


# Outline

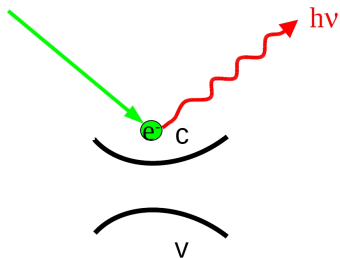
- 1 Motivation
- 2 Green's functions**
- 3 The GW Approximation
- 4 The Bethe-Salpeter Equation

# Photoemission

## Direct Photoemission



## Inverse Photoemission



# One-particle Green's function

## The one-particle Green's function $G$

Definition and meaning of  $G$ :

$$iG(\mathbf{r}_1, t_1; \mathbf{r}_2, t_2) = \langle N | T [\psi(\mathbf{r}_1, t_1) \psi^\dagger(\mathbf{r}_2, t_2)] | N \rangle$$

$$\text{for } t_1 > t_2 \Rightarrow iG(\mathbf{r}_1, t_1; \mathbf{r}_2, t_2) = \langle N | \psi(\mathbf{r}_1, t_1) \psi^\dagger(\mathbf{r}_2, t_2) | N \rangle \quad (1)$$

$$\text{for } t_1 < t_2 \Rightarrow iG(\mathbf{r}_1, t_1; \mathbf{r}_2, t_2) = -\langle N | \psi^\dagger(\mathbf{r}_2, t_2) \psi(\mathbf{r}_1, t_1) | N \rangle \quad (2)$$

# One-particle Green's function

## The one-particle Green's function $G$

Definition and meaning of  $G$ :

$$iG(\mathbf{r}_1, t_1; \mathbf{r}_2, t_2) = \langle N | T [\psi(\mathbf{r}_1, t_1) \psi^\dagger(\mathbf{r}_2, t_2)] | N \rangle$$

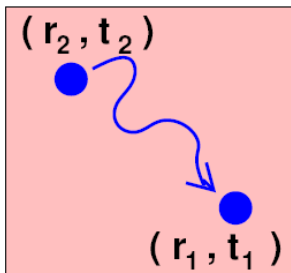
$$\text{for } t_1 > t_2 \Rightarrow iG(\mathbf{r}_1, t_1; \mathbf{r}_2, t_2) = \langle N | \psi(\mathbf{r}_1, t_1) \psi^\dagger(\mathbf{r}_2, t_2) | N \rangle \quad (1)$$

$$\text{for } t_1 < t_2 \Rightarrow iG(\mathbf{r}_1, t_1; \mathbf{r}_2, t_2) = -\langle N | \psi^\dagger(\mathbf{r}_2, t_2) \psi(\mathbf{r}_1, t_1) | N \rangle \quad (2)$$

# One-particle Green's function

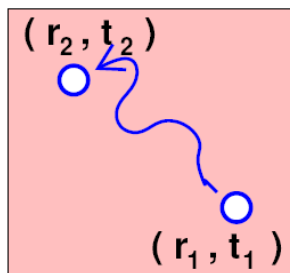
$$t_1 > t_2$$

$$\langle N | \psi(\mathbf{r}_1, t_1) \psi^\dagger(\mathbf{r}_2, t_2) | N \rangle$$



$$t_1 < t_2$$

$$-\langle N | \psi^\dagger(\mathbf{r}_2, t_2) \psi(\mathbf{r}_1, t_1) | N \rangle$$



# One-particle Green's function

## What is $G$ ?

Definition and meaning of  $G$ :

$$G(\mathbf{r}_1, t_1; \mathbf{r}_2, t_2) = -i \langle N | T \left[ \psi(\mathbf{r}_1, t_1) \psi^\dagger(\mathbf{r}_2, t_2) \right] | N \rangle$$

Insert a complete set of  $N + 1$  or  $N - 1$ -particle states. This yields

$$G(\mathbf{r}_1, t_1; \mathbf{r}_2, t_2) = -i \sum_j f_j(\mathbf{r}_1) f_j^*(\mathbf{r}_2) e^{-i\varepsilon_j(t_1 - t_2)} \times \\ \times [\theta(t_1 - t_2) \theta(\varepsilon_j - \mu) - \theta(t_2 - t_1) \Theta(\mu - \varepsilon_j)];$$

where:

$$\varepsilon_j = \begin{cases} E(N + 1, j) - E(N), & \varepsilon_j > \mu \\ E(N) - E(N - 1, j), & \varepsilon_j < \mu \end{cases}$$

$$f_j(\mathbf{r}_1) = \begin{cases} \langle N | \psi(\mathbf{r}_1) | N + 1, j \rangle, & \varepsilon_j > \mu \\ \langle N - 1, j | \psi(\mathbf{r}_1) | N \rangle, & \varepsilon_j < \mu \end{cases}$$

# One-particle Green's function

## What is G? - Fourier transform

Fourier Transform:

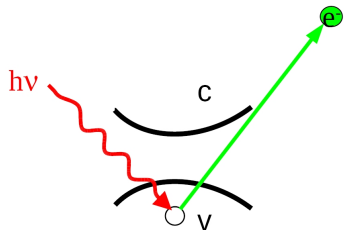
$$G(\mathbf{x}, \mathbf{x}', \omega) = \sum_j \frac{f_j(\mathbf{x})f_j^*(\mathbf{x}')}{\omega - \varepsilon_j + i\eta \operatorname{sgn}(\varepsilon_j - \mu)}.$$

Spectral function:

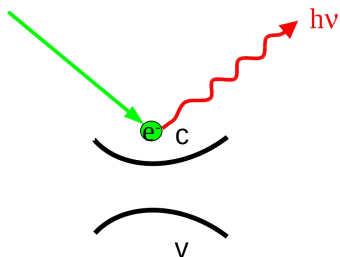
$$A(\mathbf{x}, \mathbf{x}'; \omega) = \frac{1}{\pi} |\operatorname{Im}G(\mathbf{x}, \mathbf{x}'; \omega)| = \sum_j f_j(\mathbf{x})f_j^*(\mathbf{x}')\delta(\omega - \varepsilon_j).$$

# Photoemission

## Direct Photoemission



## Inverse Photoemission



One-particle excitations  $\rightarrow$  poles of one-particle Green's function  $G$



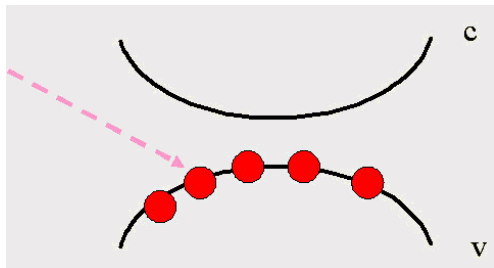
# One-particle Green's function

## One-particle Green's function

From one-particle  $G$  we can obtain:

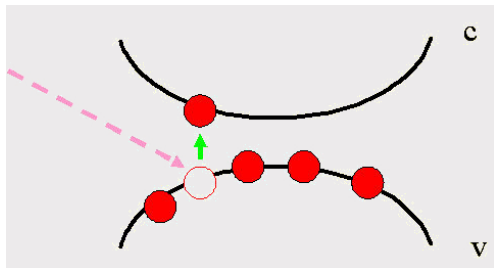
- one-particle excitation spectra
- ground-state expectation value of any one-particle operator:  
e.g. density  $\rho$  or density matrix  $\gamma$ :  
$$\rho(\mathbf{r}, t) = -iG(\mathbf{r}, \mathbf{r}, t, t^+) \quad \gamma(\mathbf{r}, \mathbf{r}', t) = -iG(\mathbf{r}, \mathbf{r}', t, t^+)$$
- ground-state total energy

# Absorption



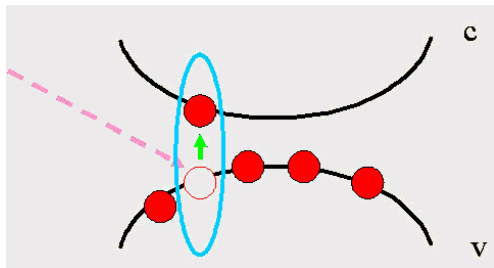
Two-particle excitations  $\rightarrow$  poles of two-particle Green's function  $L$   
Excitonic effects = electron - hole interaction

# Absorption



Two-particle excitations  $\rightarrow$  poles of two-particle Green's function  $L$   
Excitonic effects = electron - hole interaction

# Absorption



Two-particle excitations  $\rightarrow$  poles of two-particle Green's function  $L$   
Excitonic effects = electron - hole interaction

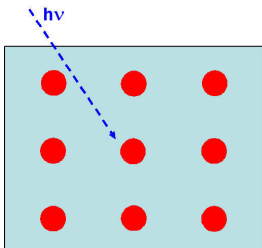
# Table of characters

- $G(1, 2)$ : one-particle Green's function (2 points)
- $L(1, 2, 3, 4)$ : two-particle Green's function (4 points)

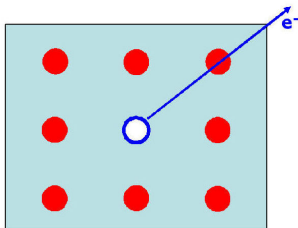
# Outline

- 1 Motivation
- 2 Green's functions
- 3 The GW Approximation**
- 4 The Bethe-Salpeter Equation

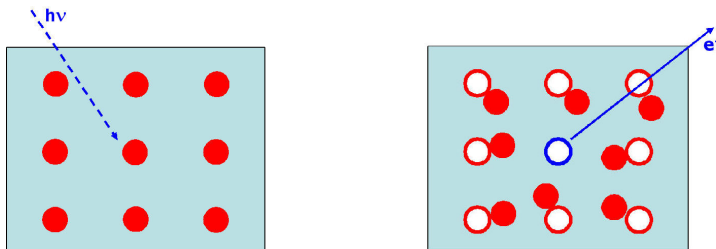
# GW bandstructure: photoemission



additional charge  $\rightarrow$



# GW bandstructure: photoemission



additional charge  $\rightarrow$  reaction: polarization, screening

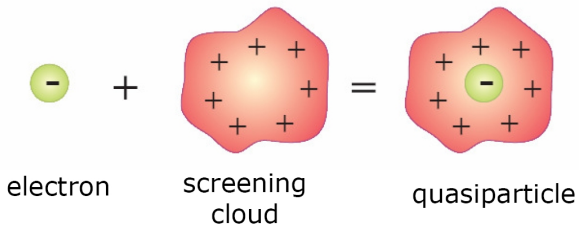
## GW approximation

- 1 polarization made of noninteracting electron-hole pairs (RPA)
- 2 classical (Hartree) interaction between additional charge and polarization charge



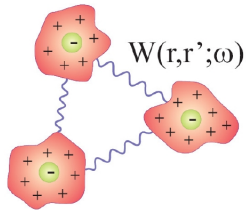
# GW and Hartree-Fock

## Quasiparticle



# GW and Hartree-Fock

## Screened potential $W$



$W$  = screened potential:  
weaker than bare Coulomb interaction

$$W(r, r', \omega) = \int dr'' \frac{\epsilon^{-1}(r, r'', \omega)}{|r'' - r'|}$$

# GW and Hartree-Fock

## Hartree-Fock

$$\Sigma(12) = iG(12)v(1+2)$$

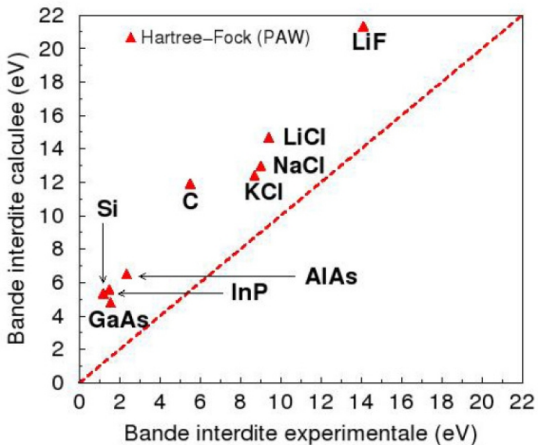
- $v$  infinite range in space
- $v$  is static
- $\Sigma$  is nonlocal, hermitian, static

## GW

$$\Sigma(12) = iG(12)W(1+2)$$

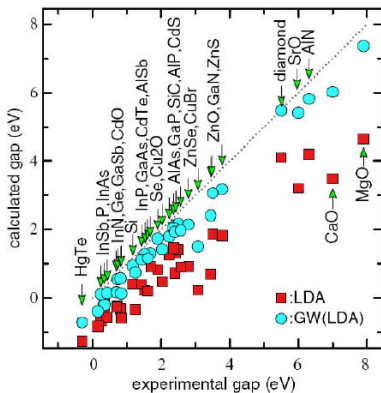
- $W$  is short ranged
- $W$  is dynamical
- $\Sigma$  is nonlocal, complex, dynamical

# GW and Hartree-Fock



from Brice Arnaud

# GW and Hartree-Fock



M. van Schilfgaarde *et al.*, PRL **96** (2006).

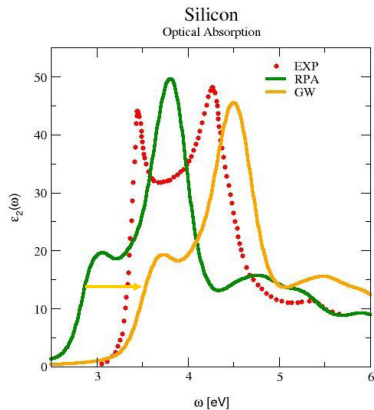
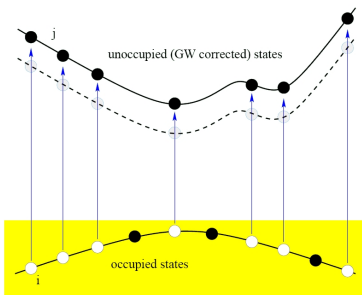
# Outline

- 1 Motivation
- 2 Green's functions
- 3 The GW Approximation
- 4 The Bethe-Salpeter Equation**

# Independent (quasi)particles: GW

Independent transitions:

$$\epsilon_2(\omega) = \frac{8\pi^2}{\Omega\omega^2} \sum_{ij} |\langle \varphi_j | \mathbf{e} \cdot \mathbf{v} | \varphi_i \rangle|^2 \delta(E_j - E_i - \omega)$$

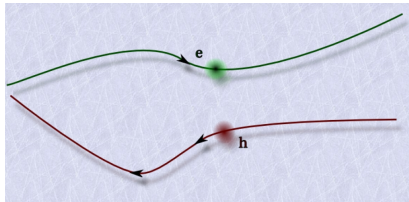


What is wrong?

What is missing?

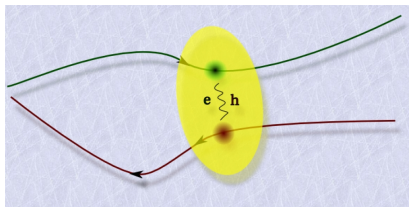


# Beyond RPA



Independent particles (RPA)

# Beyond RPA



Interacting particles (excitonic effects)

# Absorption spectra in BSE

## Independent (quasi)particles

$$Abs(\omega) \propto \sum_{vc} |\langle v|D|c \rangle|^2 \delta(E_c - E_v - \omega)$$

## Excitonic effects: the Bethe-Salpeter equation

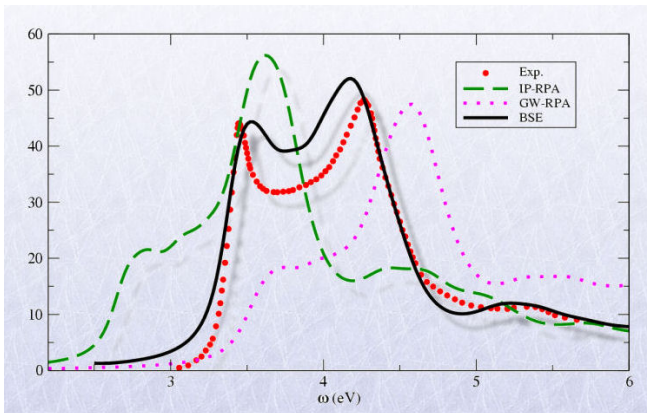
$$[H_{el} + H_{hole} + H_{el-hole}]A_\lambda = E_\lambda A_\lambda$$

$$Abs(\omega) \propto \sum_\lambda \left| \sum_{vc} A_\lambda^{(vc)} \langle v|D|c \rangle \right|^2 \delta(E_\lambda - \omega)$$

- mixing of transitions:  $|\langle v|D|c \rangle|^2 \rightarrow \left| \sum_{vc} A_\lambda^{(vc)} \langle v|D|c \rangle \right|^2$
- modification of excitation energies:  $E_c - E_v \rightarrow E_\lambda$

# Absorption spectra in BSE

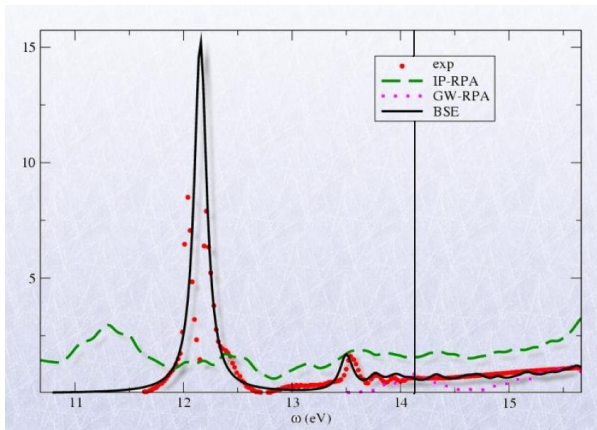
## Bulk silicon



G. Onida, L. Reining, and A. Rubio, RMP **74** (2002).

# Bound excitons

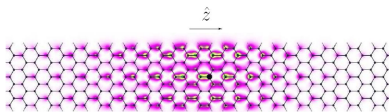
## Solid argon



F. Sottile, M. Marsili, V. Olevano, and L. Reining, PRB **76** (2007).

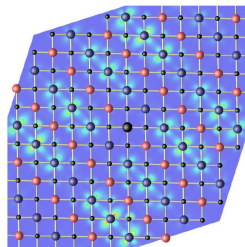
# Exciton analysis

Exciton amplitude: 
$$\Psi_\lambda(\mathbf{r}_h, \mathbf{r}_e) = \sum_{vc} A_\lambda^{(vc)} \phi_v^*(\mathbf{r}_h) \phi_c(\mathbf{r}_e)$$



Graphene nanoribbon

D. Prezzi, *et al.*, PRB **77** (2008).

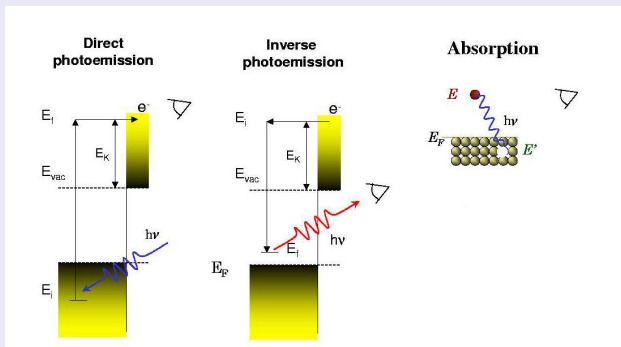


Manganese Oxide

C. Rödl, *et al.*, PRB **77** (2008).

# Electronic Spectroscopy

Here only two categories:



- 1 charged excitations: photoemission and inverse photoemission = **GW**
- 2 neutral excitations: absorption and electron energy loss = **BSE**