

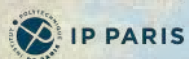
X-ray spectroscopies via Green's functions Theory

the importance of excitonic (many-body) signatures

Francesco Sottile

LSI, Ecole Polytechnique, Palaiseau and ETSF - France

19th November 2024, Campus Luminy, Marseille
Theory around XFELs



X-ray spectroscopies via Green's functions Theory

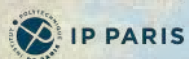
the importance of excitonic (many-body) signatures

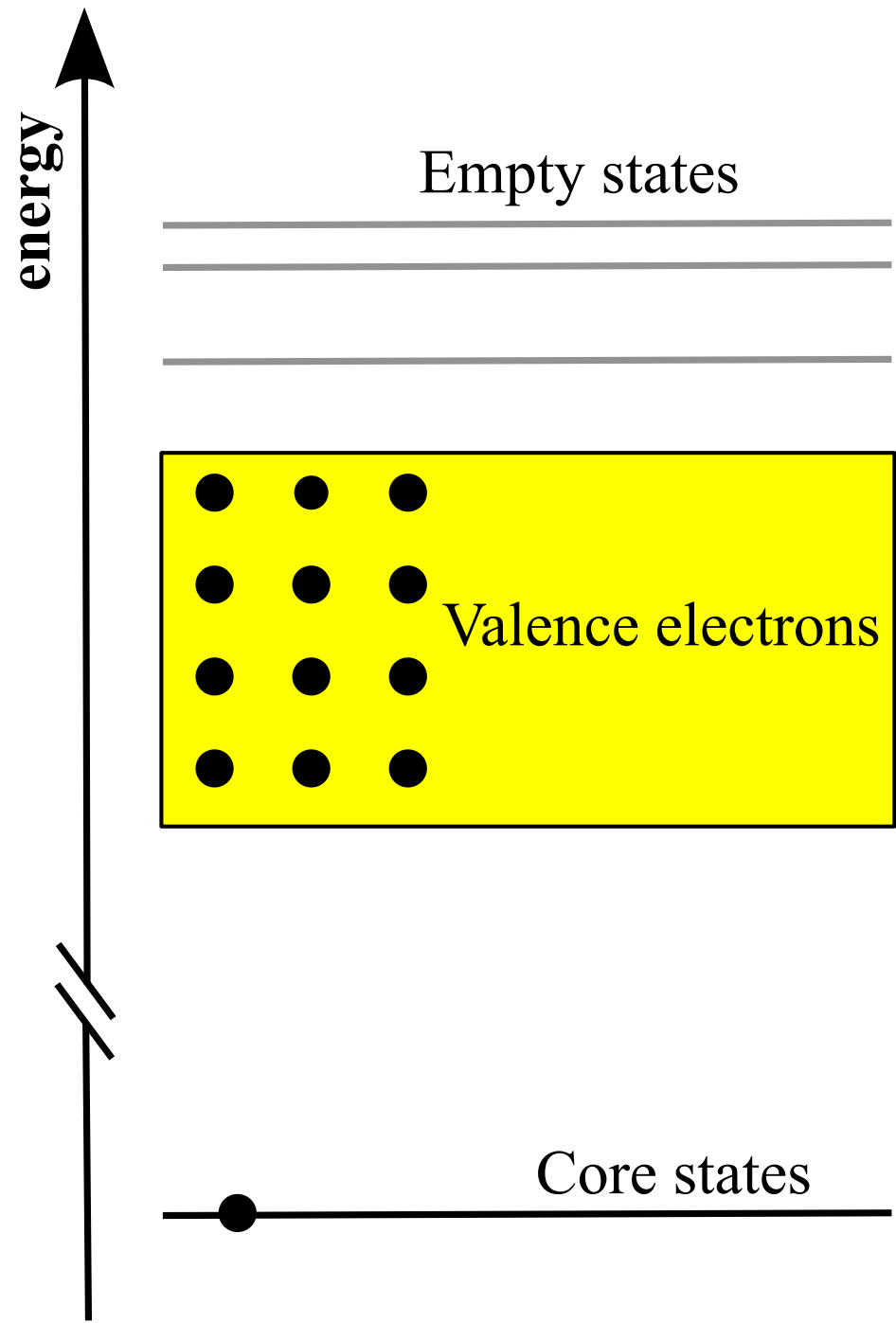
Francesco Sottile

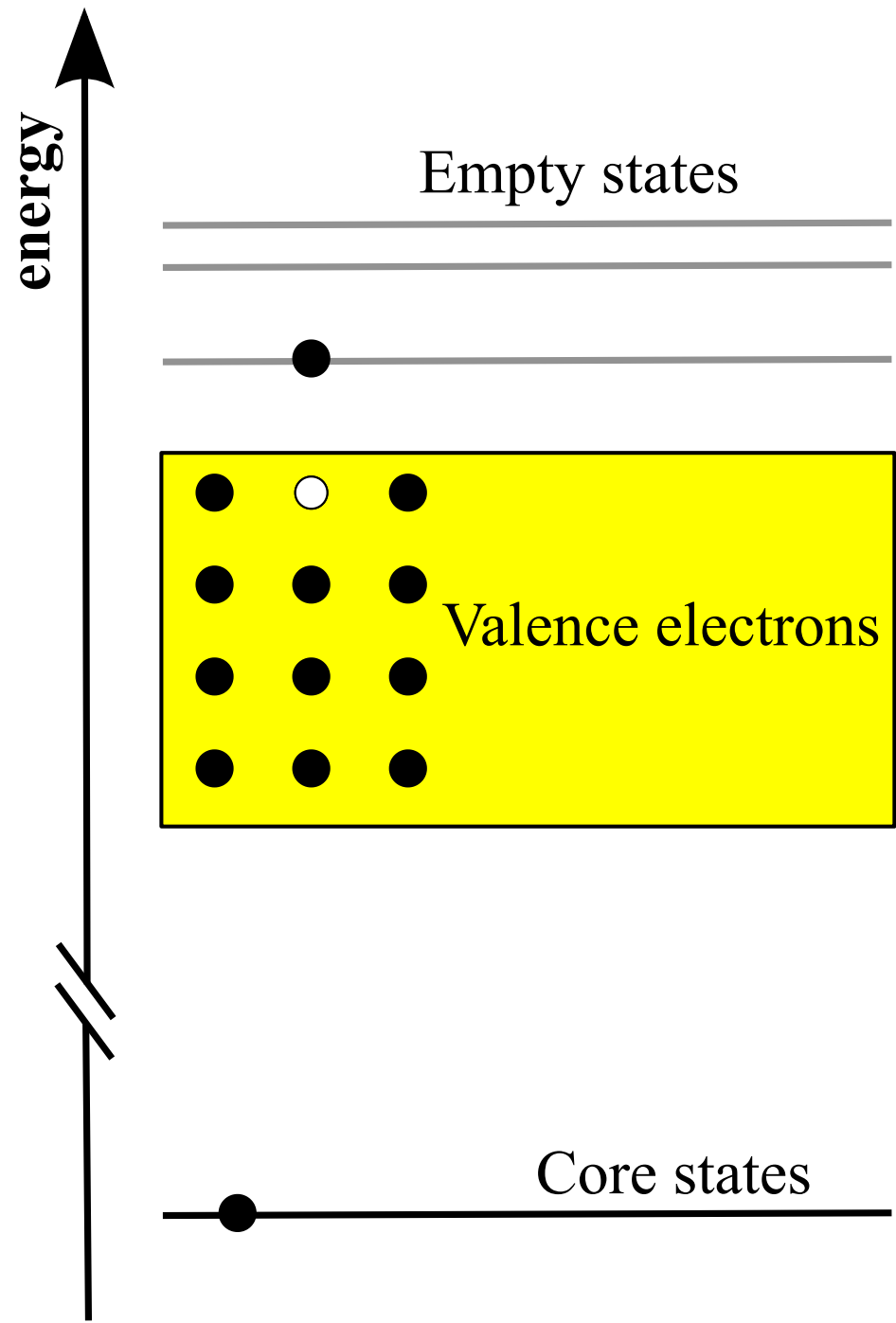
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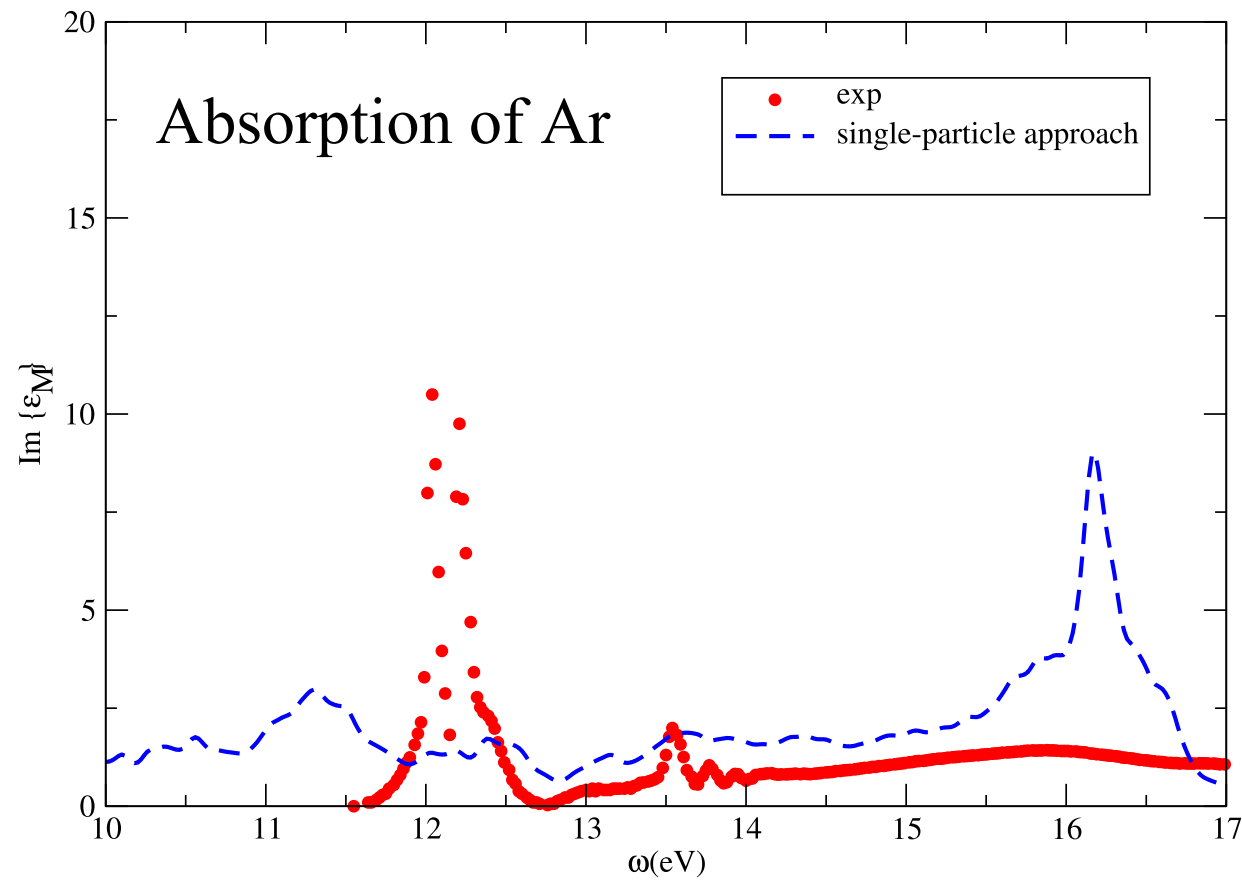
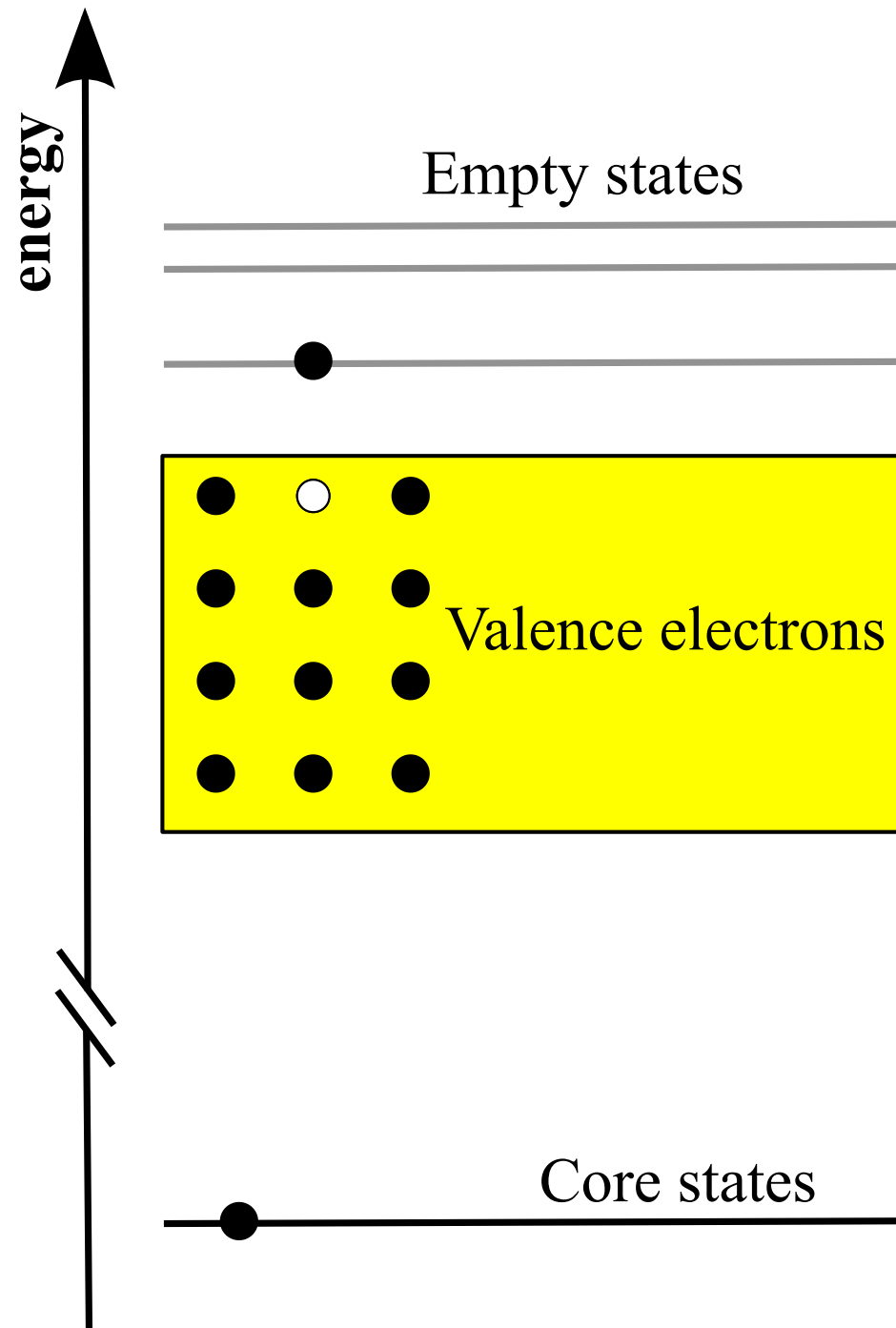
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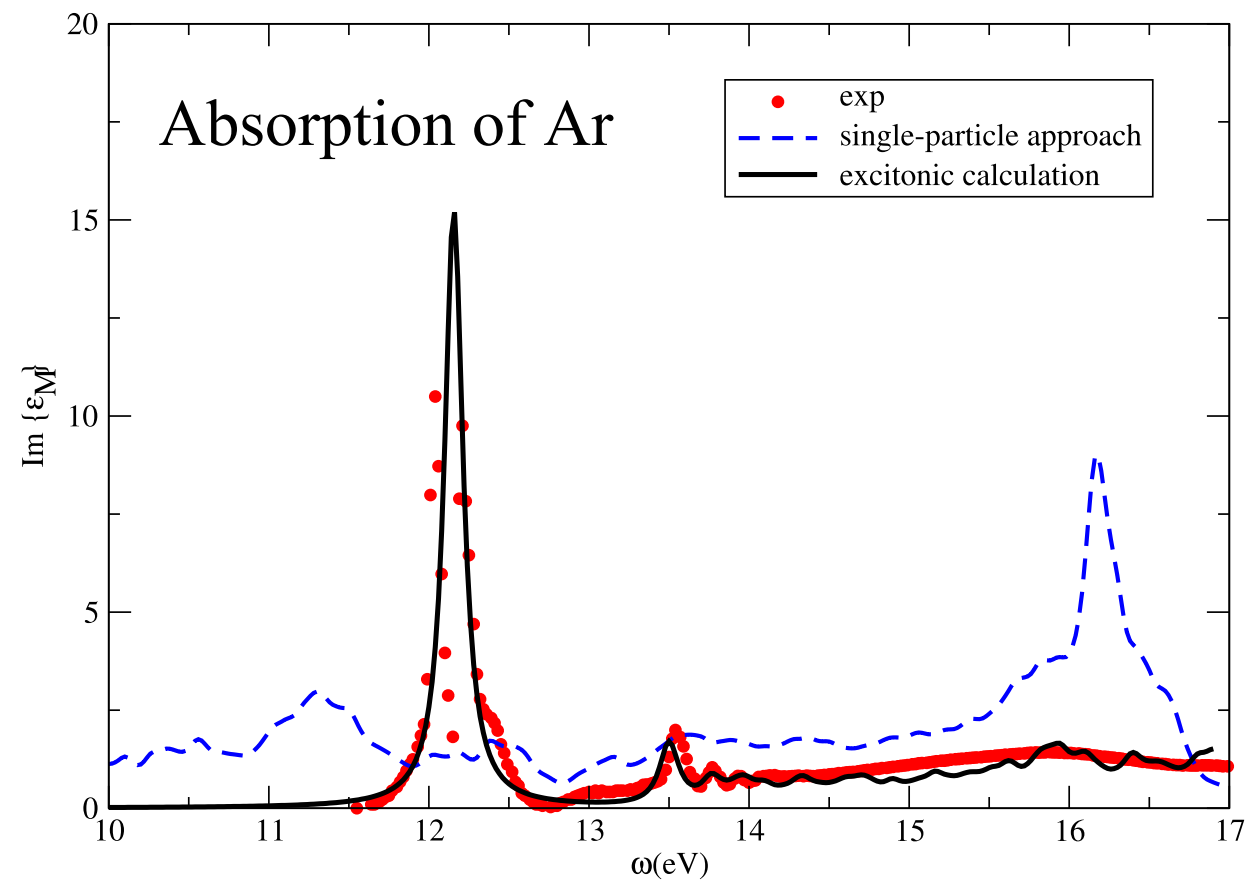
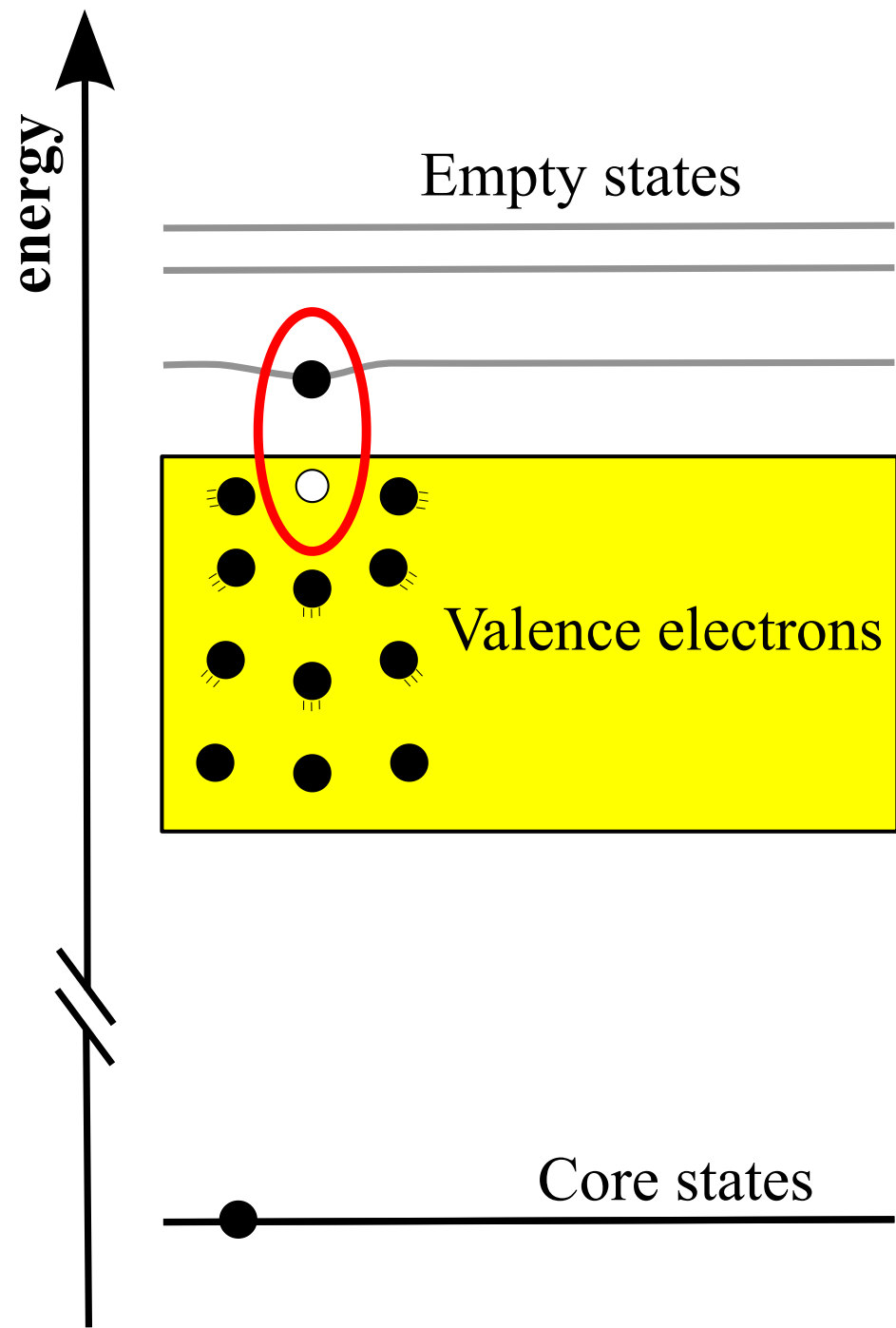
Theory around XFELs












- 
- Excitons via Green's functions many-body theory
Signatures in absorption
 - (N,C,R) Inelastic X-ray scattering within GFs
 - What about XFELs ?

The Green's functions mb formalism

$$G(1, 2)$$

$$1 = (\mathbf{r}_1, t_1, \sigma_1)$$



The Green's functions mb formalism

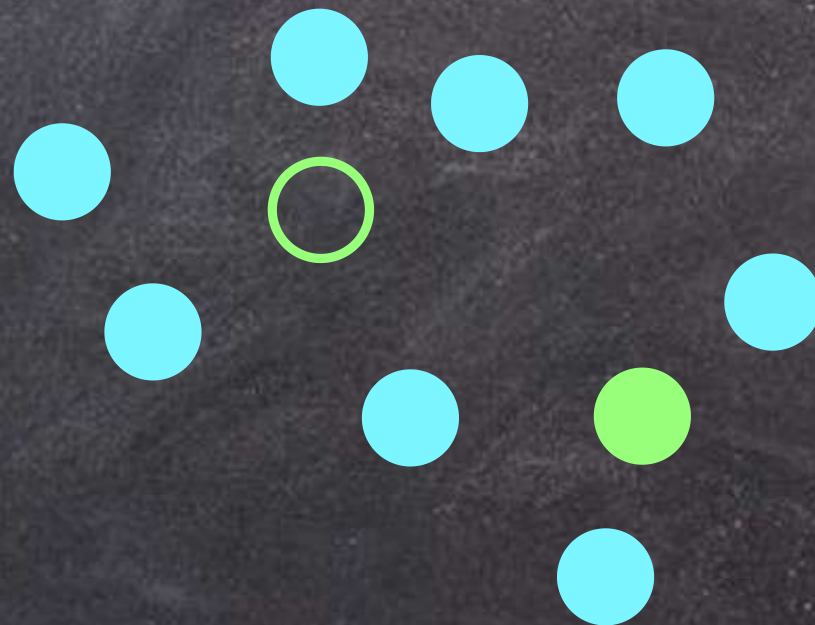
$$G(1, 2)$$



$$2 = (\mathbf{r}_2, t_2, \sigma_2)$$

The Green's functions mb formalism

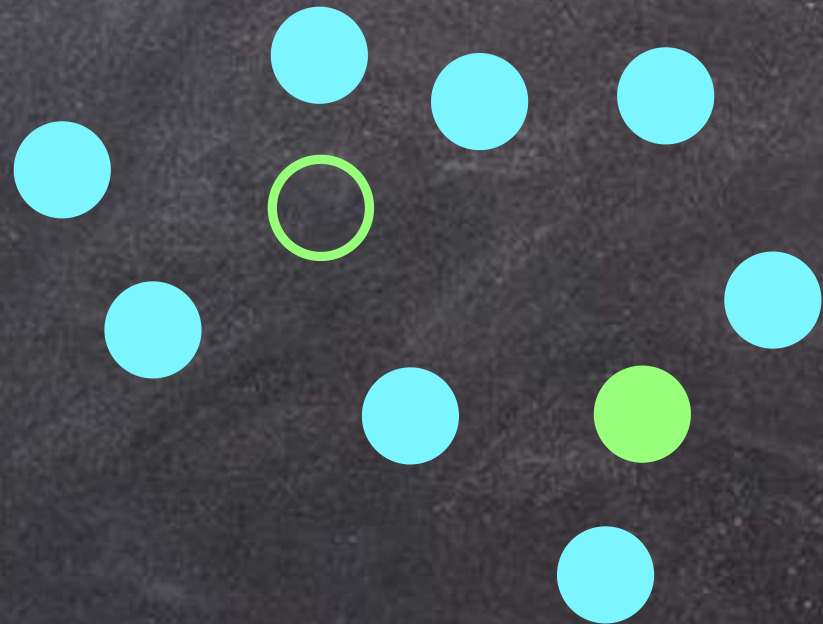
$$G(1, 2)$$



$$G^{(2)}(1, 2, 3, 4)$$

The Green's functions mb formalism

$G(1, 2)$



$$G^{(2)}(1, 2, 3, 4) - G(1, 2)G(3, 4) = -iL(1, 2, 3, 4) = \frac{\delta G(1, 3)}{\delta V_{ext}(2, 4)}$$

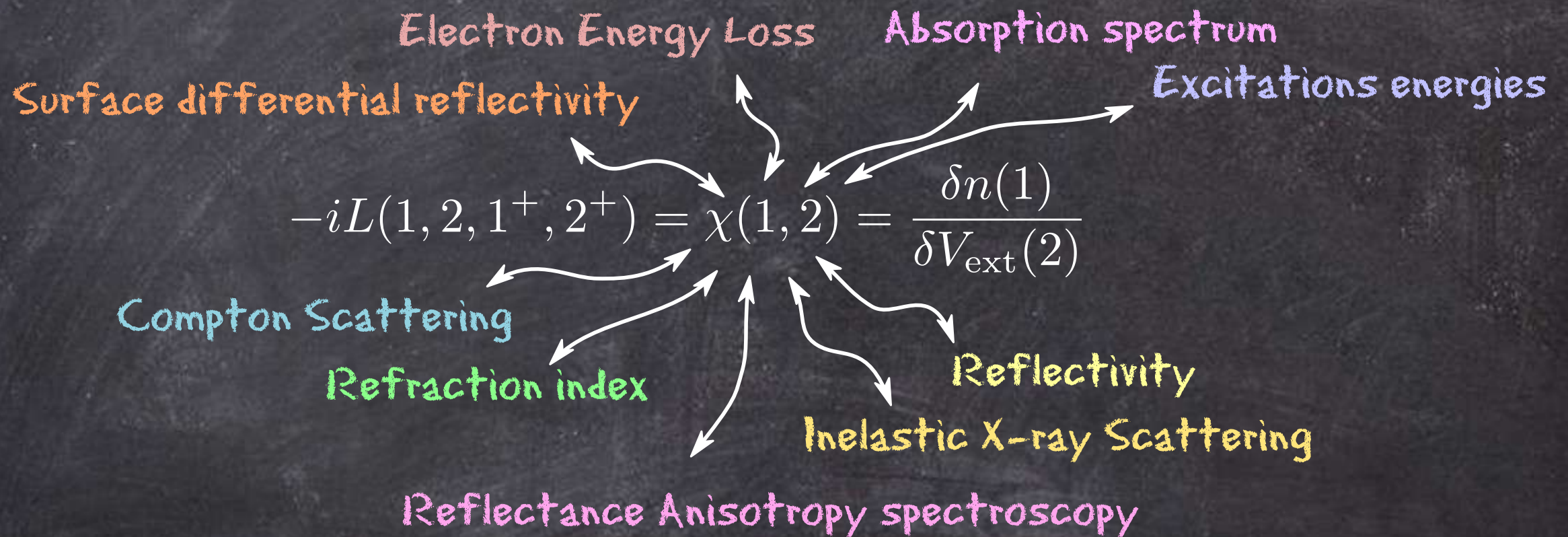
$$-iL(1, 2, 1^+, 2^+) = \chi(1, 2) = \frac{\delta n(1)}{\delta V_{\text{ext}}(2)}$$

Electron Energy Loss

Excitations energies

$$-iL(1, 2, 1^+, 2^+) = \chi(1, 2) = \frac{\delta n(1)}{\delta V_{\text{ext}}(2)}$$

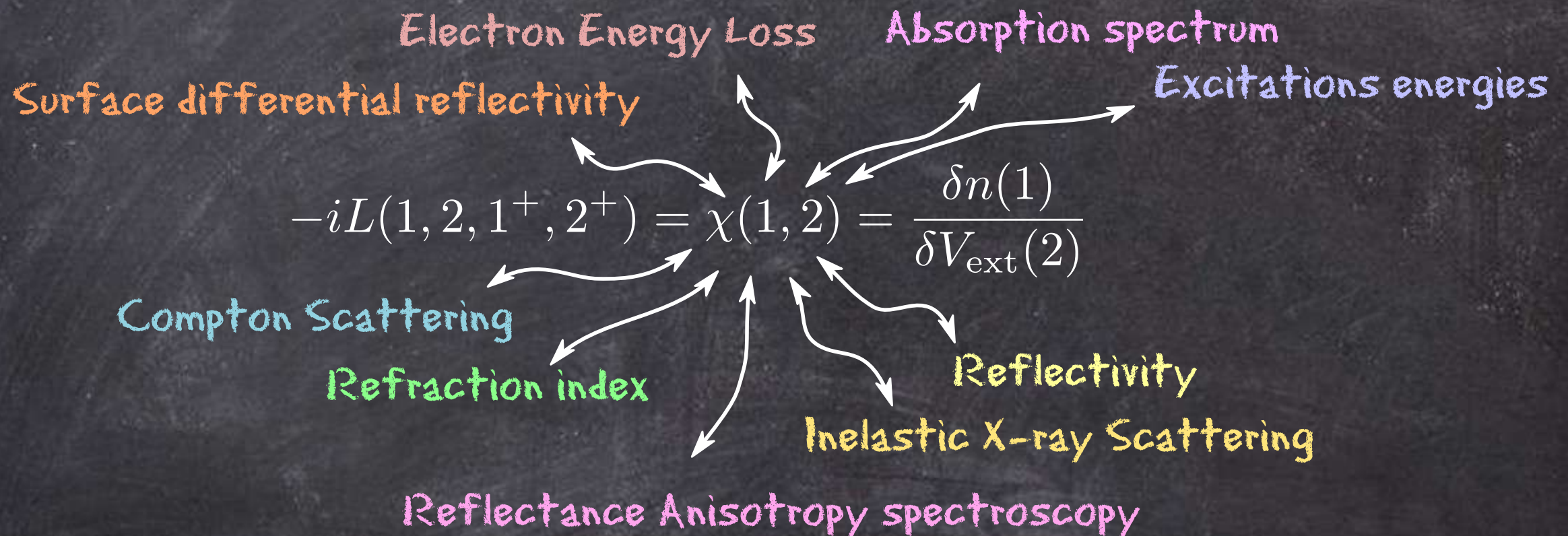
Inelastic X-ray Scattering

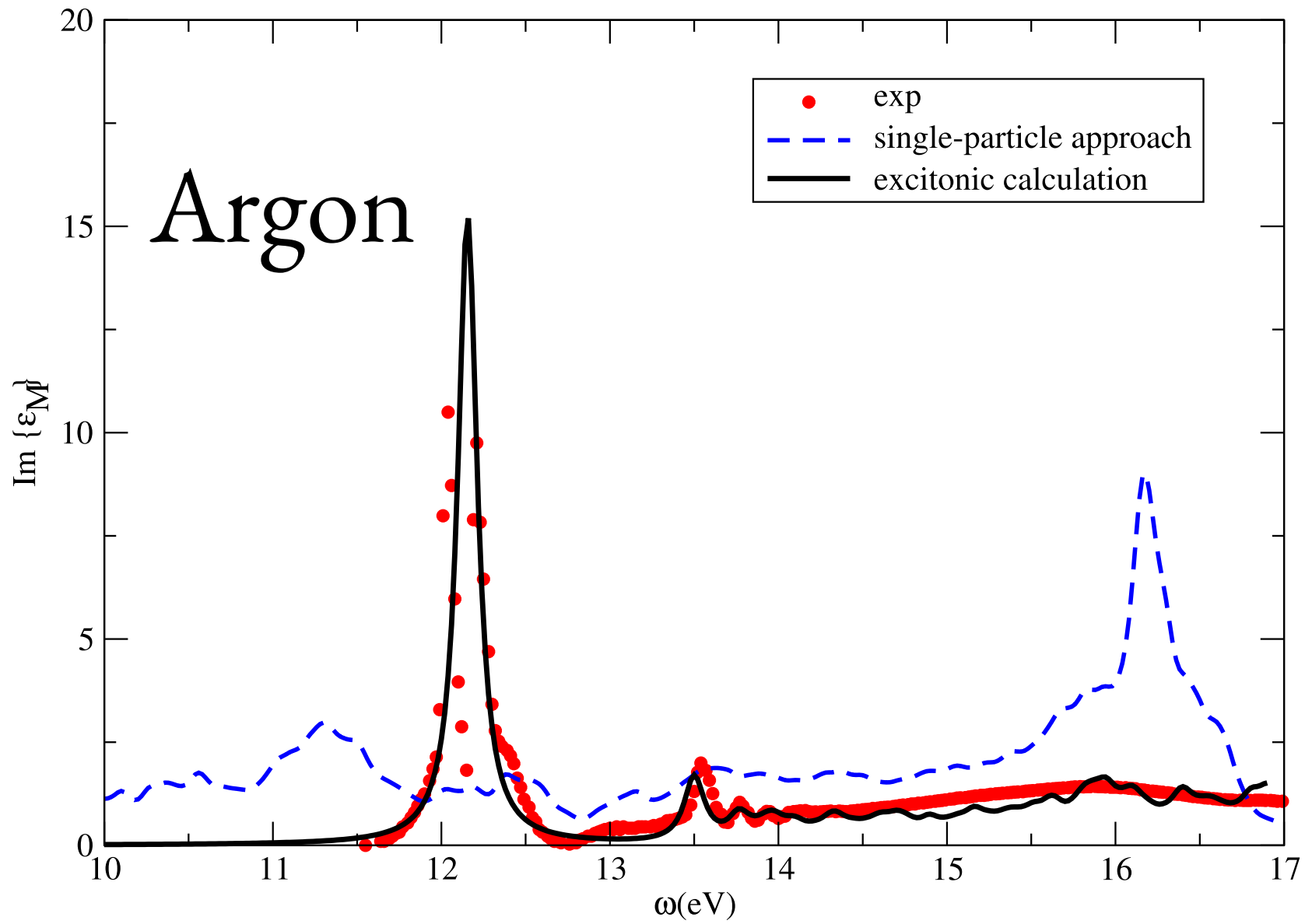


The Bethe-Salpeter Equation

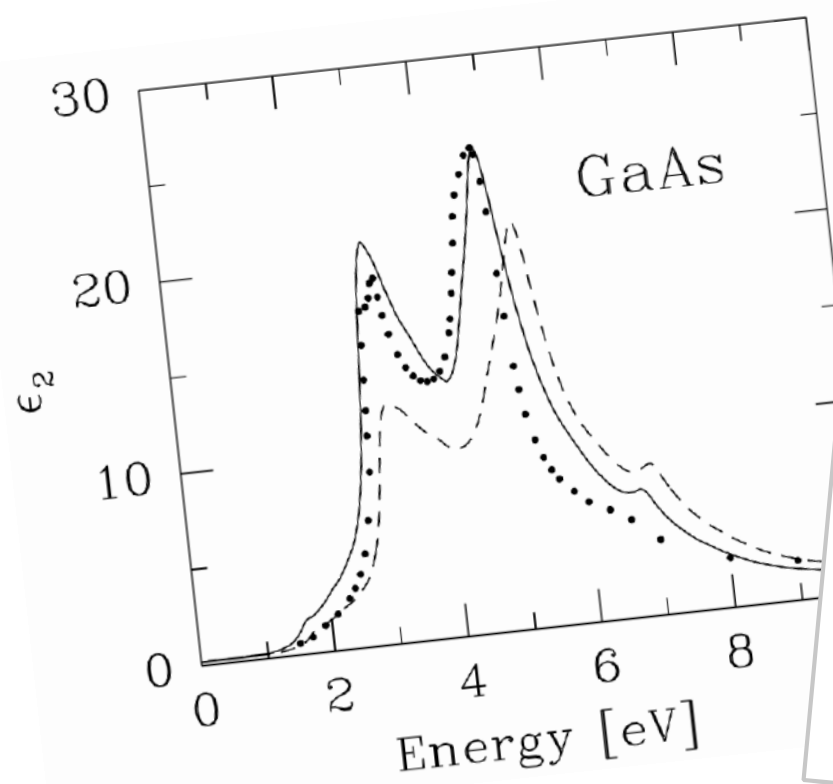
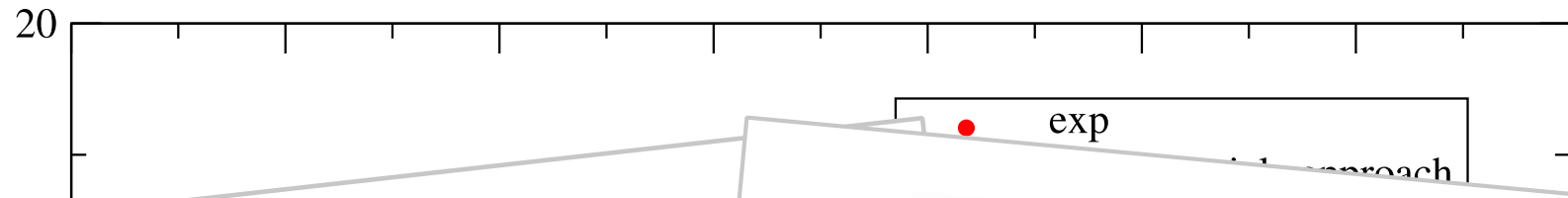
$$L(1, 2, 3, 4) = L^0(1, 2, 3, 4) + L^0(1, 2, 5, 6) \left[v(5, 7) \delta(5, 6) \delta(7, 8) + i \frac{\delta \Sigma(5, 6)}{\delta G(7, 8)} \right] L(7, 8, 3, 4)$$

$$L^0(1, 2, 3, 4) = G(1, 2)G(3, 4)$$

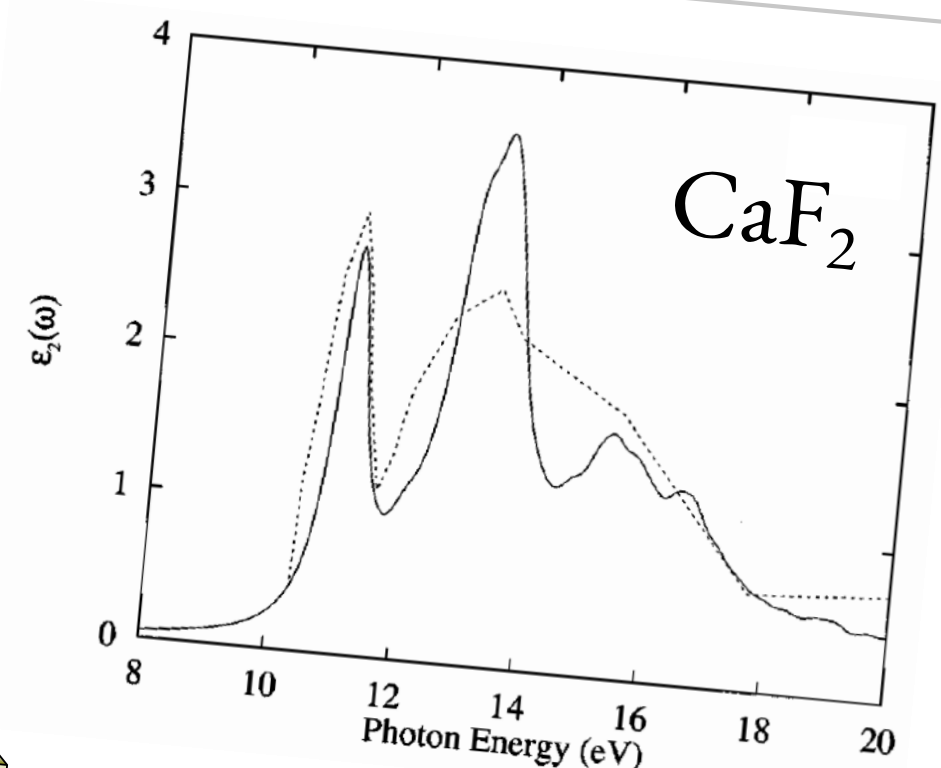





Phys. Rev. B **76** 161103 (2007)

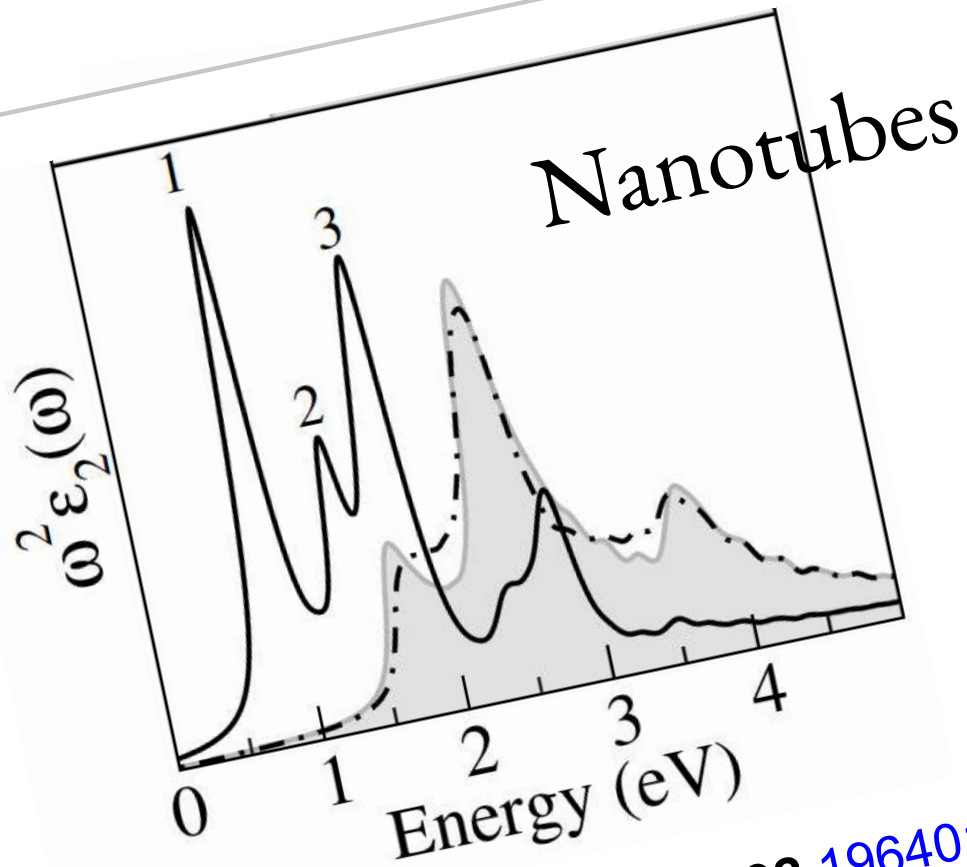
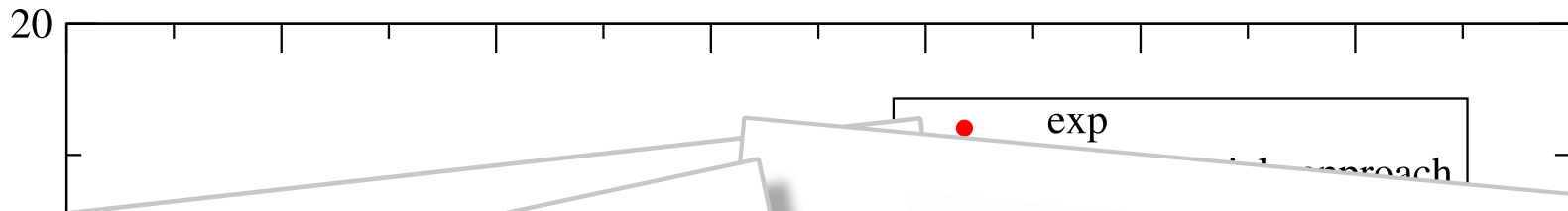


 Rohlfing and Louie Phys. Rev. Lett. **81**, 2312 (1998)



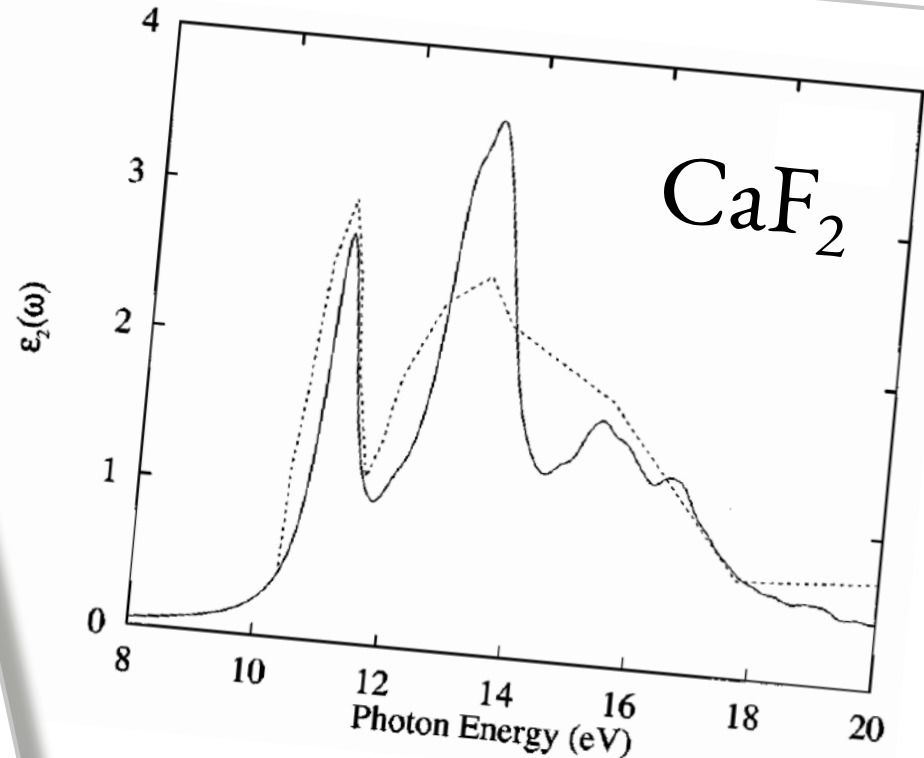
 Benedict and Shirley Phys. Rev. B **59**, 5441 (1999)

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Chang et al., Phys. Rev. Lett. **92** 196401 (2004)

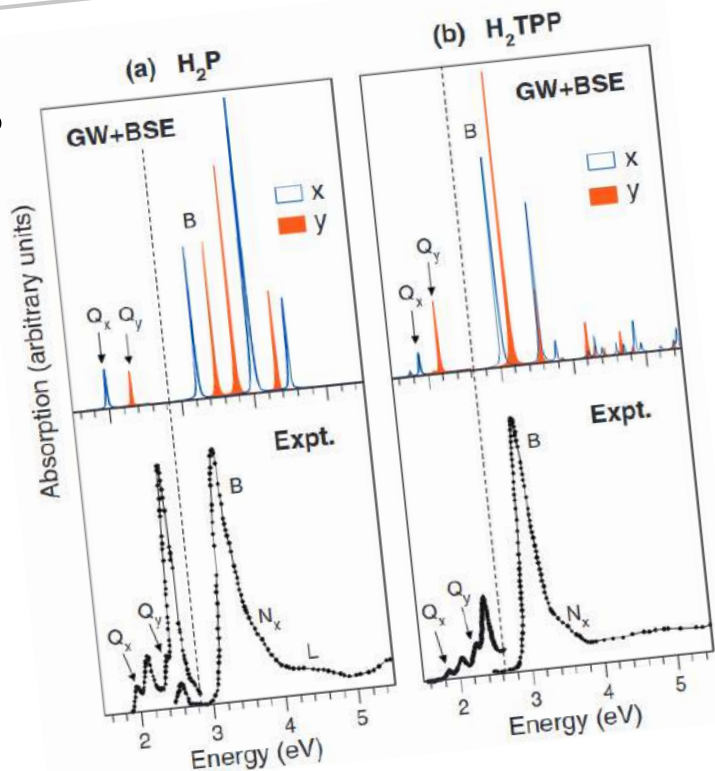
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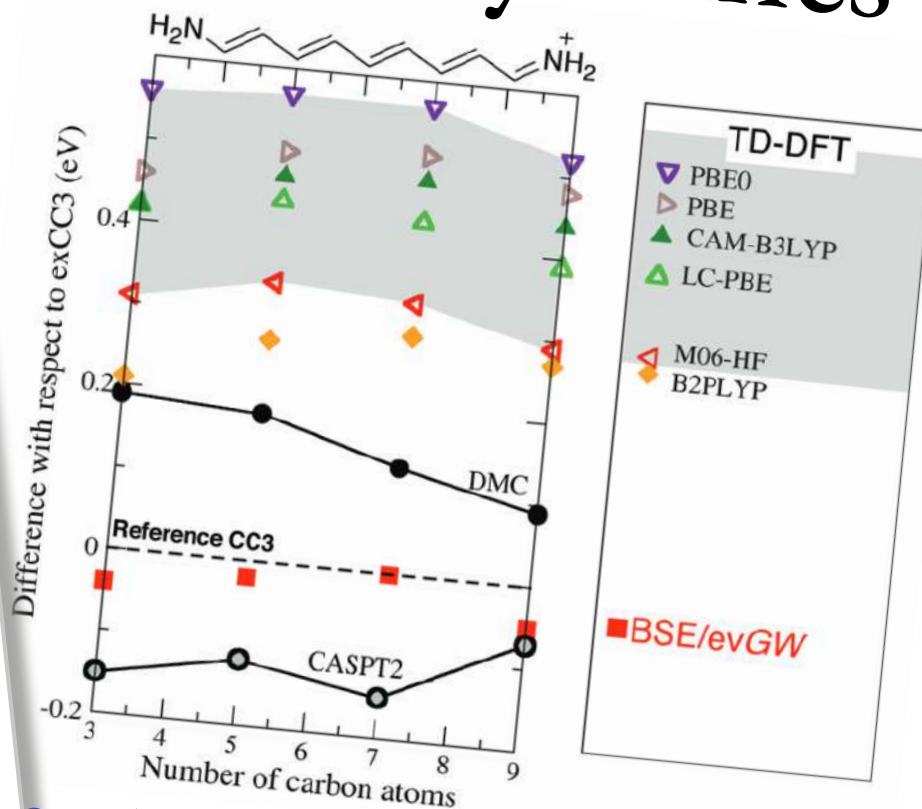
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Porphyrins



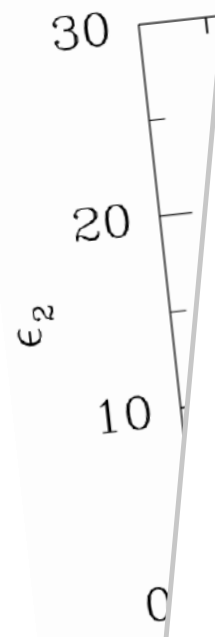
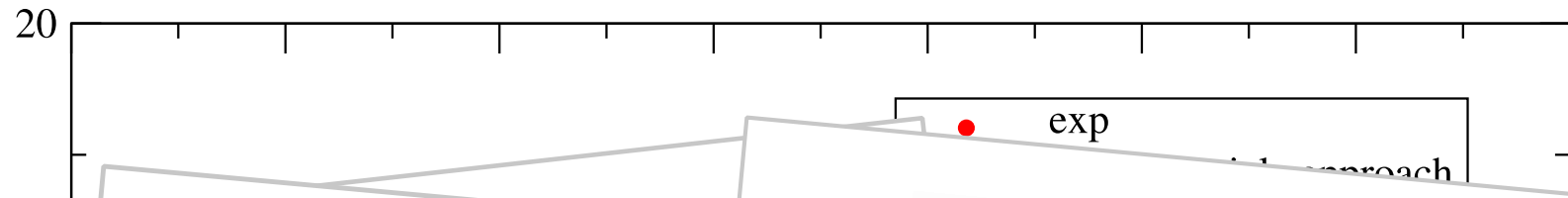
Palumbo *et al.*, *J. Chem. Phys.* **131** 084102 (2009)

streptocyanines

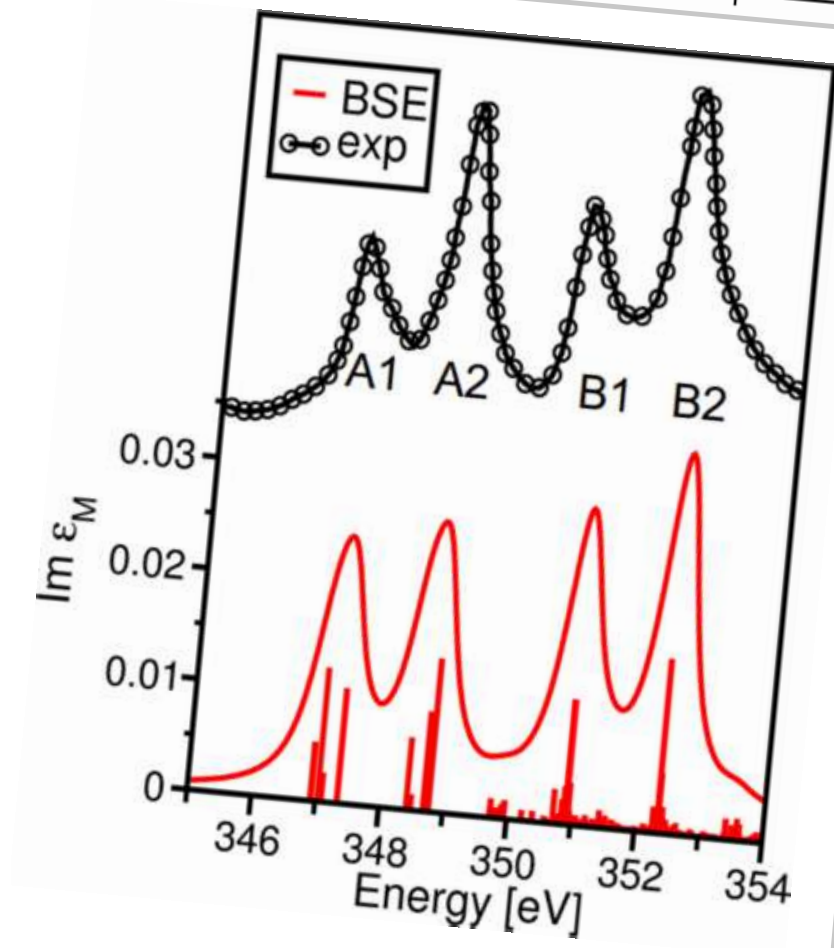


et al. *Chem. Soc. Rev.* **47**, 1022 (2018)

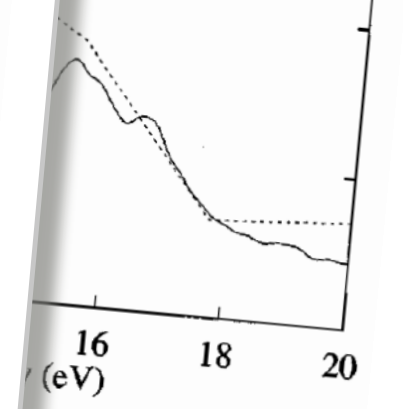
et al. *Phys. Rev. B* **76** 161103 (2007)



CaO
Ca L-edge



CaF₂



B 59, 5441 (1999)

Rohlfing

Vorwerk et al., Phys. Rev. B **95**, 155121 (2017)

Phys. Rev. B

Green's function approach has been (is) successful



Green's function approach has been (is) successful

- it captures the physics of the electron-hole interaction

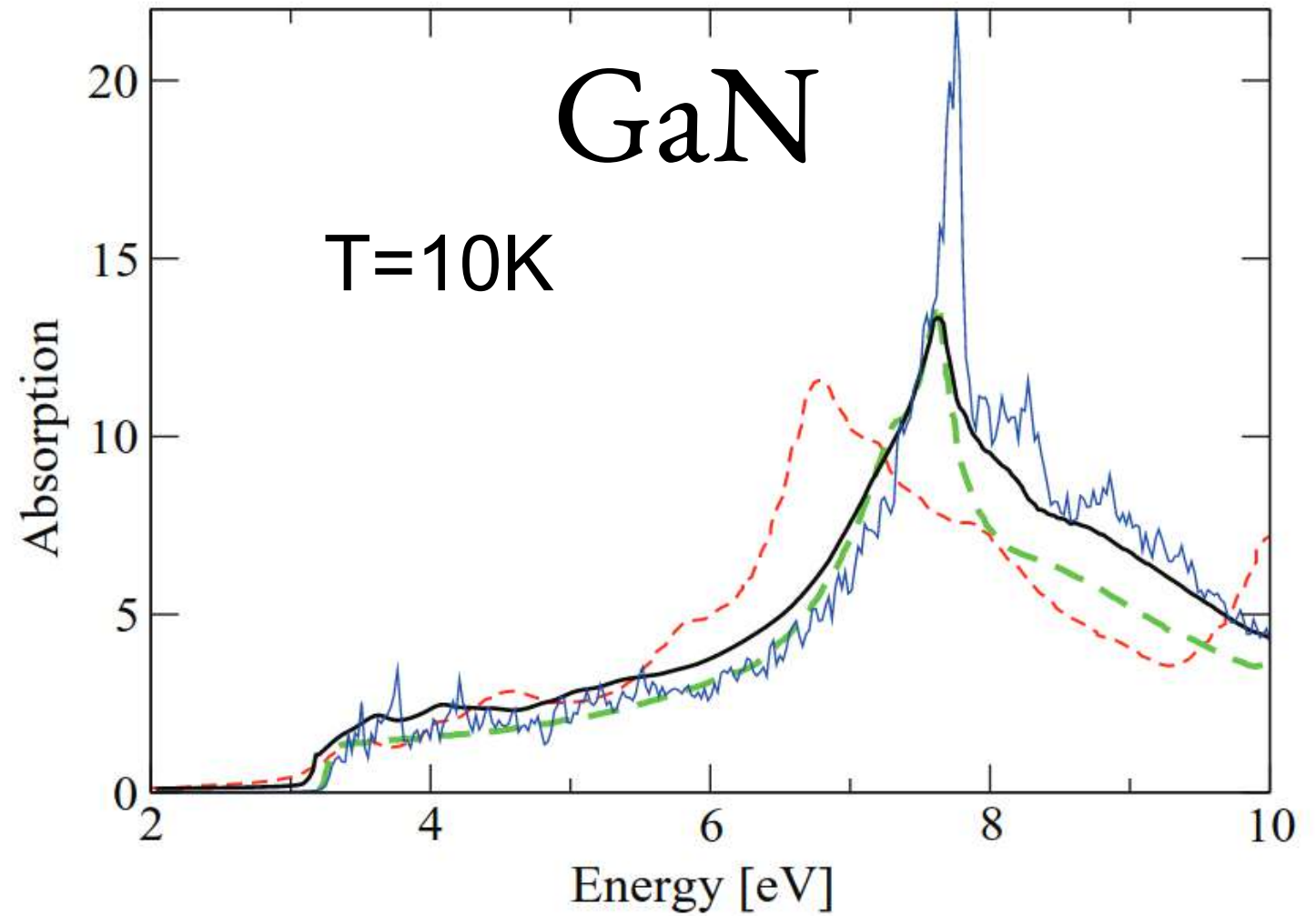


Green's function approach has been (is) successful

- it captures the physics of the electron-hole interaction
- it can (automatically) profit from extensions



Temperature electron-phonon



[Kawai *et al.* Phys. Rev. B **89**, 085202 \(2014\)](#)



Evidence of ideal excitonic insulator in bulk MoS₂ under pressure

S. Samaneh Ataei^{a,1} , Daniele Varsano^{a,1} , Elisa Molinari^{a,b}, and Massimo Rontani^{a,2} 

PNAS 2021, Vol. 118, No. 13, e2010110118

<https://doi.org/10.1073/pnas.2010110118>

PHYSICAL REVIEW B, VOLUME 65, 155332

Bethe-Salpeter equation for magnetoexcitons in quantum wells

Z. G. Koinov*

Department of Physics & Astronomy, University of Texas at San Antonio, San Antonio, Texas 78249

(Received 10 December 2001; published 11 April 2002)

PRL 116, 196804 (2016)

PHYSICAL REVIEW LETTERS

week ending
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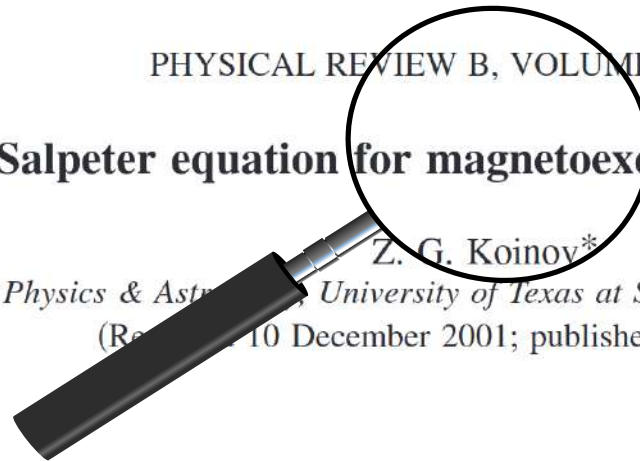
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Green's function approach has been (is) successful

- it captures the physics of the electron-hole interaction
- it can (automatically) profit from extensions
- *ab initio* \longrightarrow predictions

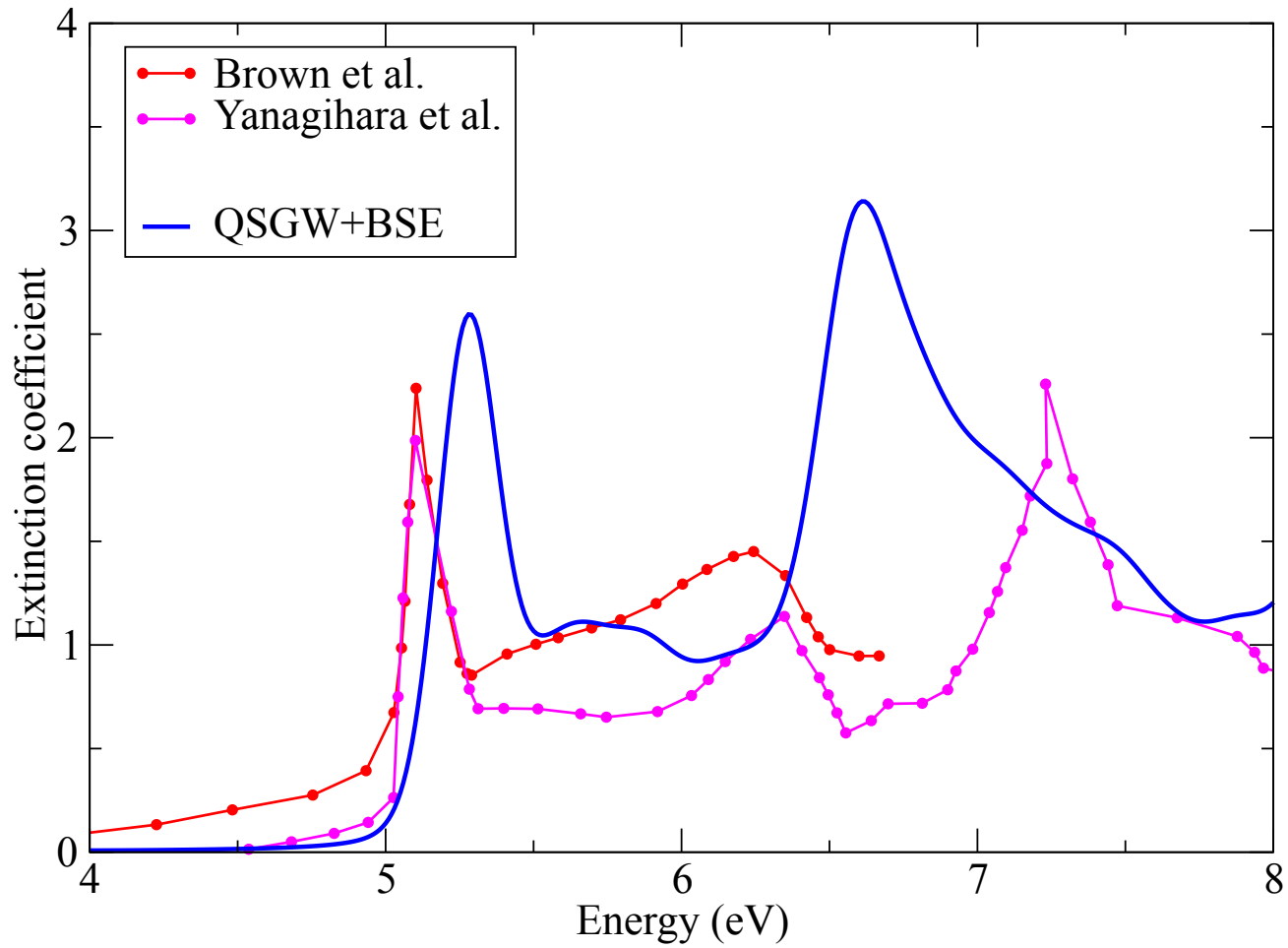


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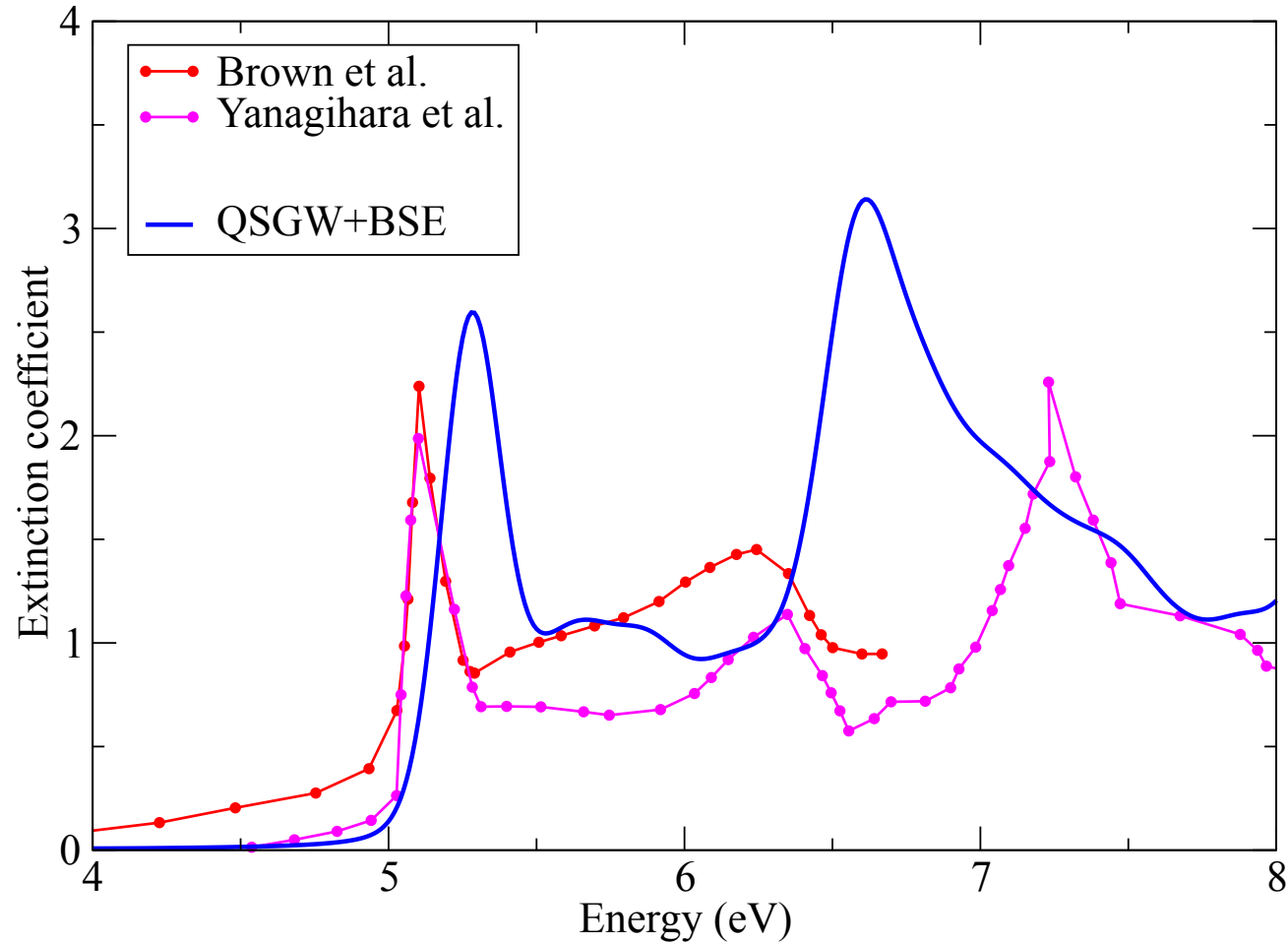
- it captures the physics of the electron-hole interaction
- it can (automatically) profit from extensions
- *ab initio* → predictions
- analysis tools (why? how? who is responsible?)



AgCl extinction coefficient



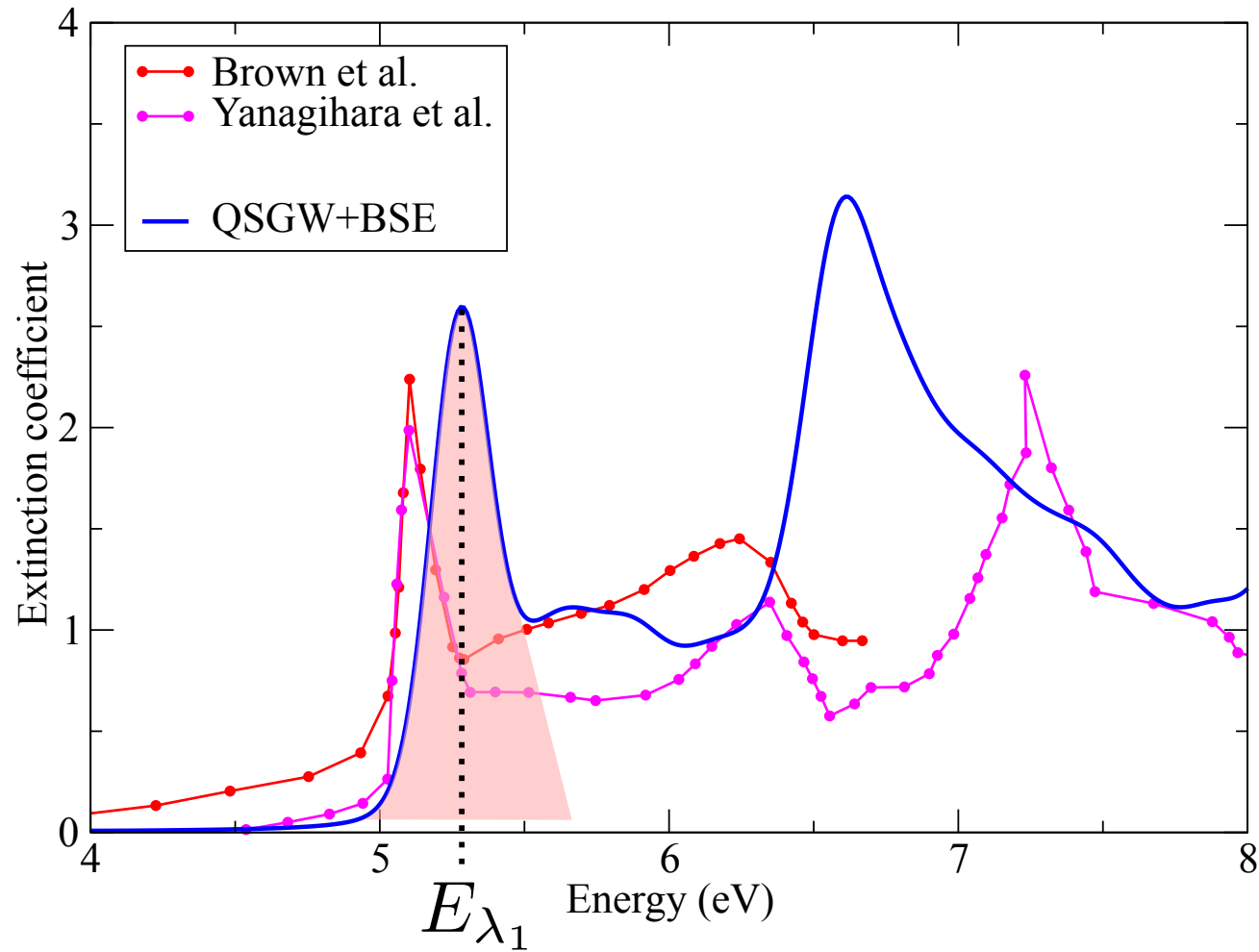
AgCl extinction coefficient



$$\chi_M = \sum_{\lambda} \frac{\left| \sum_{v\mathbf{c}\mathbf{k}} A_{\lambda}^{v\mathbf{c}\mathbf{k}} \langle \mathbf{c}\mathbf{k} | \hat{\mathbf{d}} | v\mathbf{k} \rangle \right|^2}{\omega - E_{\lambda} + i\eta}$$

$$\kappa = \text{Im} \sqrt{\frac{1}{1 + v_0 \chi_M}}$$

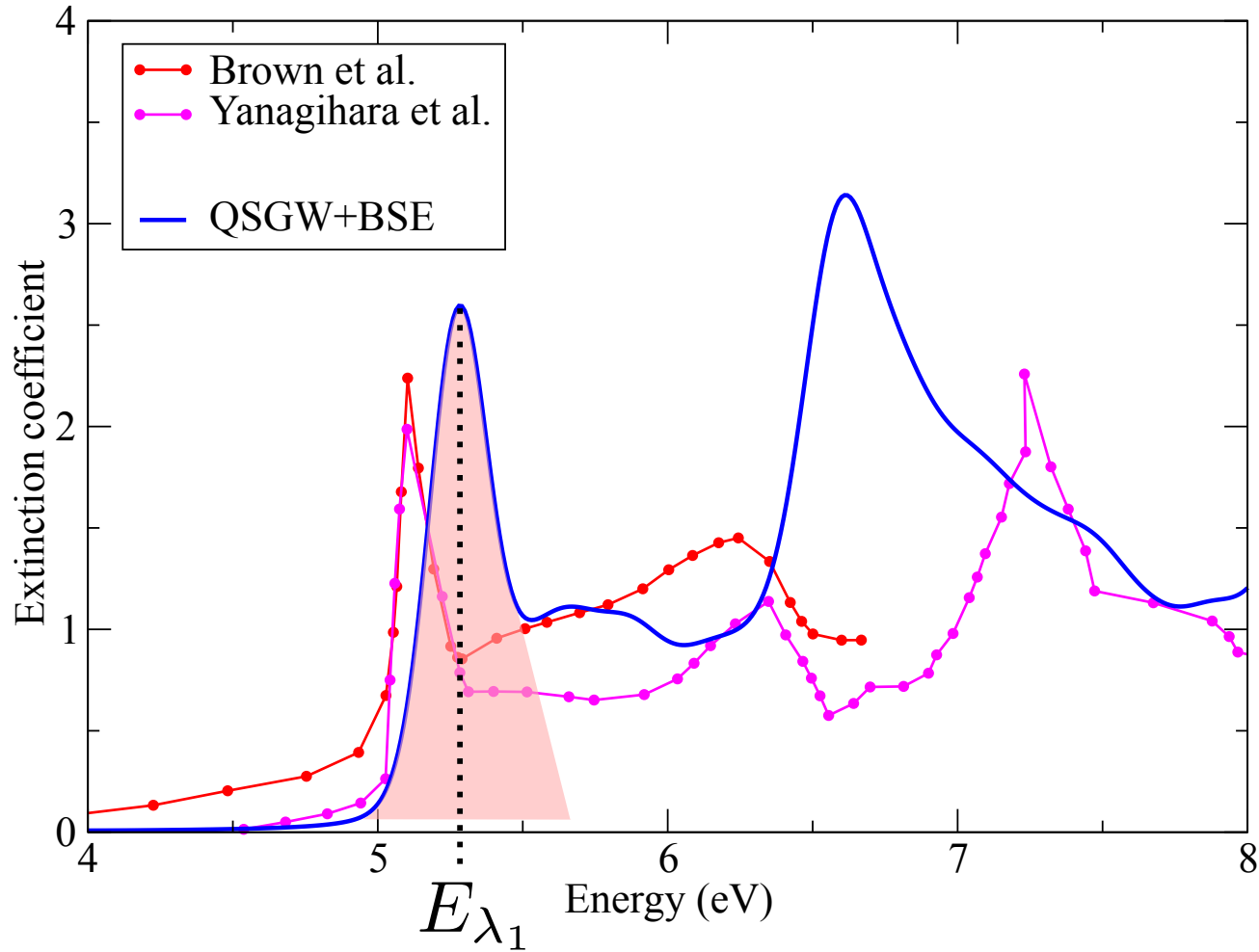
AgCl extinction coefficient



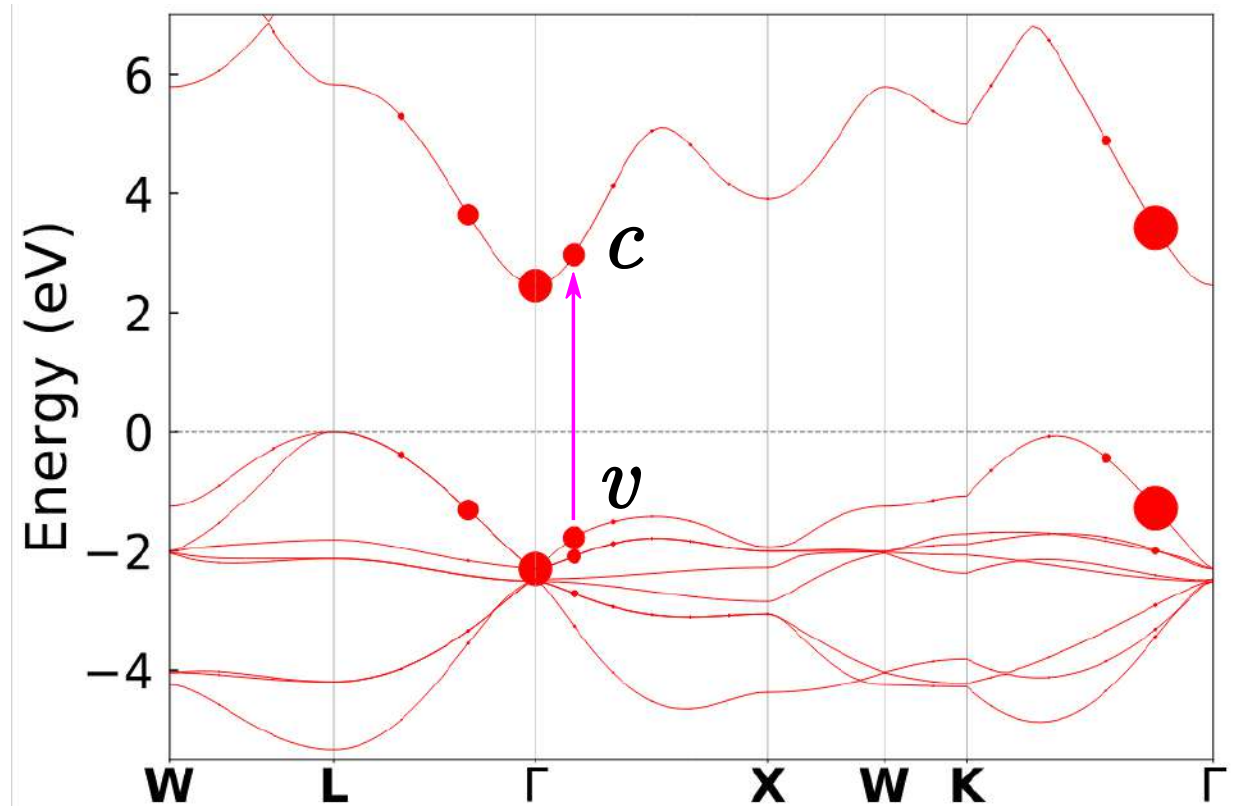
$$\chi_M = \sum_{\lambda} \frac{\left| \sum_{v\mathbf{ck}} A_{\lambda}^{v\mathbf{ck}} \langle \mathbf{ck} | \hat{\mathbf{d}} | v\mathbf{k} \rangle \right|^2}{\omega - E_{\lambda} + i\eta}$$

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AgCl extinction coefficient



$$\chi_M = \sum_{\lambda} \frac{\left| \sum_{vck} A_{\lambda_1}^{vck} \langle ck | \hat{d} | vk \rangle \right|^2}{\omega - E_{\lambda} + i\eta}$$



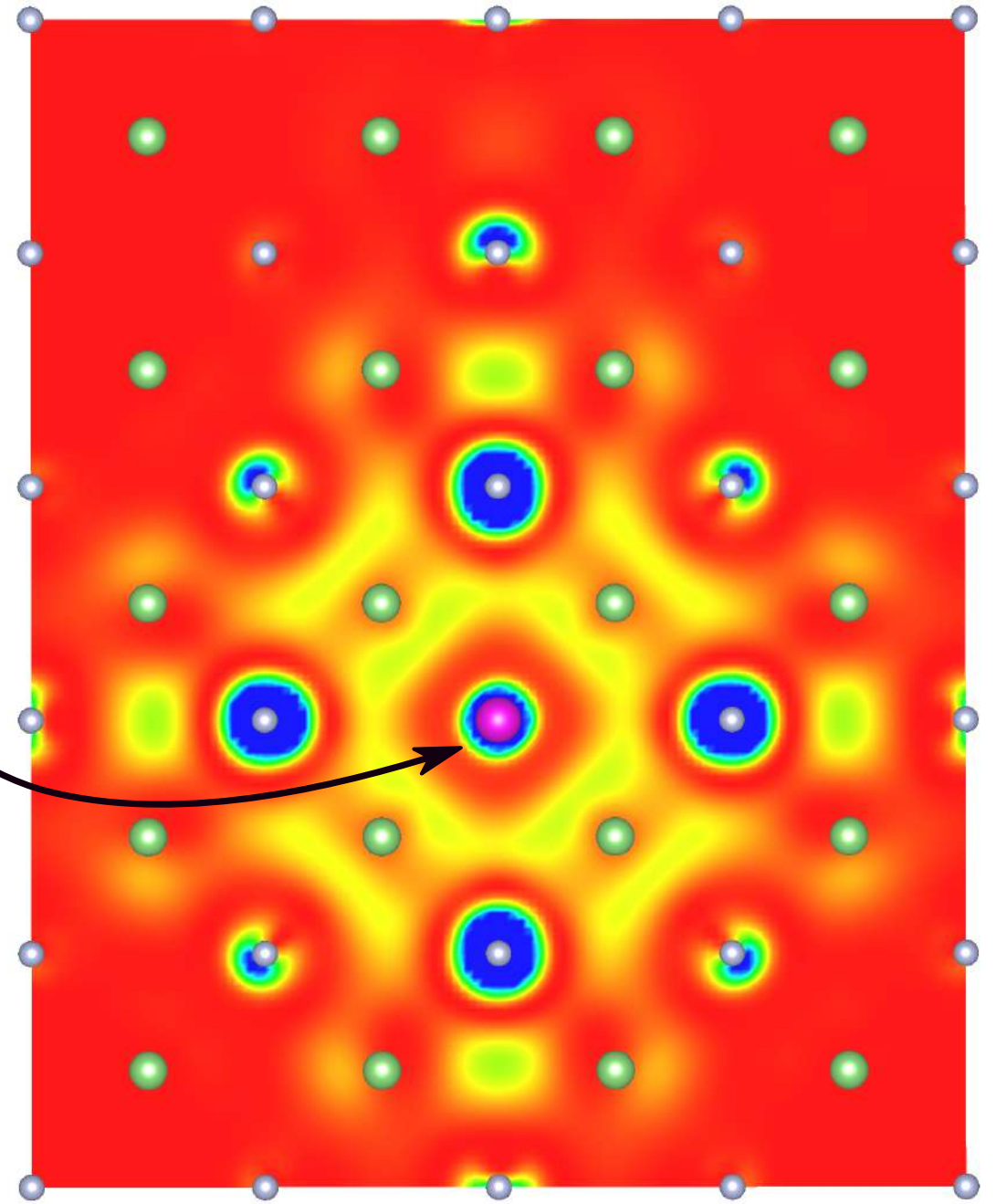
Excitonic wavefunction of LiF

$$\Psi_{\lambda}(\mathbf{r}_e, \mathbf{r}_h) = \sum_{v\mathbf{k}} A_{\lambda}^{v\mathbf{k}} \psi_{c\mathbf{k}}^*(\mathbf{r}_e) \psi_{v\mathbf{k}}(\mathbf{r}_h)$$

Excitonic wavefunction of LiF

$$|\Psi_\lambda(\mathbf{r}_e, \mathbf{r}_h)|^2 = \left| \sum_{v\mathbf{k}} A_\lambda^{v\mathbf{k}} \psi_{c\mathbf{k}}^*(\mathbf{r}_e) \psi_{v\mathbf{k}}(\mathbf{r}_h) \right|^2$$

- where is the exciton localised ?
- how much ?



- 
- Excitons via Green's functions many-body theory
Signatures in absorption
 - (N,C,R) Inelastic X-ray scattering within GFs
 - What about XFELs ?



Christian Vorwerk (thesis with C. Draxl)



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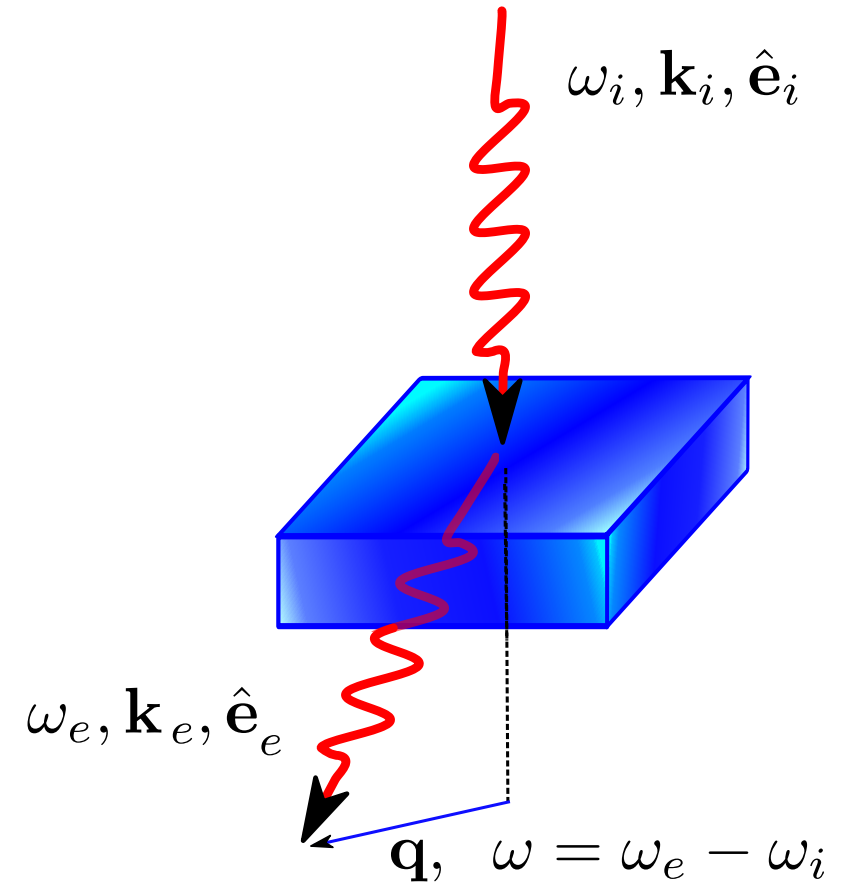


Matteo Gatti



Laura Urquiza

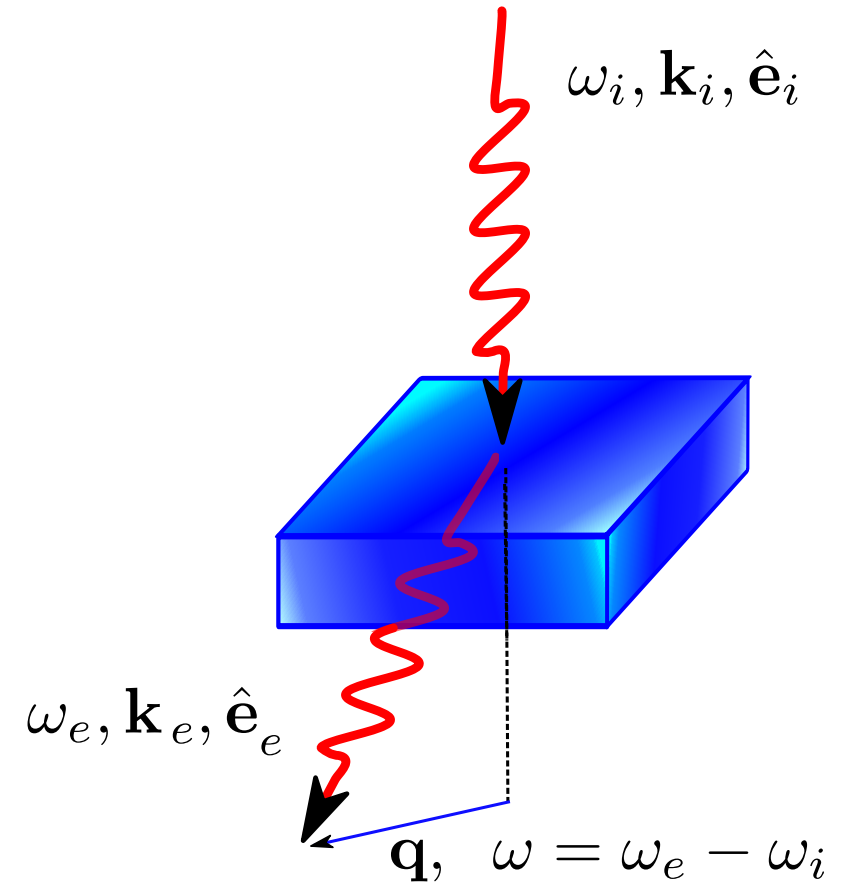
X-ray scattering



$$\frac{d^2\sigma}{d\Omega_2 d\omega_e} \propto \sum_f \left| \langle f | e^{i\mathbf{q}\cdot\mathbf{r}} | 0 \rangle - \frac{i\omega_i/e}{2mc^2} \sum_n \frac{\langle f | e^{-i\mathbf{k}_f\cdot\mathbf{r}} \nabla | n \rangle \langle n | e^{i\mathbf{k}_i\cdot\mathbf{r}} \nabla | 0 \rangle}{\omega_i - (E_n - E_0)} \right|^2 \times \delta(\omega - (E_f - E_0))$$

X-ray scattering

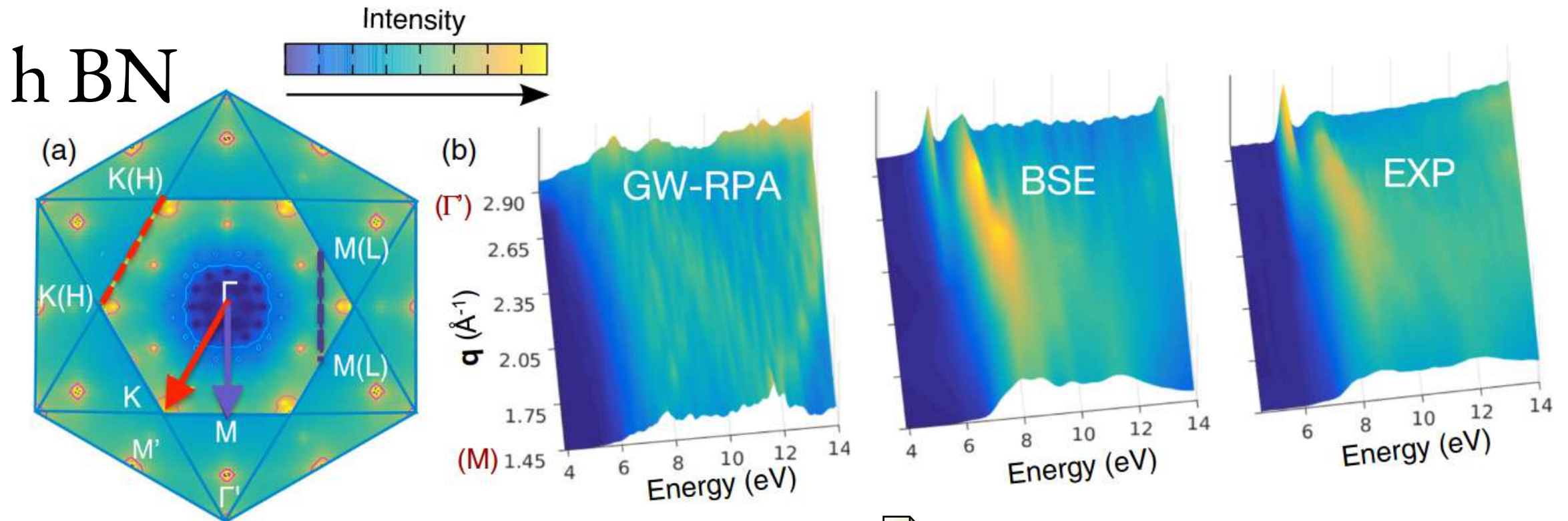
non-Resonant IXS



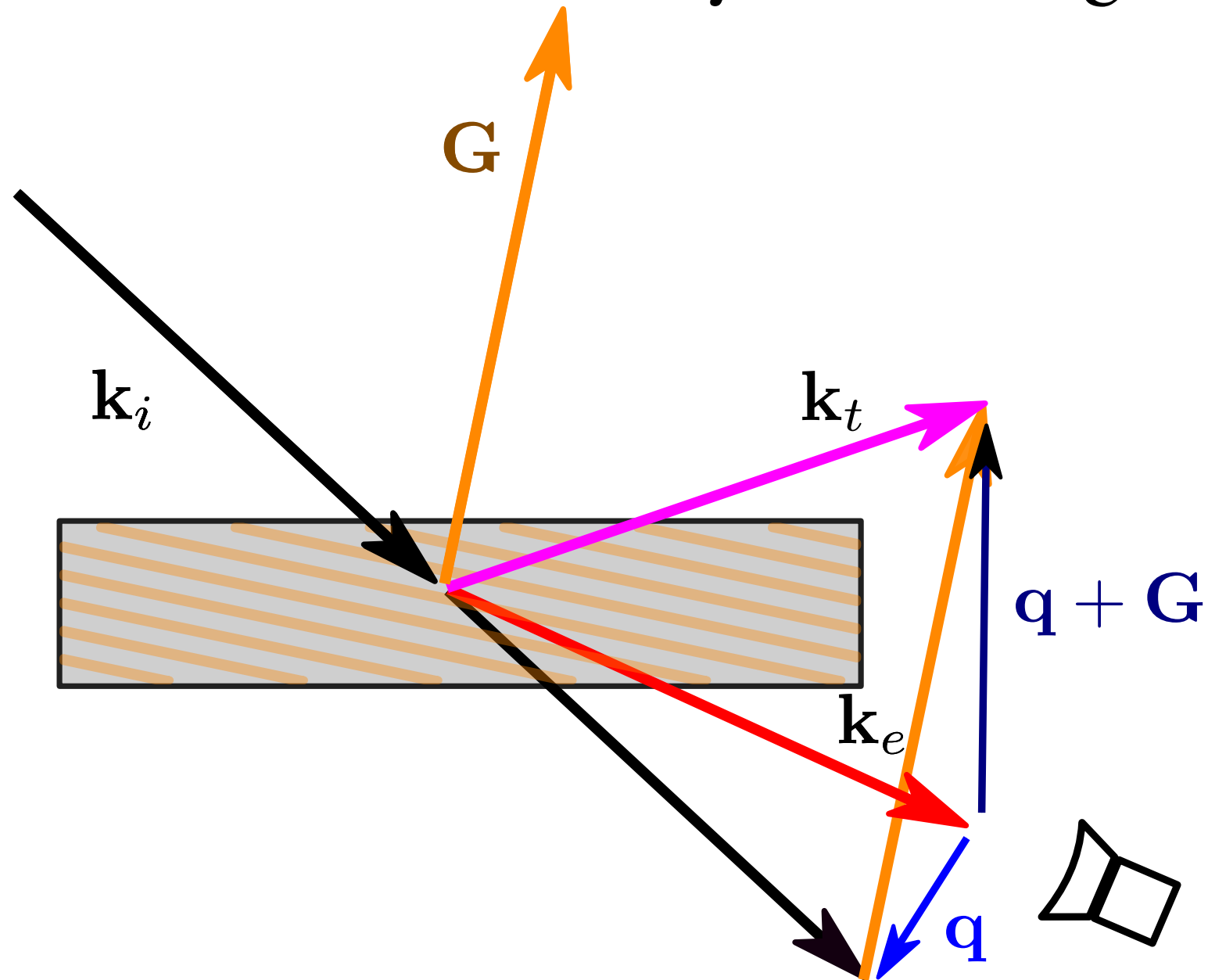
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Bethe-Salpeter Equation - finite momentum transfer

$$S(\mathbf{q}, \omega) \propto \chi_M(\mathbf{q}, \omega) = \sum_{\lambda} \frac{|\sum_{vc} A_{\lambda}^{vc, \mathbf{q}} \langle c | e^{i\mathbf{q} \cdot \mathbf{r}} | v \rangle|^2}{\omega - E_{\lambda}(\mathbf{q}) + i\eta}$$




Coherent Inelastic X-ray scattering

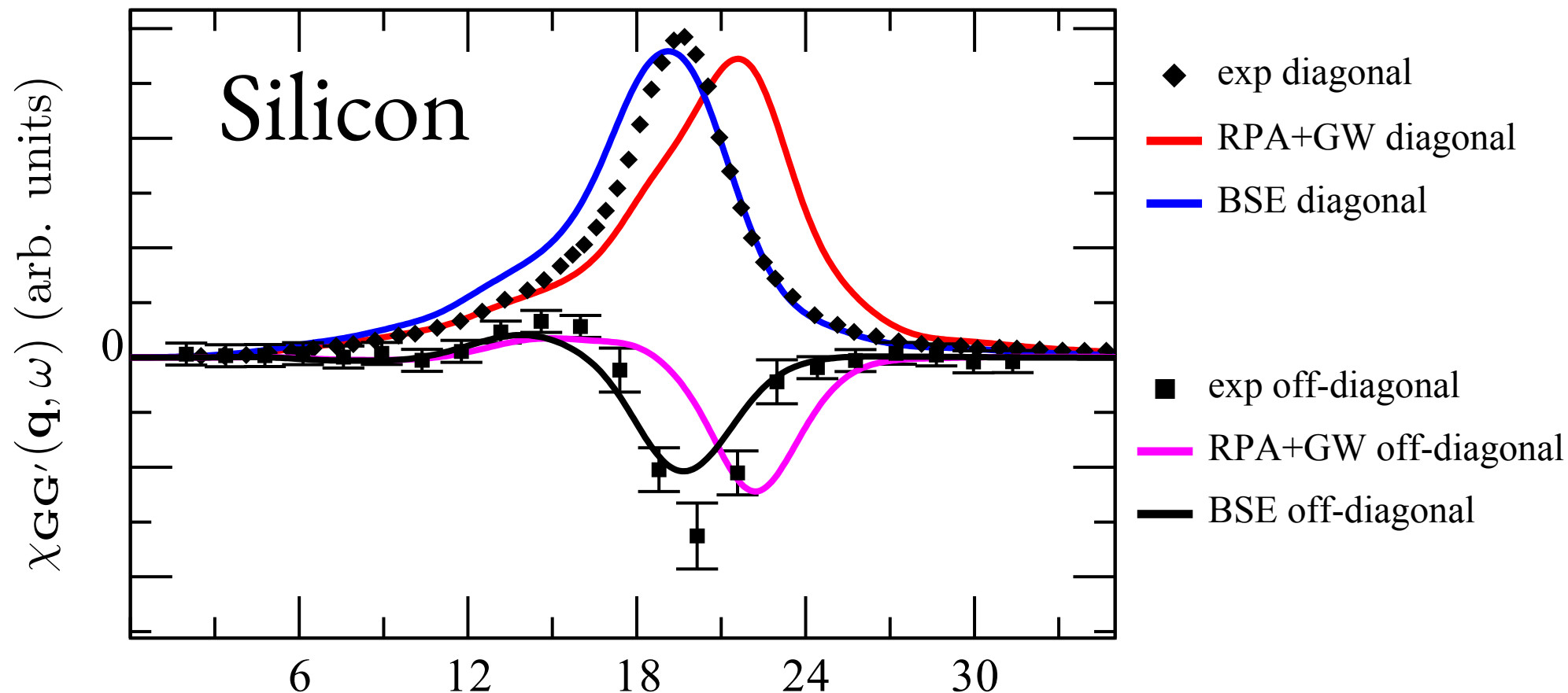


Coherent Inelastic X-ray scattering

$$\chi(\mathbf{q}, \mathbf{q} + \mathbf{G}, \omega) = \sum_{\lambda\lambda'} \frac{\sum_{vc} A_{\lambda}^{vc, \mathbf{q}} \langle c | e^{-i\mathbf{q}\cdot\mathbf{r}} | v \rangle S_{\lambda\lambda'}^{-1} \sum_{v'c'} A_{\lambda}^{*, v'c', \mathbf{q}} \langle v' | e^{i(\mathbf{q}+\mathbf{G})\cdot\mathbf{r}} | c' \rangle}{\omega - E_{\lambda}(\mathbf{q}) + i\eta}$$



 Igor Reshetnyak *et al.*
Phys. Rev. Research **1**,
032010(R) (2019)




 Schülke and Kaprolat, Phys. Rev. Lett. **67**, 879 (1991).

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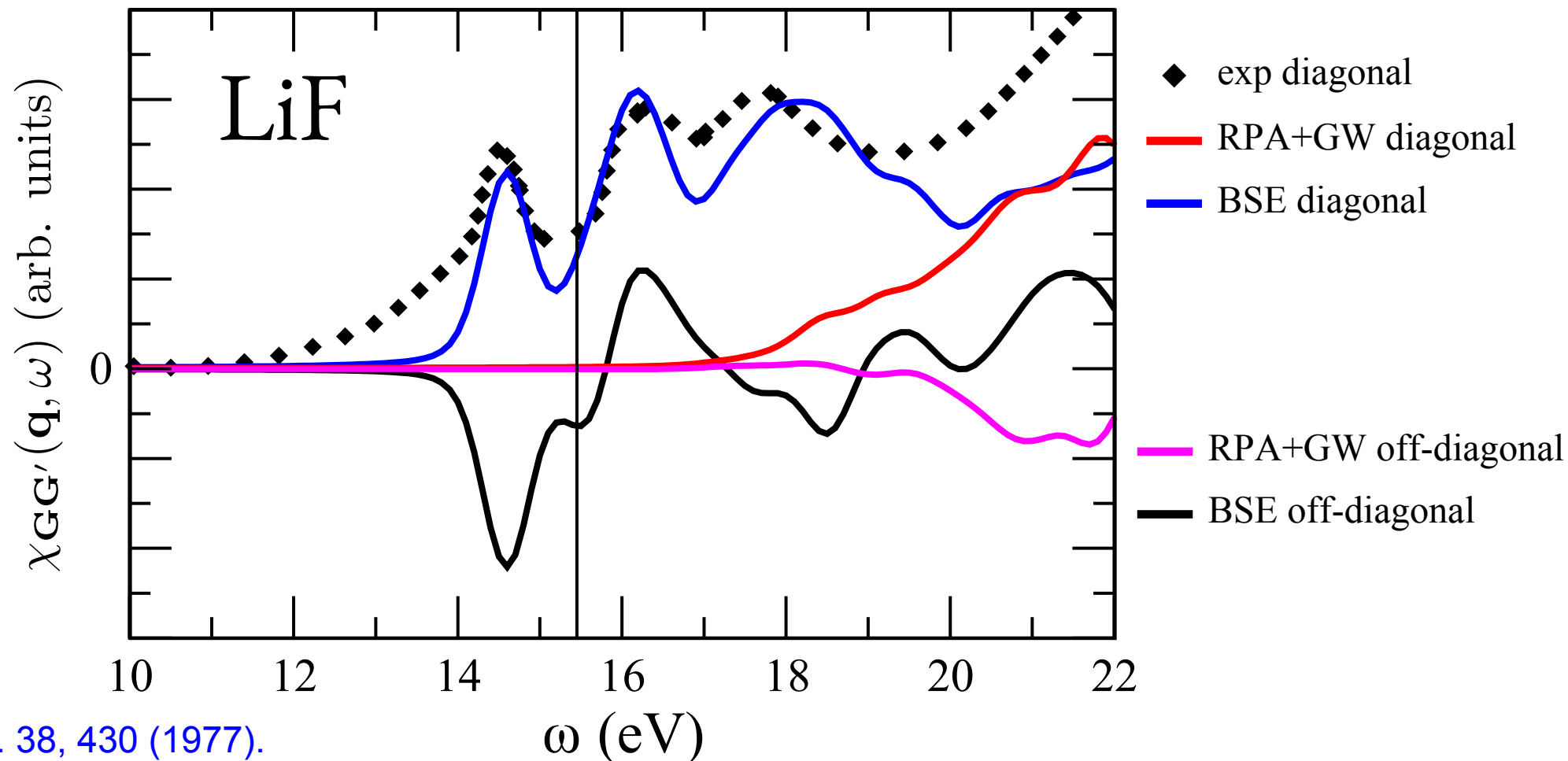
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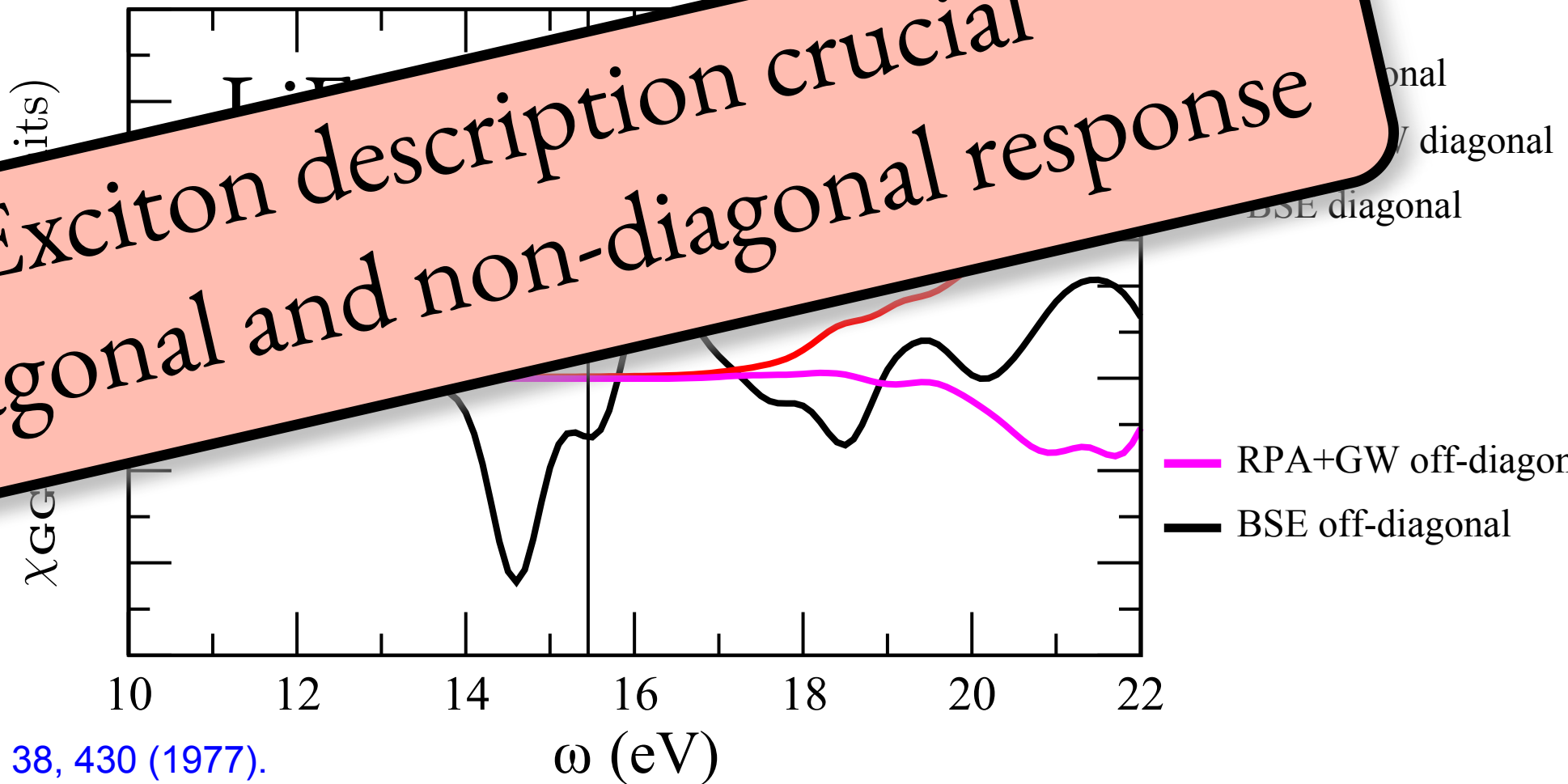
 Fields *et al.* Phys. Rev. Lett. **38**, 430 (1977).



Coherent Inelastic X-ray scattering

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Exciton description crucial for diagonal and non-diagonal response

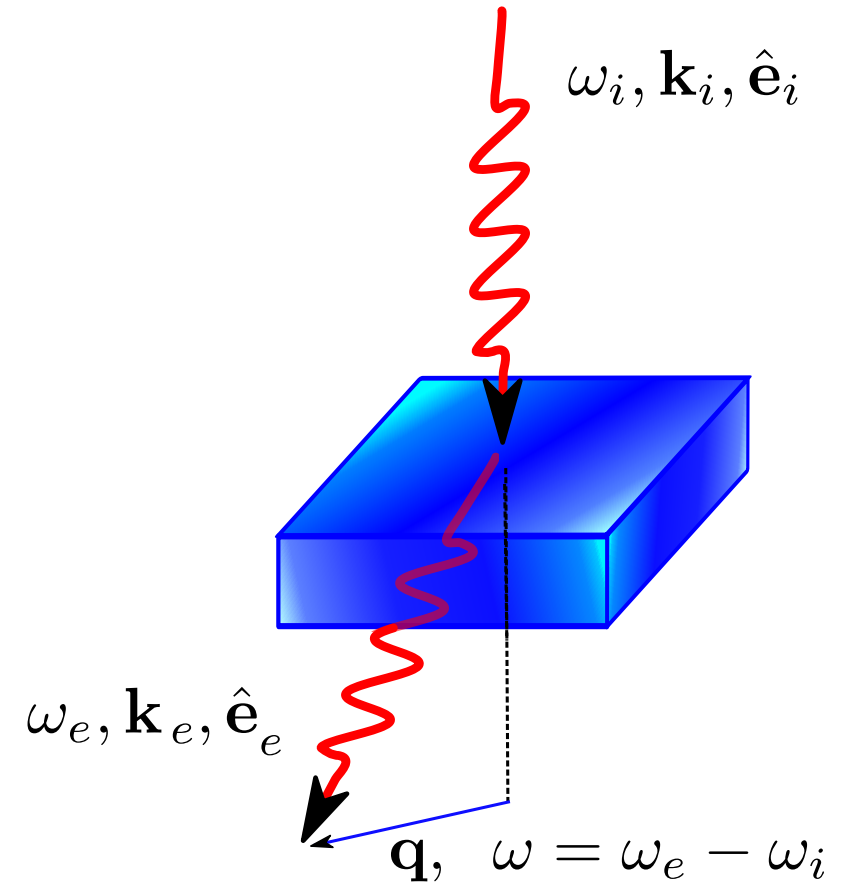


Igor Re...
Phys. R...
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X-ray scattering

non-Resonant IXS

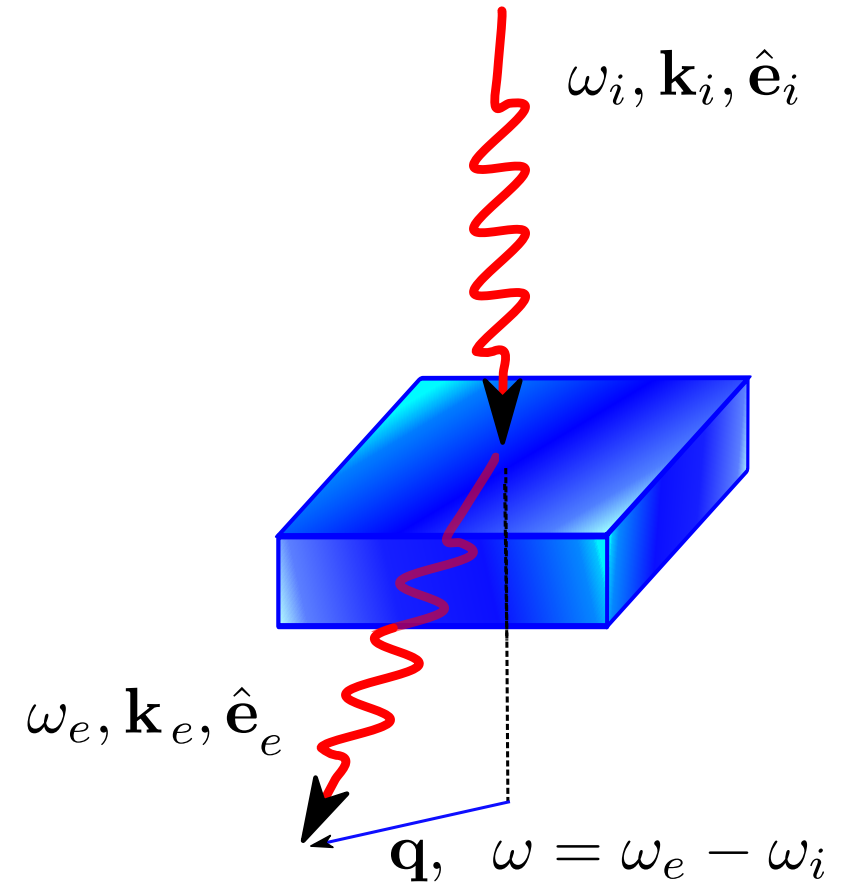


$$\frac{d^2\sigma}{d\Omega_2 d\omega_e} \propto \sum_f \left| \langle f | e^{i\mathbf{q}\cdot\mathbf{r}} | 0 \rangle - \frac{i\omega_i/e}{2mc^2} \sum_n \frac{\langle f | e^{-i\mathbf{k}_f\cdot\mathbf{r}} \nabla | n \rangle \langle n | e^{i\mathbf{k}_i\cdot\mathbf{r}} \nabla | 0 \rangle}{\omega_i - (E_n - E_0)} \right|^2 \times \delta(\omega - (E_f - E_0))$$

X-ray scattering

non-Resonant IXS

Resonant IXS



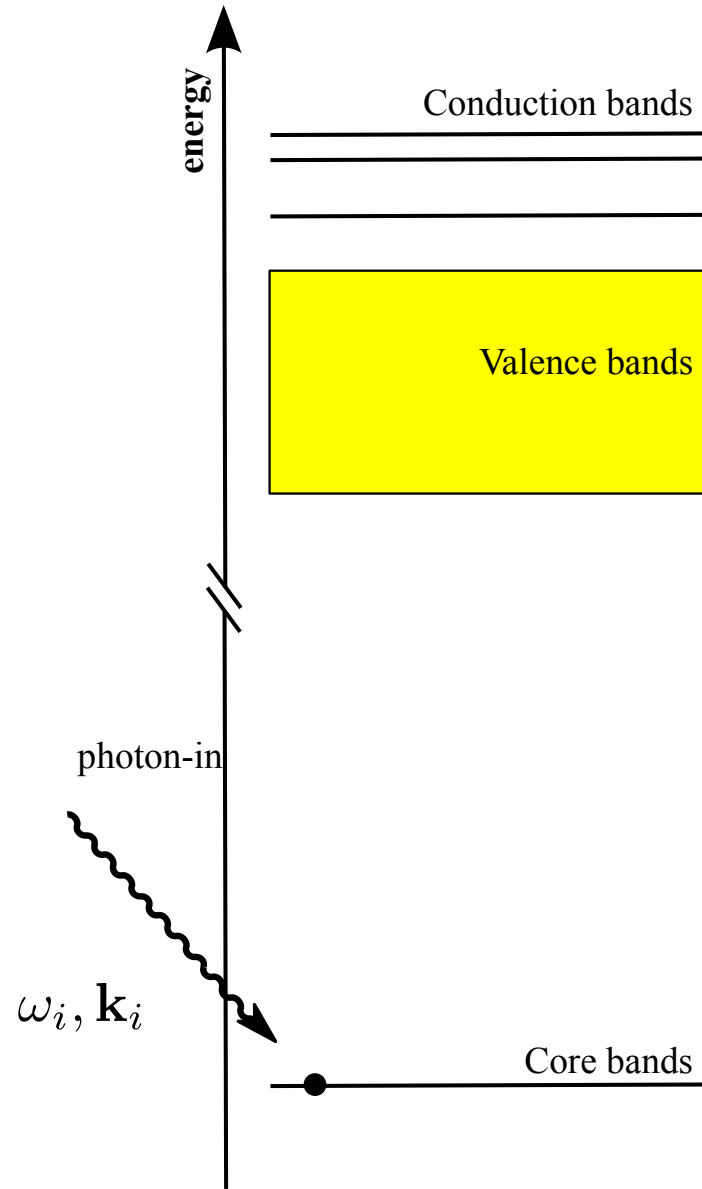
$$\frac{d^2\sigma}{d\Omega_2 d\omega_e} \propto \sum_f \left| \langle f | e^{i\mathbf{q}\cdot\mathbf{r}} | 0 \rangle - \frac{i\omega_i/e}{2mc^2} \sum_n \frac{\langle f | e^{-i\mathbf{k}_f\cdot\mathbf{r}} \nabla | n \rangle \langle n | e^{i\mathbf{k}_i\cdot\mathbf{r}} \nabla | 0 \rangle}{\omega_i - (E_n - E_0)} \right|^2 \times \delta(\omega - (E_f - E_0))$$

Resonant IXS

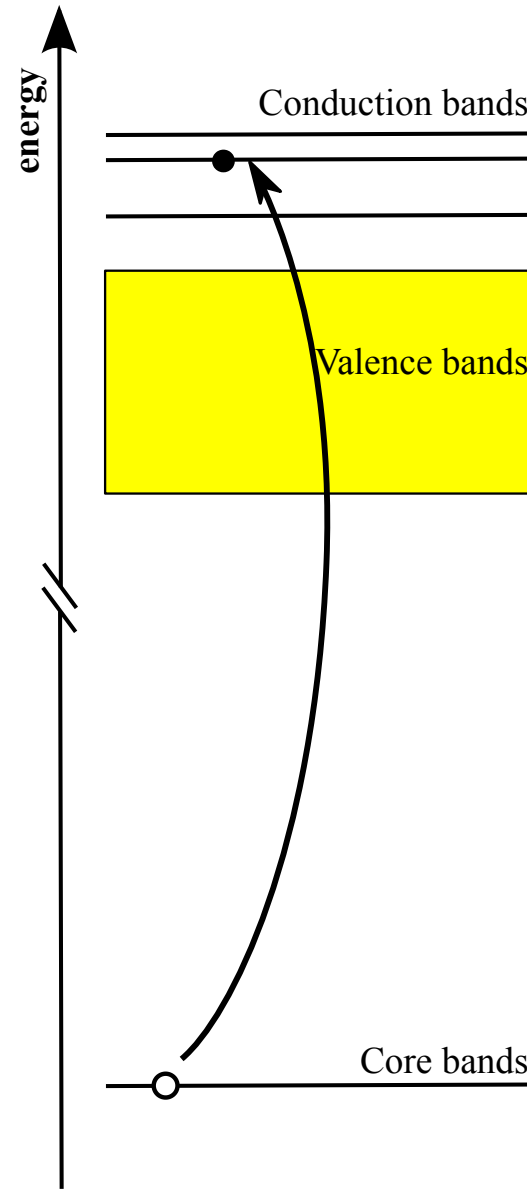


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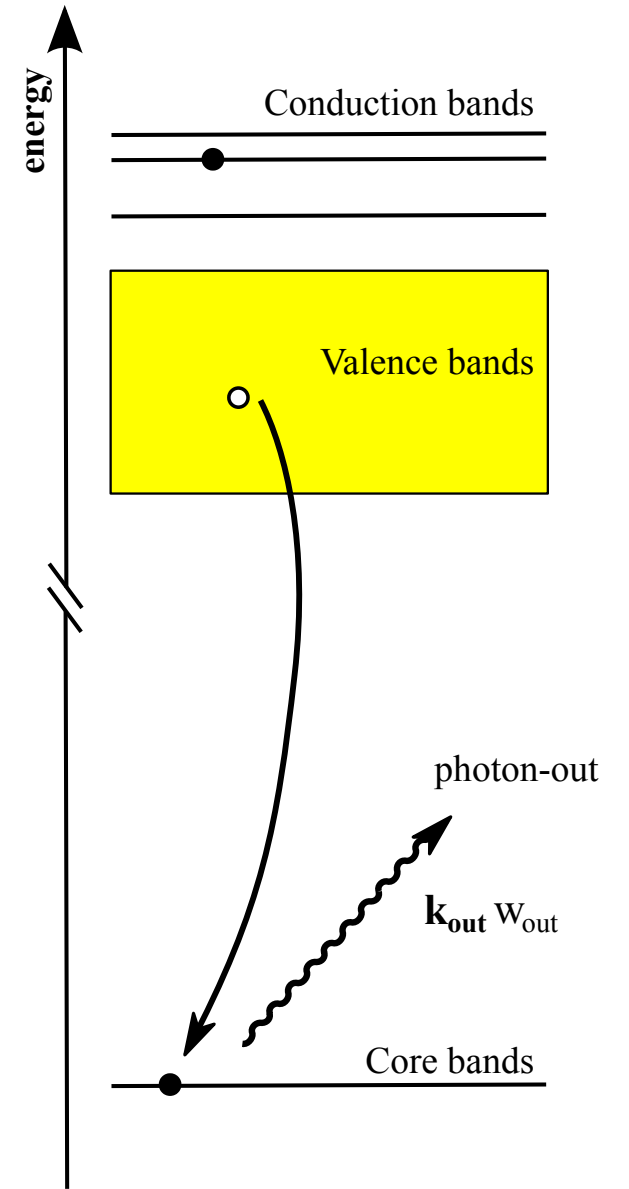
Initial state



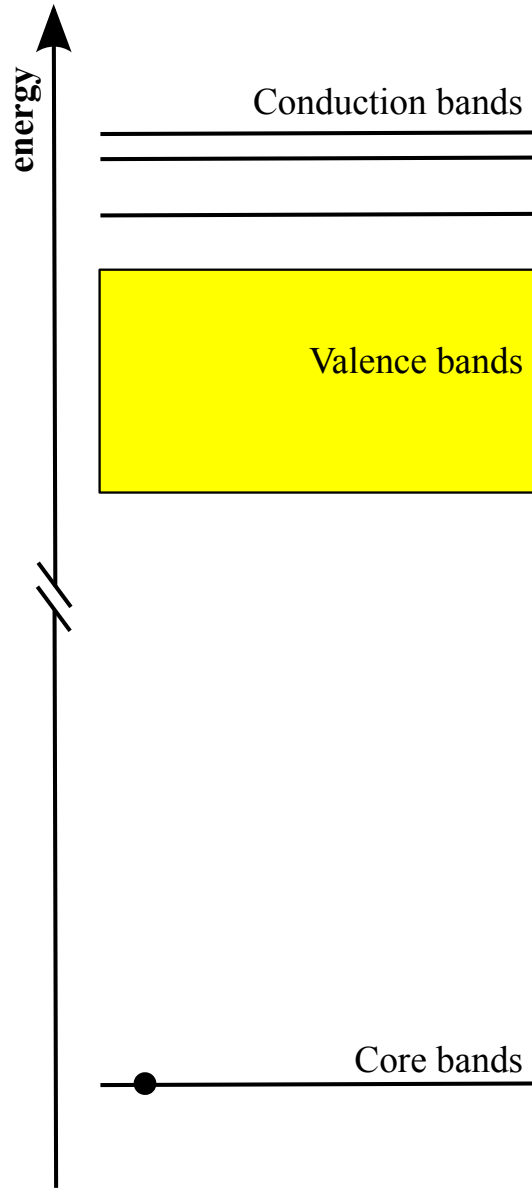
Intermediate state



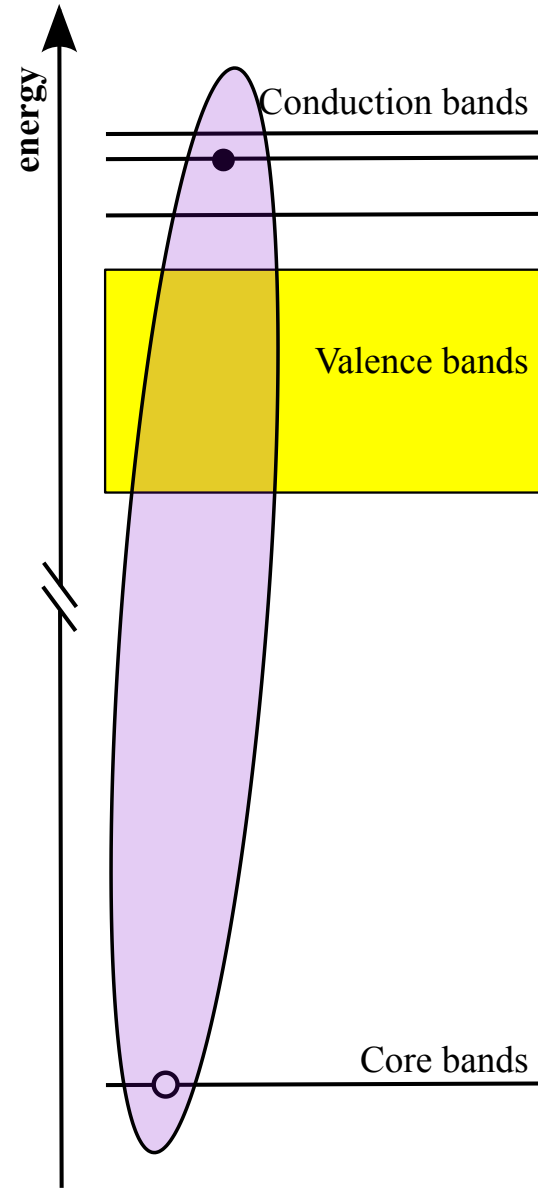
Final state



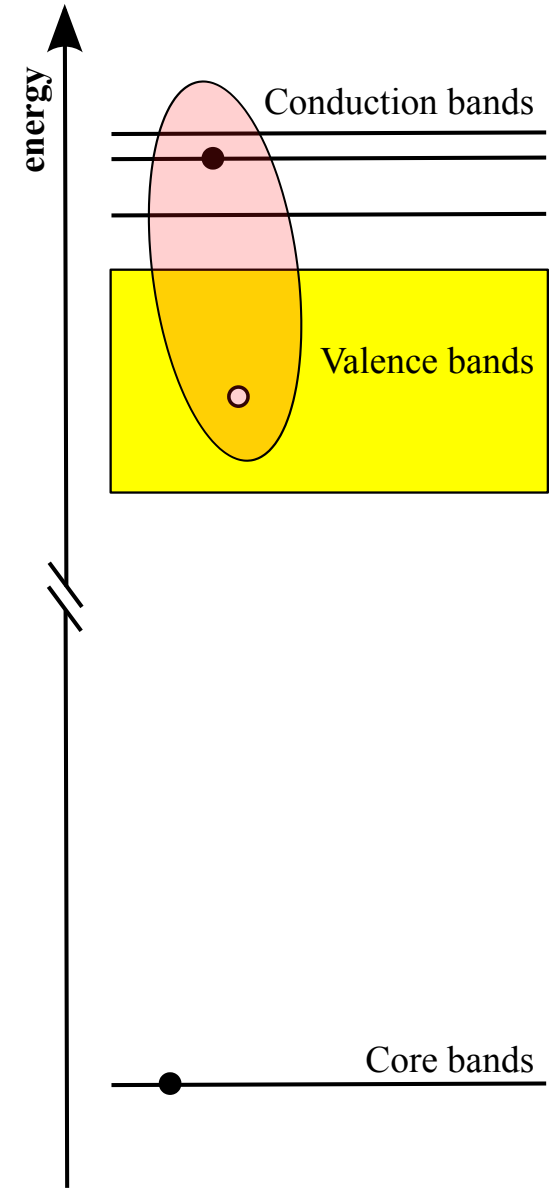
Initial state



Intermediate state



Final state



Resonant Inelastic X-ray scattering via excitonic pathways in BSE

$$\frac{d^2\sigma}{d\Omega_2 d\omega_e} \propto \sum_f \left| \sum_n \frac{\langle f | e^{-i\mathbf{k}_f \cdot \mathbf{r}} \nabla | n \rangle \langle n | e^{i\mathbf{k}_i \cdot \mathbf{r}} \nabla | 0 \rangle}{\omega_i - (E_n - E_0)} \right|^2 \times \delta(\omega - (E_f - E_0))$$



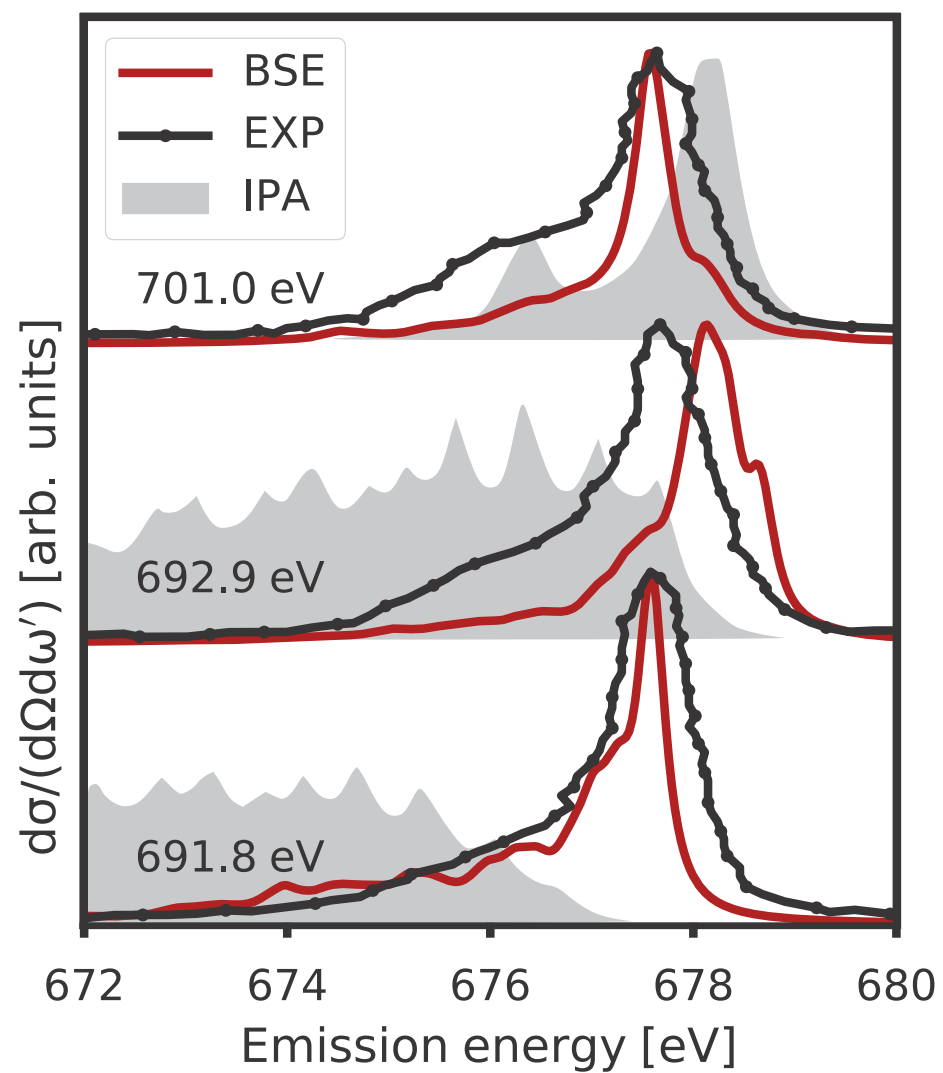
$$\frac{d^2\sigma}{d\Omega_2 d\omega_e} \propto \text{Im} \sum_{\substack{\mu\mu'' \\ \lambda'_c \lambda_c \lambda}} \sum_{\substack{vv' \\ cc''}} \left[\frac{t_{\lambda'_c}^{(1)} A_{\lambda'_c}^{\mu c} \tilde{\rho}_{\mu v}}{\omega_i - E_{\lambda'_c} + i\eta} \right]^* \frac{A_{\lambda}^{vc} A_{\lambda}^{*v'c''}}{\omega - E_{\lambda} + i\eta} \left[\frac{\tilde{\rho}_{\mu''v'}^* A_{\lambda_c}^{\mu''c''} t_{\lambda_c}^{(1)}}{\omega_i - E_{\lambda_c} + i\eta} \right]$$



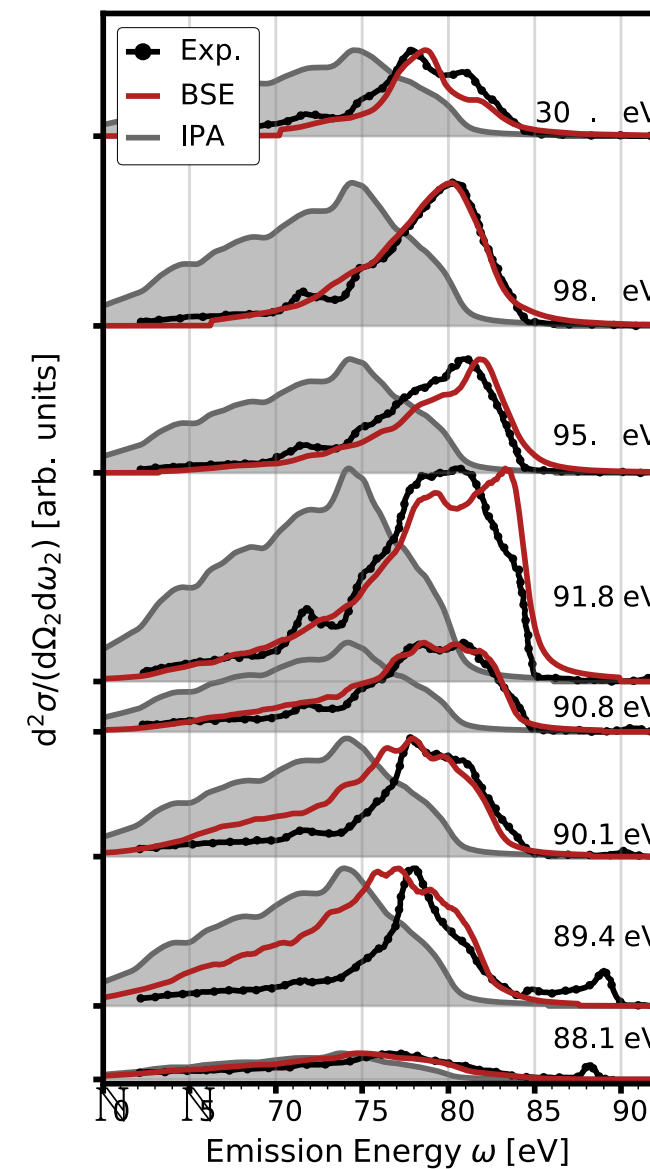
Vorwerk, Sottile, Draxl, Phys. Rev. Research **2**, 042003(R) (2020)

<https://github.com/exciting/BRIXS>

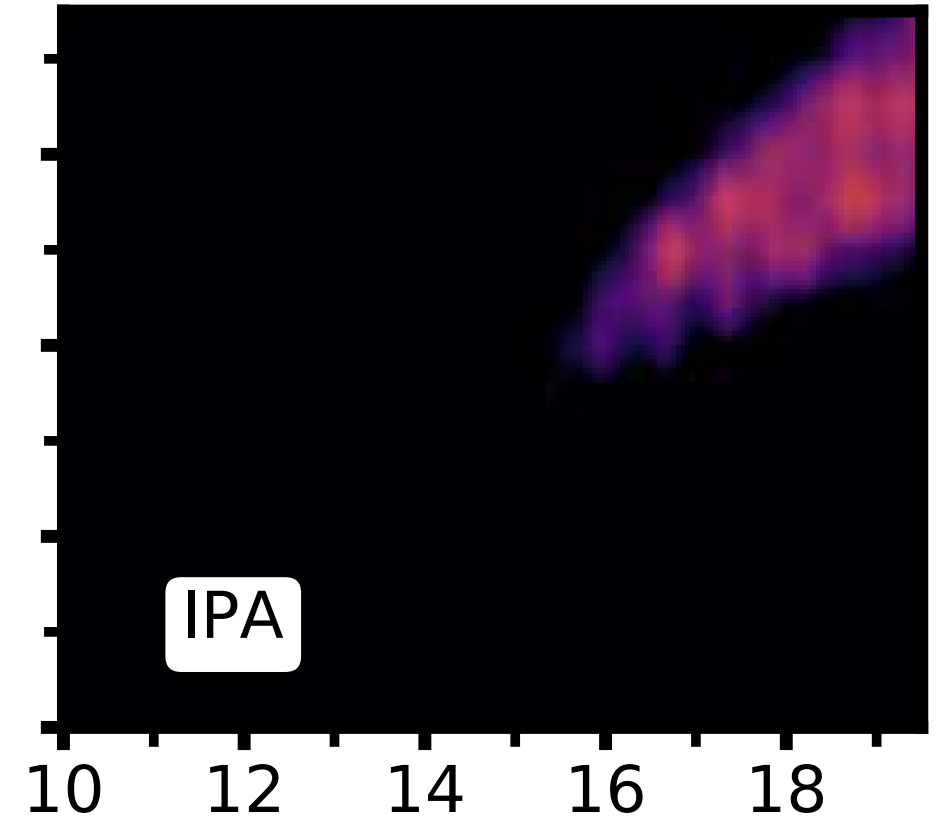
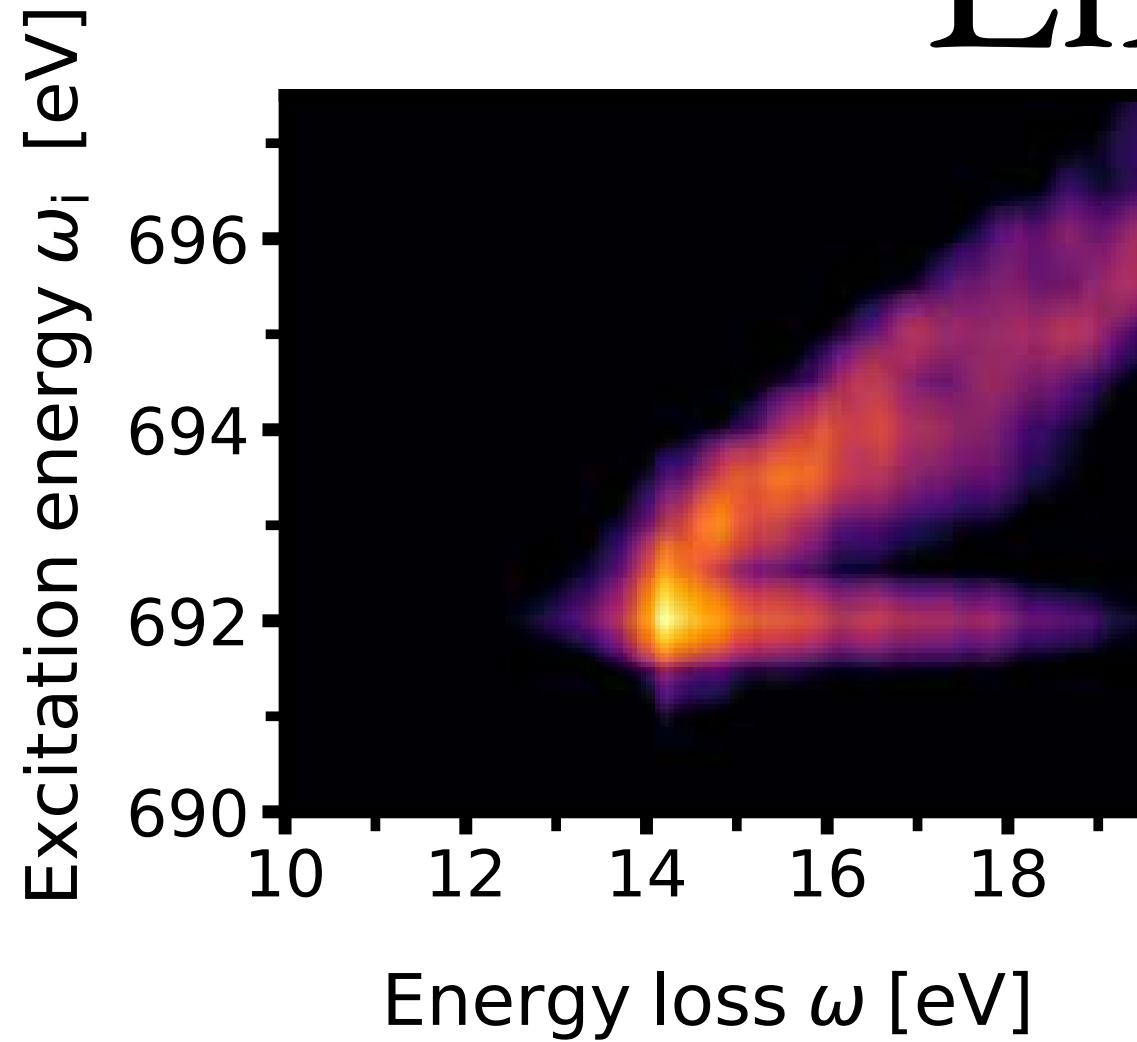
LiF



Diamond



LiF



Conclusions

- Absorption and RIXS accurate within BSE
- Ab initio and predictive, and permits analysis

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- Ab initio and predictive, and permits analysis
- Rely on the description of the initial (ground) state
- Cumbersome calculations

Perspectives (with XFEL in mind)

- Tackle small (fragile) systems
probe-before-damage



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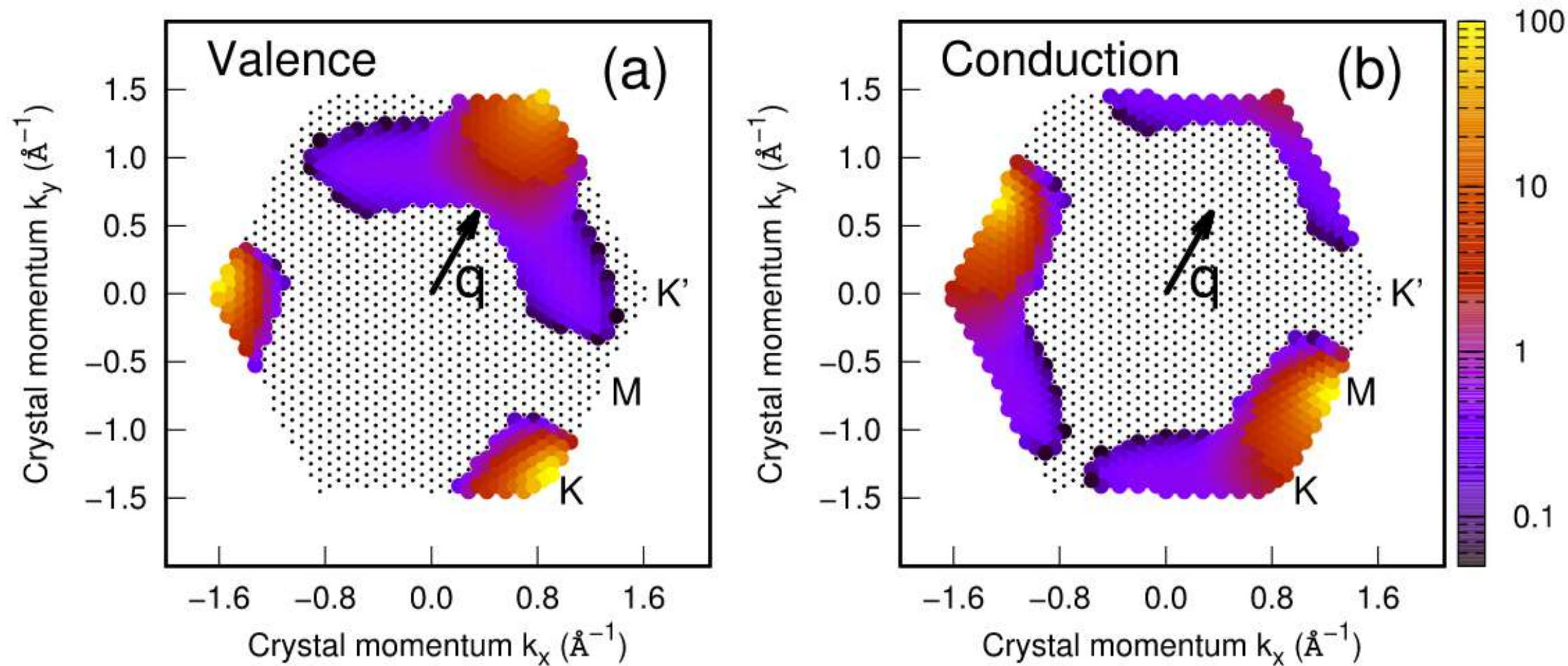
- Tackle small (fragile) systems
probe-before-damage
- Develop GFs theory (ab-initio with excitonic effects)
 - time-dependent RIXS (fs resolution)
 - stimulated RIXS (non-linear regime)

Thanks to the Theoretical Spectroscopy Group



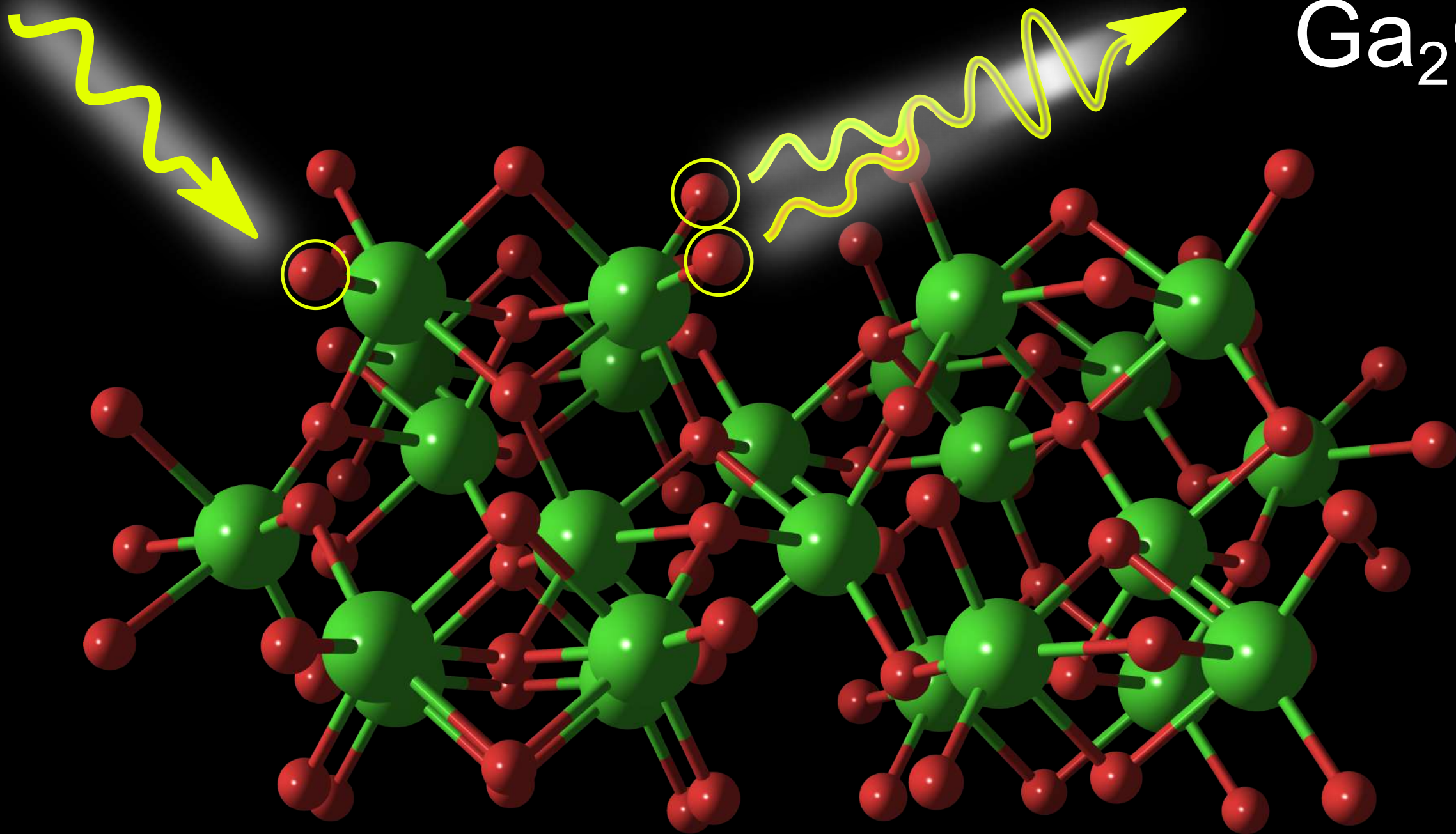
and to You

Excitons in hexagonal BN



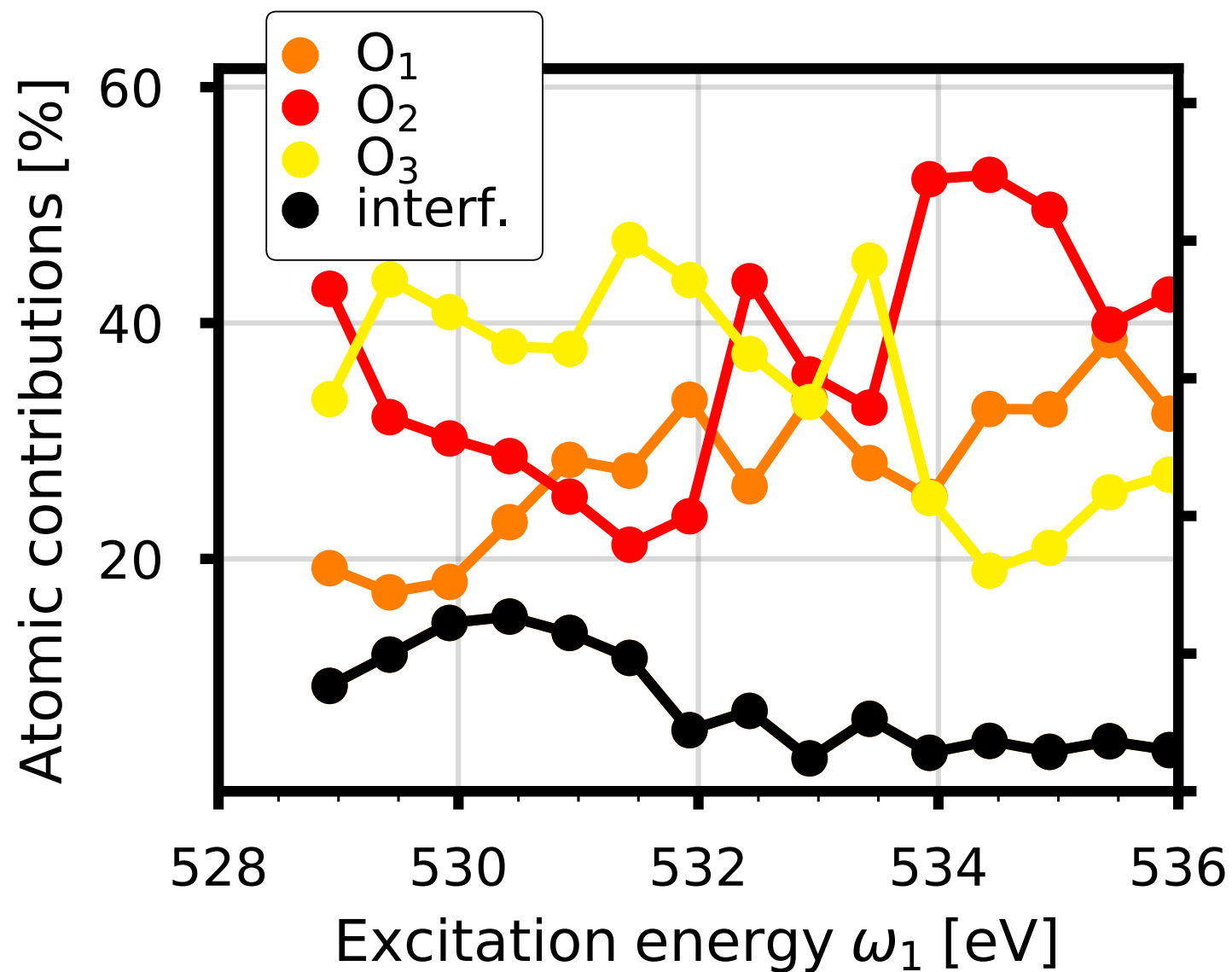
Sponza *et al.* Phys. Rev. B **97**, 075121 (2018)

Ga_2O_3



O-K Ga_2O_3

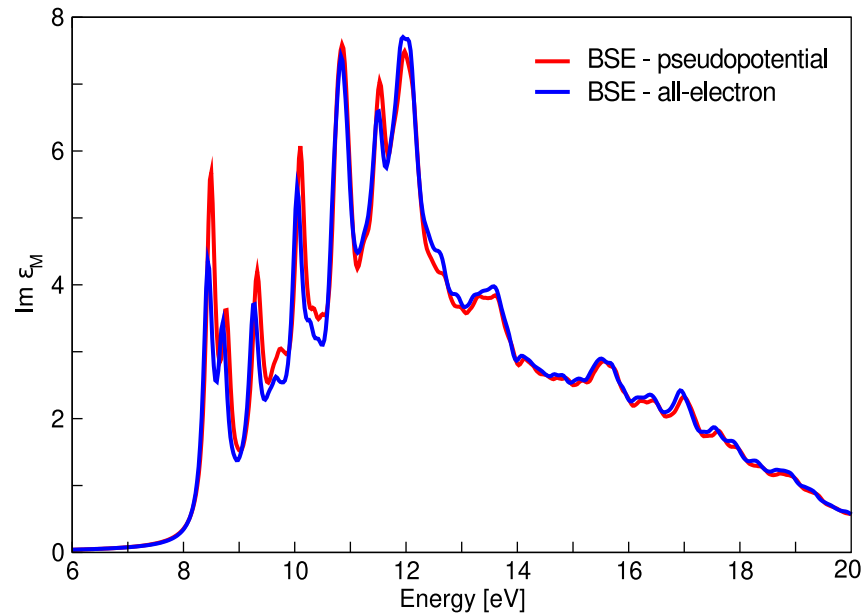
3 inequivalent oxigens



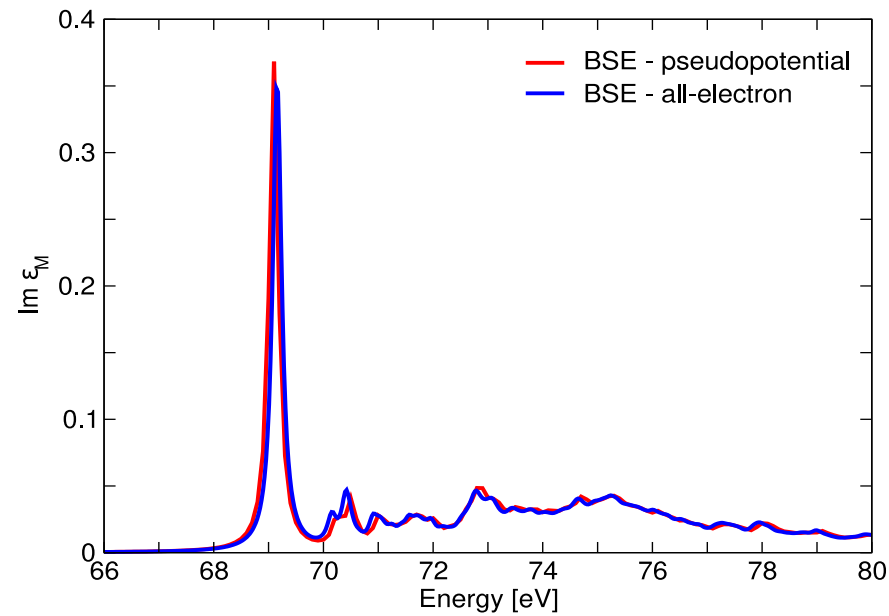
Vorwerk *et al.* Phys. Chem. Chem. Phys. **24**, 17439 (2022).

Optical and X-ray absorption of Al_2O_3

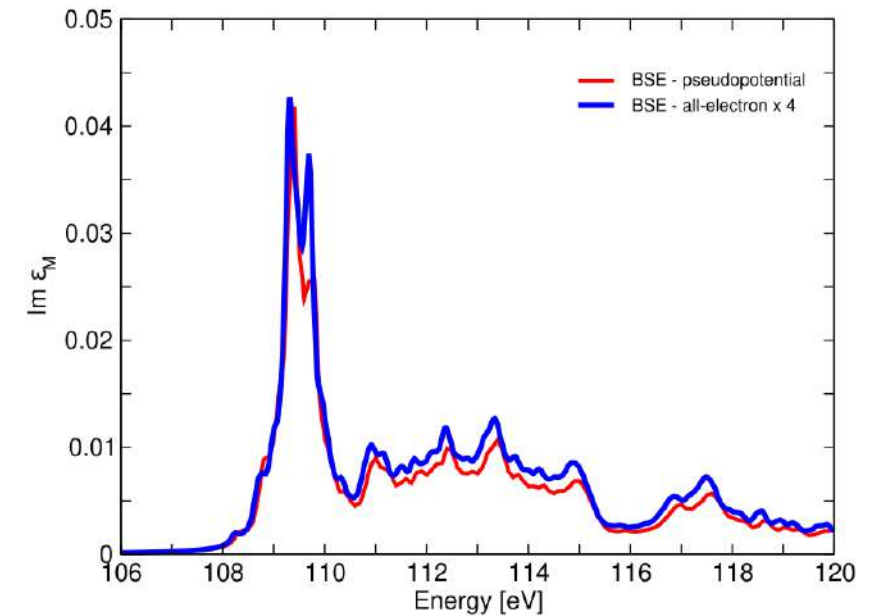
All-electron vs pseudo-potential



optical



$L_{2,3}$



L_1



